

To Higgs or not to Higgs
vacuum stability and the origin of mass

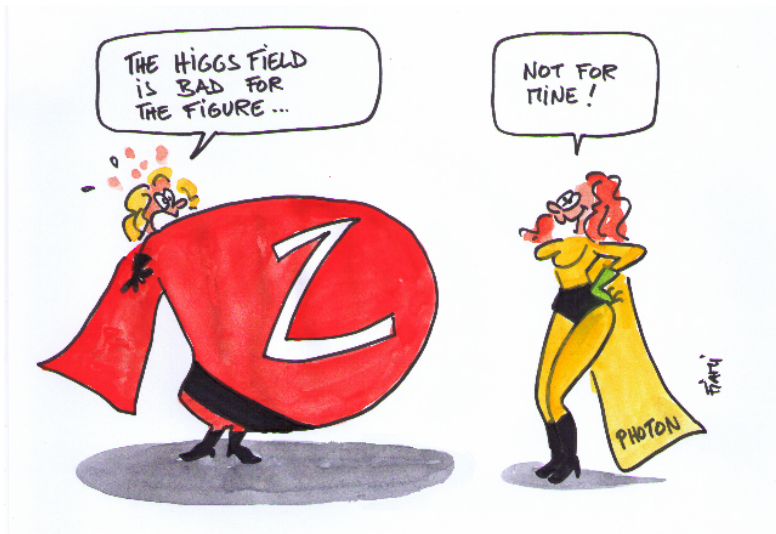


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Dec 12 2016 | MU Programtag Mainz

The Higgs mechanism and the origin of mass



[CERN bulletin]



The tale of the Higgs dale

- potential instable at the origin
- ground state (= *vacuum*) breaks symmetry

Massive Gauge Bosons

spontaneously broken $SU(2)_L \times U(1)_Y$

$$\Phi(x) = \begin{pmatrix} \phi^+ \\ \phi^0 \end{pmatrix} \rightarrow \begin{pmatrix} 0 \\ v + \frac{1}{\sqrt{2}}h(x) \end{pmatrix}$$

$(D_\mu \phi)^\dagger D^\mu \phi \hookrightarrow$ quadratic terms for gauge fields:

$$\frac{1}{4}v^2 (gA_\mu^3 - g'B_\mu) (gA^{3\mu} - g'B^\mu) + \frac{1}{2}g^2v^2 A_\mu^+ A^{-\mu},$$

$$\text{with } A_\mu^\pm = \frac{1}{\sqrt{2}}(A_\mu^1 \pm iA_\mu^2).$$

Masses for W and Z

$$m_W^2 = \frac{v^2 g^2}{2}, \quad m_Z^2 = \frac{v^2}{2} (g^2 + g'^2),$$

weak mixing angle

$$\tan \theta_w = \frac{g'}{g}, \quad m_W = \cos \theta_w m_Z.$$

Fundamental masses

all fundamental masses $\sim v = 246 \text{ GeV}$ in the Standard Model

Masses for fermions

Yukawa interactions:

$$\mathcal{L}_{\text{Yuk}} = Y \bar{\Psi}_L \cdot \Phi \Psi_R + \text{h.c.} = Y \begin{pmatrix} \bar{\psi}_L^u & \bar{\psi}_L^d \end{pmatrix} \cdot \begin{pmatrix} 0 \\ v/\sqrt{2} \end{pmatrix} \psi_R^d + \text{h.c.}$$

(for up-type fields: $\Phi \rightarrow \tilde{\Phi} = i\sigma_2 \Phi^*$ and $\psi_R^d \rightarrow \psi_R^u$)

$$m_\psi = \frac{v}{\sqrt{2}} Y$$

In the Standard Model, by construction,

$$V_{\text{SM}} = -\mu^2 H^\dagger H + \lambda (H^\dagger H)^2$$

leads to $v^2 = \mu^2/\lambda$ (with $\mu^2 > 0$, $\lambda > 0$).

- measuring gauge boson masses + gauge couplings: v^2
- measuring Higgs boson mass: λ

So what?

