

Quantum corrections to neutrino mixing in the MSSM with righthanded neutrinos

International School Cargese 2012: Across the TeV frontier with the LHC

Wolfgang Gregor Hollik | August 2012

INSTITUT FÜR THEORETISCHE TEILCHENPHYSIK | KIT CAMPUS SÜD



Motivation

- Neutrinos seem to have mass
- Oscillations:
 - $\Delta m_{21}^2 = 7.58 \times 10^{-5} \text{ eV}^2$
 - $\Delta m_{31}^2 = 2.35 \times 10^{-3} \text{ eV}^2$
 - large mixing angles



- Unknown: Absolute neutrino mass scale \rightarrow KATRIN
 - possible upper limit: 0.2 eV, discovery: 0.35 eV

Compare CKM and PMNS mixing matrix:

$$V_{\text{CKM}} = \begin{pmatrix} \bullet & \cdot & \cdot \\ \cdot & \bullet & \cdot \\ \cdot & \cdot & \bullet \end{pmatrix} \quad U_{\text{PMNS}} = \begin{pmatrix} \bullet & \cdot & \cdot \\ \cdot & \bullet & \cdot \\ \cdot & \cdot & \bullet \end{pmatrix}$$

Neutrino masses and seesaw

Standard Model + righthanded Neutrinos = Seesaw Type I

$$-\mathcal{L}_{\nu, \text{mass}} = \underbrace{\bar{\nu}_L m_D \nu_R}_{\text{Dirac mass}} + \frac{1}{2} \underbrace{\overline{\nu_L^c} m_R \nu_R}_{\text{Majorana mass}} + \text{h. c.}$$



Neutrino masses and seesaw

Standard Model + righthanded Neutrinos = Seesaw Type I

$$-\mathcal{L}_{\nu, \text{mass}} = \underbrace{\bar{\nu}_L m_D \nu_R}_{\text{Dirac mass}} + \frac{1}{2} \underbrace{\overline{\nu_L^c} m_R \nu_R}_{\text{Majorana mass}} + \text{h. c.}$$

Neutrino mass matrix:

$$\mathcal{M}_\nu = \begin{pmatrix} 0 & m_D \\ m_D^T & m_R \end{pmatrix}.$$

What about m_R ?

- righthanded neutrinos are SM singlets \rightarrow no constraint for mass
- seesaw  : $m_\nu = -m_D m_R^{-1} m_D \approx \mathcal{O}(0.1 \text{ eV})$
- assumption: Dirac mass of order EW scale ($\mathcal{O}(10 \dots 100 \text{ GeV})$):
 $m_R \sim \mathcal{O}(10^{13 \dots 14} \text{ GeV})$

The MSSM with righthanded neutrinos

Superpotential of the ν MSSM

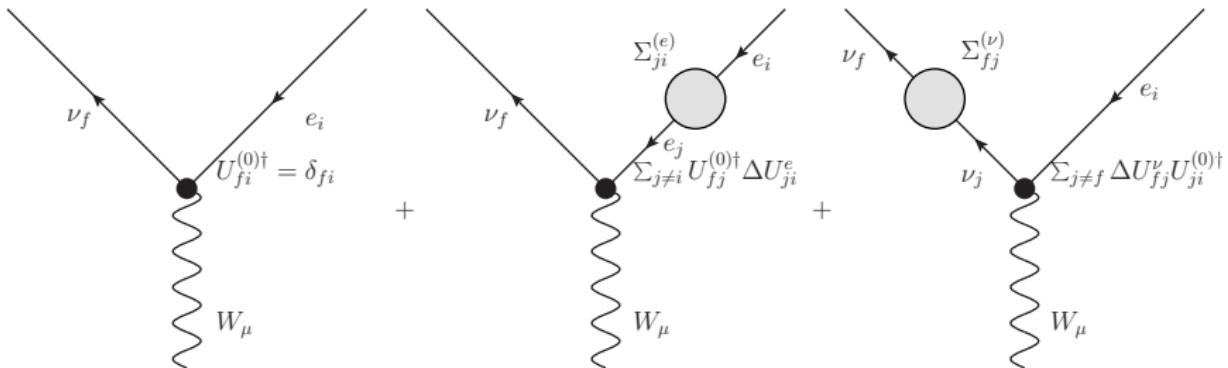
$$\mathcal{W}^\ell = \mu H_d \cdot H_u - Y_\ell^{IJ} H_d \cdot L_L^I E_R^J + Y_\nu^{IJ} H_u \cdot L_L^I N_R^J + \frac{1}{2} m_R^{IJ} N_R^I N_R^J,$$

with $L_L = (\ell_L, \tilde{\ell}_L) \in SU(2)_L$ and $E_R = (e_L^c, \tilde{e}_R^*)$, $N_R = (\nu_L^c, \tilde{\nu}_R^*)$.