- only tree level
- no new physics?
- V_{SM} put by hand...
- going beyond tree and SM: minimum may get unstable
↔ “new” minimum: different v = different mass
may break additional “good” symmetries of the SM

The Standard Model (In)Stability

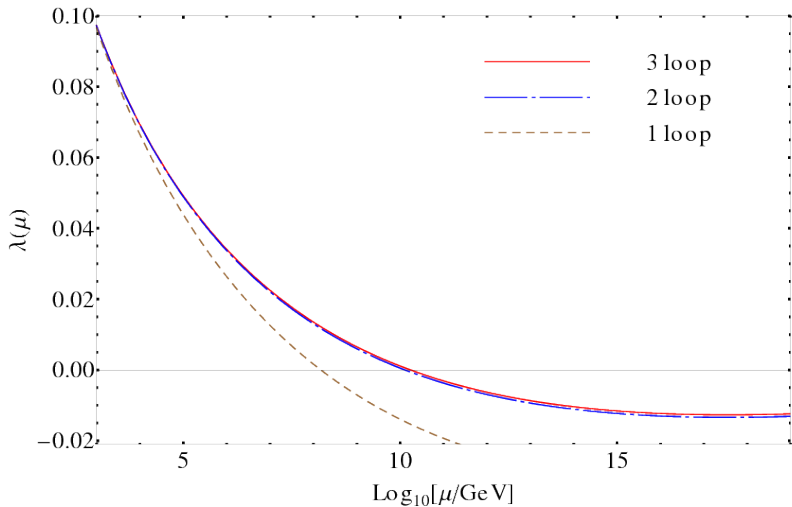
$$V_{\text{SM}} = -\mu^2 H^\dagger H + \lambda (H^\dagger H)^2$$

- large field values: $V \sim \lambda (H^\dagger H)^2$
- RGE: $\lambda \rightarrow \lambda(Q)$, where $Q \sim H$
- $\lambda \rightarrow 0$ around $Q \sim 10^{10}$ GeV, new minimum beyond M_{Planck}

Connection to Matter & Universe

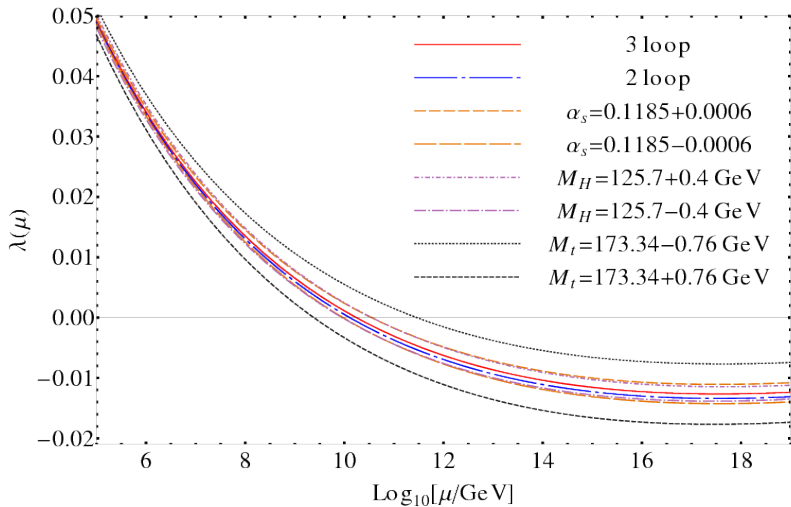
- transplanckian vev: different physics
- particle masses beyond M_{Planck} : link to gravitational physics
- Higgs inflation? [Bezrukov, Rubio, Shaposhnikov '14]
- main sources for uncertainty: m_t and α_S

Precise analysis: up to three loops!



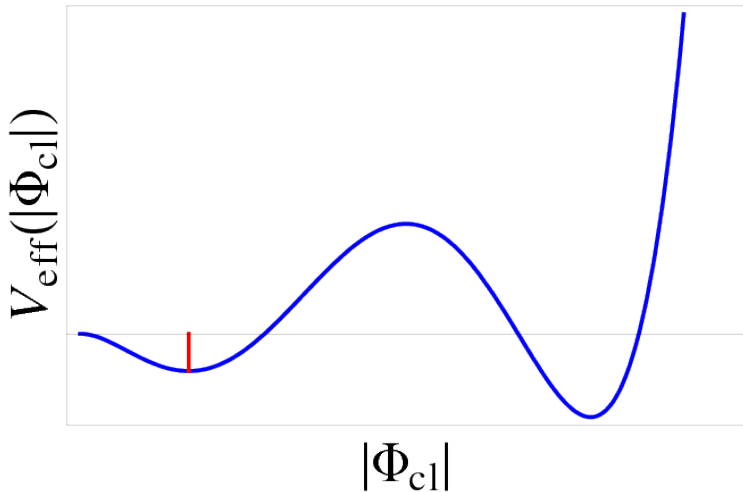
[Zoller 2014]

Precise analysis: up to three loops!



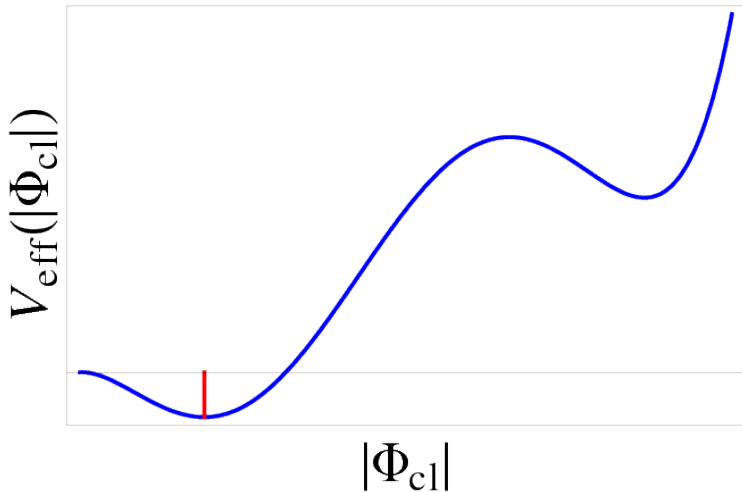
[Zoller 2014]

Stability, instability or metastability?



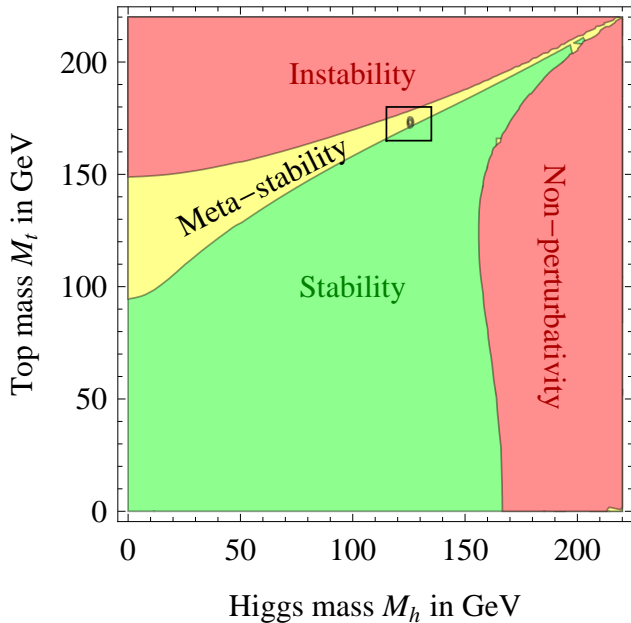
[Courtesy of Max Zoller]

Stability, instability or metastability?



[Courtesy of Max Zoller]

The SM phase diagram



[Degrassi et al. JHEP 1208 (2012) 098]

- Quantum effects: new particles in the loop
- tree level: multi-scalar potentials
 - Two Higgs Doublet Models (2HDM): no charge breaking
 - 2HDM + Singlet: more involved
- Minimal Supersymmetric Standard Model (MSSM):
 - sfermions: additional electrically and color charged directions
 - NMSSM: additional non-trivial neutral vacua
- severe constraints for any model, can be tested numerically
[cf. Vevacious]
- generically difficult to get practical limits
 - i.e. noanalytical bounds
 - if so, with many simplifications (see next slides)