Soft-breaking terms

$$\begin{aligned} \mathcal{V}_{\text{soft}} = & (\mathcal{M}_{\tilde{\ell}}^2)^{IJ} \tilde{L}_L^{I*} \tilde{L}_L^J + (\mathcal{M}_{\tilde{e}}^2)^{IJ} \tilde{e}_R^I \tilde{e}_R^{J*} + (\mathcal{M}_{\tilde{\nu}}^2)^{IJ} \tilde{\nu}_R^I \tilde{\nu}_R^{J*} \\ & - \left[(B_\nu)^{IJ} \tilde{\nu}_R^{I*} \tilde{\nu}_R^{J*} + A_e^{IJ} H_1 \cdot \tilde{L}_L^I \tilde{e}_R^{J*} - A_\nu^{IJ} H_2 \cdot \tilde{L}_L^I \tilde{\nu}_R^{J*} + \text{h.c.} \right], \end{aligned}$$

radiative flavour violation in the lepton sector



PMNS matrix renormalization

$$i \frac{g}{\sqrt{2}} \gamma^\mu P_L U_{\text{PMNS}}^\dagger \rightarrow i \frac{g}{\sqrt{2}} \gamma^\mu P_L (\mathbb{1} + \Delta U^e + \Delta U^\nu),$$

flavour changing self energies and sensitivity to neutrino mass

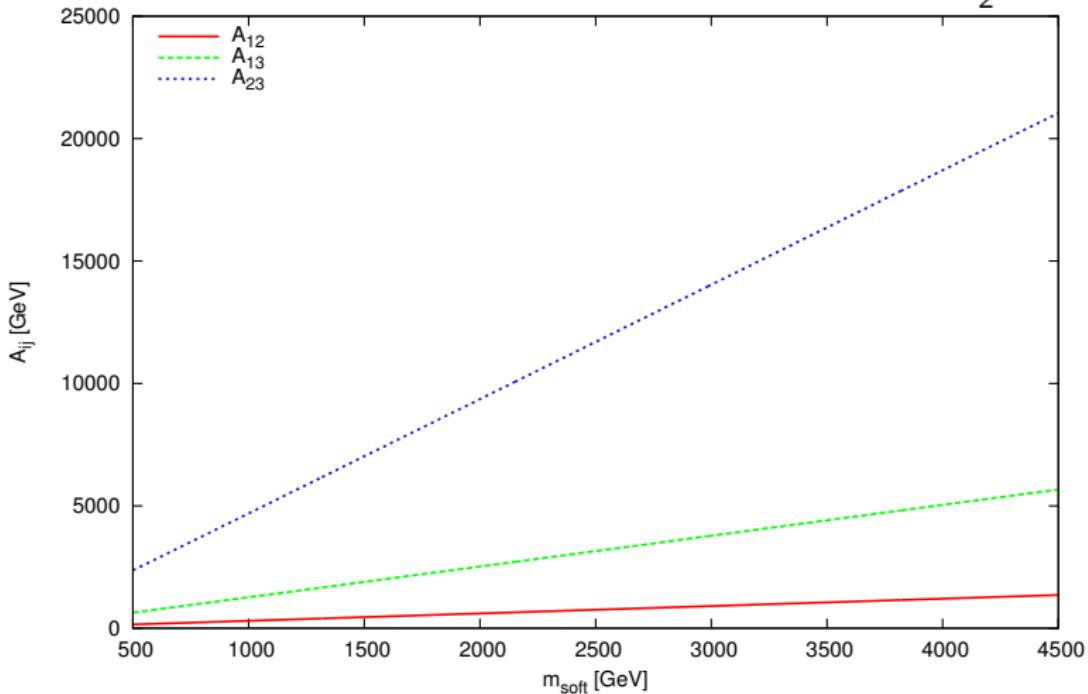
$$\Delta U_{fi}^\nu \sim \frac{m_{\nu_f} \Sigma_{fi}}{\Delta m_\nu^2}$$

Numerical results

Flavour violation only in the neutrino A-Terms!

$$\mathcal{M}_{\text{soft}}^2 = m_0 \mathbb{1}.$$
$$M_1 = \frac{1}{2} M_2 = m_0.$$

trilinear flavour changing coupling A_{ij}

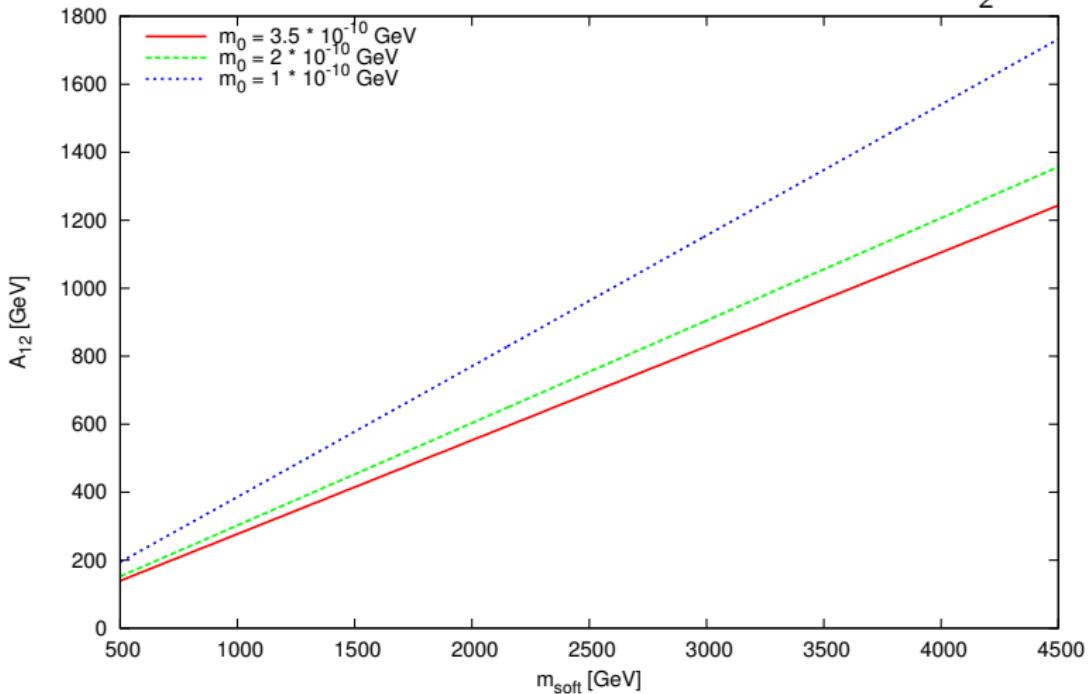


Numerical results

Flavour violation only in the neutrino A-Terms!

$$\mathcal{M}_{\text{soft}}^2 = m_0 \mathbb{1}.$$
$$M_1 = \frac{1}{2} M_2 = m_0.$$

compare neutrino spectra



Conclusion

- supersymmetric seesaw incorporates additional flavour structures
- seesaw-like structure in sneutrino squared mass matrix
- sensitive to the degree of degeneracy of neutrino mass spectrum
- sensitive to scale of righthanded neutrinos
- possibly radiative generation of neutrino mixing
- no severe enhancement of LFV obervables

Conclusion

- supersymmetric seesaw incorporates additional flavour structures
- seesaw-like structure in sneutrino squared mass matrix
- sensitive to the degree of degeneracy of neutrino mass spectrum
- sensitive to scale of righthanded neutrinos
- possibly radiative generation of neutrino mixing
- no severe enhancement of LFV obervables



Backup

Slides

effects on sneutrino mass matrix

- charged slepton mass matrix as in the MSSM
- sneutrino mass matrix in the MSSM: simple

$$\mathcal{M}_{\tilde{\nu}}^2 = \begin{pmatrix} \mathcal{M}_{\tilde{\ell}}^2 + M_Z^2 T_{3L}^{\tilde{\nu}} \cos 2\beta \mathbb{1} & 1 \\ 1 & 0 \end{pmatrix}$$

- Majorana mass term $\nu_R^T m_R \nu_R$ inflates sneutrino mass matrix:
additional terms $\sim \tilde{\nu}_R \tilde{\nu}_R, \tilde{\nu}_R^* \tilde{\nu}_R^*$

$$\mathcal{M}_{\tilde{\nu}}^2 = \begin{pmatrix} \mathcal{M}_{L^* L}^2 & \mathcal{M}_{L^* L^*}^2 & \mathcal{M}_{L^* R^*}^2 & \mathcal{M}_{L^* R}^2 \\ \mathcal{M}_{LL}^2 & \mathcal{M}_{LL^*}^2 & \mathcal{M}_{LR^*}^2 & \mathcal{M}_{LR}^2 \\ \mathcal{M}_{RL}^2 & \mathcal{M}_{RL^*}^2 & \mathcal{M}_{RR^*}^2 & \mathcal{M}_{RR}^2 \\ \mathcal{M}_{R^* L}^2 & \mathcal{M}_{R^* L^*}^2 & \mathcal{M}_{R^* R^*}^2 & \mathcal{M}_{R^* R}^2 \end{pmatrix}$$