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$$\begin{aligned}
 V_{\tilde{q},h} = & \tilde{t}_L^* (\tilde{m}_L^2 + |y_t h_2|^2) \tilde{t}_L + \tilde{t}_R^* (\tilde{m}_t^2 + |y_t h_2|^2) \tilde{t}_R \\
 & + \tilde{b}_L^* (\tilde{m}_L^2 + |y_b h_1|^2) \tilde{b}_L + \tilde{b}_R^* (\tilde{m}_b^2 + |y_b h_1|^2) \tilde{b}_R \\
 & - [\tilde{t}_L^* (\mu^* y_t h_1^* - A_t h_2) \tilde{t}_R + \text{h.c.}] \\
 & - [\tilde{b}_L^* (\mu^* y_b h_2^* - A_b h_1) \tilde{b}_R + \text{h.c.}] \\
 & + |y_t|^2 |\tilde{t}_L|^2 |\tilde{t}_R|^2 + |y_b|^2 |\tilde{b}_L|^2 |\tilde{b}_R|^2 + |y_t|^2 |\tilde{b}_L|^2 |\tilde{t}_R|^2 + |y_b|^2 |\tilde{t}_L|^2 |\tilde{b}_R|^2 \\
 & + \frac{g_1^2}{8} \left(|h_2|^2 - |h_1|^2 + \frac{1}{3} |\tilde{b}_L|^2 + \frac{2}{3} |\tilde{b}_R|^2 + \frac{1}{3} |\tilde{t}_L|^2 - \frac{4}{3} |\tilde{t}_R|^2 \right)^2 \\
 & + \frac{g_2^2}{8} \left(|h_2|^2 - |h_1|^2 + |\tilde{b}_L|^2 - |\tilde{t}_L|^2 \right)^2 + \frac{g_2^2}{2} |\tilde{t}_L|^2 |\tilde{b}_L|^2 \\
 & + \frac{g_3^2}{8} \left(|\tilde{t}_L|^2 - |\tilde{t}_R|^2 + |\tilde{b}_L|^2 - |\tilde{b}_R|^2 \right)^2 \\
 & + (m_{h_2}^2 + |\mu|^2) |h_2|^2 + (m_{h_1}^2 + |\mu|^2) |h_1|^2 - 2 \text{Re}(B_\mu h_1 h_2).
 \end{aligned}$$

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 \end{aligned}$$

$$|\tilde{t}_L| = |\tilde{t}_R| = |\tilde{t}|, \quad |\tilde{b}_L| = |\tilde{b}_R| = |\tilde{b}|$$

The MSSM tree-level scalar potential, 3rd generation squarks

$$\begin{aligned}
 V_{\tilde{q},h} = & \tilde{t}^* (\tilde{m}_L^2 + |y_t h_2|^2) \tilde{t} + \tilde{t}^* (\tilde{m}_t^2 + |y_t h_2|^2) \tilde{t} \\
 & + \tilde{b}^* (\tilde{m}_L^2 + |y_b h_1|^2) \tilde{b} + \tilde{b}^* (\tilde{m}_b^2 + |y_b h_1|^2) \tilde{b} \\
 & - [\tilde{t}^* (\mu^* y_t h_1^* - A_t h_2) \tilde{t} + \text{h.c.}] \\
 & - [\tilde{b}^* (\mu^* y_b h_2^* - A_b h_1) \tilde{b} + \text{h.c.}] \\
 & + |y_t|^2 |\tilde{t}_L|^2 |\tilde{t}_R|^2 + |y_b|^2 |\tilde{b}_L|^2 |\tilde{b}_R|^2 + |y_t|^2 |\tilde{b}_L|^2 |\tilde{t}_L|^2 + |y_b|^2 |\tilde{t}_L|^2 |\tilde{b}_L|^2 \\
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$$|\tilde{t}_L| = |\tilde{t}_R| = |\tilde{t}|, |\tilde{b}_L| = |\tilde{b}_R| = |\tilde{b}|; |\tilde{b}| = |h_1| = |\phi_1|, |\tilde{t}| = |h_2| = |\phi_2|$$

The MSSM tree-level scalar potential, 3rd generation squarks

$$\begin{aligned}
 V_{\tilde{q},h} = & \phi_2^* (\tilde{m}_L^2 + |y_t \phi_2|^2) \phi_2 + \phi_2^* (\tilde{m}_t^2 + |y_t \phi_2|^2) \phi_2 \\
 & + \phi_1^* (\tilde{m}_L^2 + |y_b \phi_1|^2) \phi_1 + \phi_1^* (\tilde{m}_b^2 + |y_b \phi_1|^2) \phi_1 \\
 & - [\phi_2^* (\mu^* y_t \phi_1^* - A_t \phi_2) \phi_2 + \text{h.c.}] \\
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 & + |y_t|^2 |\phi_2|^2 |\phi_2|^2 + |y_b|^2 |\phi_1|^2 |\phi_1|^2 + |y_t|^2 |\tilde{b}^-|^2 |\tilde{t}^-|^2 + |y_b|^2 |\tilde{t}^-|^2 |\tilde{b}^-|^2 \\
 & + \frac{g_2^2}{2} |\tilde{t}^-|^2 |\tilde{b}^-|^2 \\
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 \end{aligned}$$

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$$- [\phi_2^* (\quad - A_t \phi_2) \phi_2 + \text{h.c.}]$$

$$+ |y_t|^2 |\phi_2|^2 |\phi_2|^2$$

$$+ \cancel{|y_t|^2 |\tilde{b}|^2 |\tilde{t}|^2} + \cancel{|y_b|^2 |\tilde{t}|^2 |\tilde{b}|^2}$$

$$+ \cancel{\frac{g_2^2}{2} |\tilde{t}|^2 |\tilde{b}|^2}$$

$$+ (m_{h_2}^2 + |\mu|^2) |\phi_2|^2$$

$$|\tilde{t}_L| = |\tilde{t}_R| = |\tilde{t}|, |\tilde{b}_L| = |\tilde{b}_R| = |\tilde{b}|; \cancel{|\tilde{b}| = |h_1| = |\phi_1|}, |\tilde{t}| = |h_2| = |\phi_2|$$

Minimize the potential

$$V(\phi) = m^2\phi^2 - A\phi^3 + \lambda\phi^4,$$

with $m^2 = m_{h_2}^2 + |\mu|^2 + \tilde{m}_L^2 + \tilde{m}_t^2$, $A = -A_t$ and $\lambda = 3y_t^2$.

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Answer:

$$\phi_0 = 0, \quad \phi_{\pm} = \frac{3A \pm \sqrt{9A^2 - 32\lambda m^2}}{8\lambda}.$$

Condition to be safe from non-standard (i.e. non-trivial) minima:

$$V(\phi_{\pm}) > 0 \quad \Leftrightarrow \quad m^2 > \frac{A^2}{4\lambda}$$

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Well-known constraints

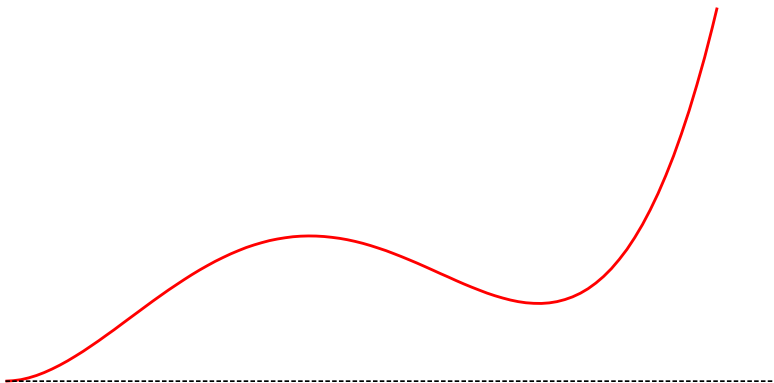
[Gunion, Haber, Sher '88]

$$|A_t|^2 < 3y_t^2 (m_{h_2}^2 + |\mu|^2 + \tilde{m}_L^2 + \tilde{m}_t^2)$$

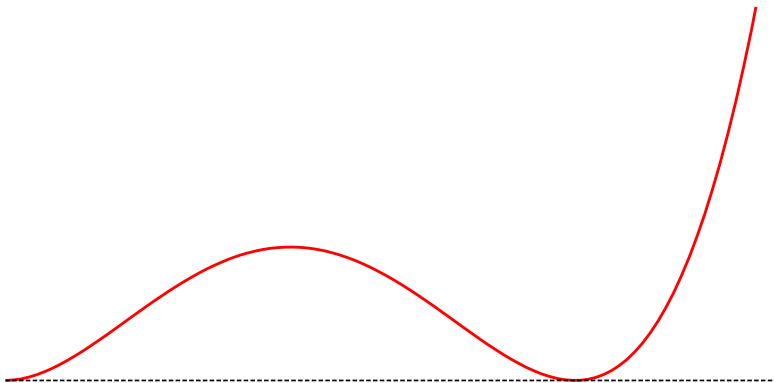
$$|A_b|^2 < 3y_b^2 (m_{h_1}^2 + |\mu|^2 + \tilde{m}_L^2 + \tilde{m}_b^2)$$

for the limiting cases $|\tilde{t}_L| = |\tilde{t}_R| = |h_2|$ and $|\tilde{b}_L| = |\tilde{b}_R| = |h_1|$!