12 × 12-Matrix

effects on sneutrino mass matrix

- charged slepton mass matrix as in the MSSM
- sneutrino mass matrix in the MSSM: simple

$$\mathcal{M}_{\tilde{\nu}}^2 = \begin{pmatrix} \mathcal{M}_{\tilde{\ell}}^2 + M_Z^2 T_{3L}^{\tilde{\nu}} \cos 2\beta & 1 \\ 1 & 0 \end{pmatrix}$$

- Majorana mass term $\nu_R^T m_R \nu_R$ inflates sneutrino mass matrix:
additional terms $\sim \tilde{\nu}_R \tilde{\nu}_R, \tilde{\nu}_R^* \tilde{\nu}_R^*$

$$\mathcal{M}_{\tilde{\nu}}^2 = \begin{pmatrix} \mathcal{M}_{LL}^2 & \mathcal{M}_{LR}^2 \\ (\mathcal{M}_{LR}^2)^\dagger & \mathcal{M}_{RR}^2 \end{pmatrix}$$

12 × 12-Matrix

full sneutrino squared mass matrix in the ν MSSM

$$\mathcal{M}_{\tilde{\nu}}^2 = \frac{1}{2} \begin{pmatrix} \mathcal{M}_{LL}^2 & \mathcal{M}_{LR}^2 \\ (\mathcal{M}_{LR}^2)^\dagger & \mathcal{M}_{RR}^2 \end{pmatrix}$$

$$\mathcal{M}_{LL}^2 = \begin{pmatrix} \mathcal{M}_{\tilde{\ell}}^2 + \frac{1}{2} M_Z^2 \cos 2\beta \mathbf{1} + \mathbf{m}_\nu \mathbf{m}_\nu^\dagger & \mathbf{0} \\ \mathbf{0} & (\searrow)^* \end{pmatrix},$$

$$\mathcal{M}_{RL}^2 = \begin{pmatrix} \frac{1}{2} \mathbf{m}_\nu \mathbf{m}_R & -\mu \cot \beta \mathbf{m}_\nu - v_2 \mathbf{A}_\nu^* \\ -\mu^* \cot \beta \mathbf{m}_\nu^* - v_2 \mathbf{A}_\nu & \frac{1}{2} \mathbf{m}_\nu^* \mathbf{m}_R^* \end{pmatrix},$$

$$\mathcal{M}_{RR}^2 = \begin{pmatrix} (\mathcal{M}_{\tilde{\nu}}^2)^T + \mathbf{m}_\nu^T \mathbf{m}_\nu^* + \frac{1}{2} \mathbf{m}_R^* \mathbf{m}_R & -2 \mathbf{B}^* \\ -2 \mathbf{B} & \mathcal{M}_{\tilde{\nu}}^2 + \mathbf{m}_\nu^\dagger \mathbf{m}_\nu + \frac{1}{2} \mathbf{m}_R \mathbf{m}_R^* \end{pmatrix}.$$

effective sneutrino mass matrix

$$\mathcal{M}_{\tilde{\nu}\ell}^2 = \begin{pmatrix} \mathbf{m}_{\Delta L=0}^2 & (\mathbf{m}_{\Delta L=2}^2)^* \\ \mathbf{m}_{\Delta L=2}^2 & (\mathbf{m}_{\Delta L=0}^2)^* \end{pmatrix} + \mathcal{O}(M_{\text{SUSY}}^2 m_R^{-2}),$$

$$\begin{aligned} \mathbf{m}_{\Delta L=0}^2 &= \text{MSSM} + \mathbf{m}_\nu^D \mathbf{m}_\nu^{D\dagger} - \mathbf{m}_\nu^D \mathbf{m}_R (\mathbf{m}_R^2 + \mathcal{M}_{\tilde{\nu}}^2)^{-1} \mathbf{m}_R \mathbf{m}_\nu^D, \\ \mathbf{m}_{\Delta L=2}^2 &= X_\nu \mathbf{m}_\nu^D (\mathbf{m}_R^2 + \mathcal{M}_{\tilde{\nu}}^2)^{-1} \mathbf{m}_R \mathbf{m}_\nu^{DT} + (\rightarrow)^T \\ &\quad - 2 \mathbf{m}_\nu^{D*} \mathbf{m}_R \left[\mathbf{m}_R^2 + (\mathcal{M}_{\tilde{\nu}}^2)^T \right]^{-1} \mathbf{B} (\mathbf{m}_R^2 + \mathcal{M}_{\tilde{\nu}}^2)^{-1} \mathbf{m}_R \mathbf{m}_\nu^{D\dagger}. \end{aligned}$$

$$X_\nu \mathbf{m}_\nu^D = -\mu^* \cot \beta \mathbf{m}_\nu^{D*} - v_2 \mathbf{A}_\nu$$