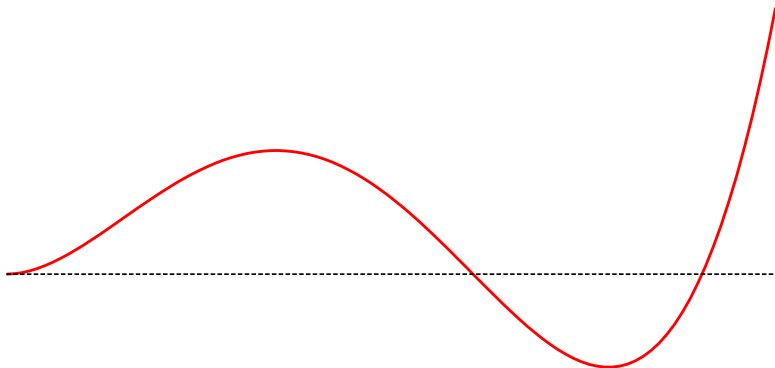
$$A^2 < 4\lambda m^2$$



$$A^2 = 4\lambda m^2$$



$$A^2 > 4\lambda m^2$$



A simple view of a complicated object

$$h_2 = \phi, \quad |\tilde{t}| = \alpha|\phi|, \quad h_1 = \eta\phi, \quad |\tilde{b}| = \beta|\phi|$$

$$\begin{aligned} V_\phi = & (m_{h_2}^2 + \eta^2 m_{h_1}^2 + (1 + \eta^2)\mu^2 - 2B_\mu\eta \\ & + (\alpha^2 + \beta^2)\tilde{m}_L^2 + \alpha^2\tilde{m}_t^2 + \beta^2\tilde{m}_b^2)\phi^2 \\ & - 2(\alpha^2(\mu y_t \eta - A_t) + \beta^2(\mu y_t - \eta A_b))\phi^3 + (\alpha^2 y_t^2 + \beta^4 y_b^2)\phi^4 \\ & + \left(\frac{g_1^2 + g_2^2}{8} (1 - \eta^2 + \beta^2 - \alpha^2)^2 + 2\alpha^2 y_t^2 + 2\beta^2 y_b^2 \right) \phi^4 \\ \equiv & M^2(\eta, \alpha, \beta)\phi^2 - \mathcal{A}(\eta, \alpha, \beta)\phi^3 + \lambda(\eta, \alpha, \beta)\phi^4, \end{aligned}$$

A simple view of a complicated object

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with

$$\begin{aligned} M^2 = & m_{h_2}^2 + \eta^2 m_{h_1}^2 + (1 + \eta^2)\mu^2 - 2B_\mu\eta \\ & + (\alpha^2 + \beta^2)\tilde{m}_L^2 + \alpha^2\tilde{m}_t^2 + \beta^2\tilde{m}_b^2, \end{aligned}$$

$$\mathcal{A} = 2\alpha^2\eta\mu y_t - 2\alpha^2 A_t + 2\beta^2\mu y_b - 2\eta\beta^2 A_b,$$

$$\begin{aligned} \lambda = & \frac{g_1^2 + g_2^2}{8}(1 - \eta^2 + \beta^2 - \alpha^2)^2 \\ & + (2 + \alpha^2)\alpha^2 y_t^2 + (2\eta^2 + \beta^2)\beta^2 y_b^2. \end{aligned}$$

[Gunion, Haber, Sher '88; Casas, Lleyda, Muñoz '96]

The same but different (“ A -parameter bounds”)

$$\mathcal{A}^2 < 4\lambda M^2$$

$$\downarrow$$

$$4\lambda(\eta, \alpha, \beta)M^2(\eta, \alpha, \beta) > (\mathcal{A}(\eta, \alpha, \beta))^2$$

$$h_u = \tilde{b}, h_d^0 = 0$$

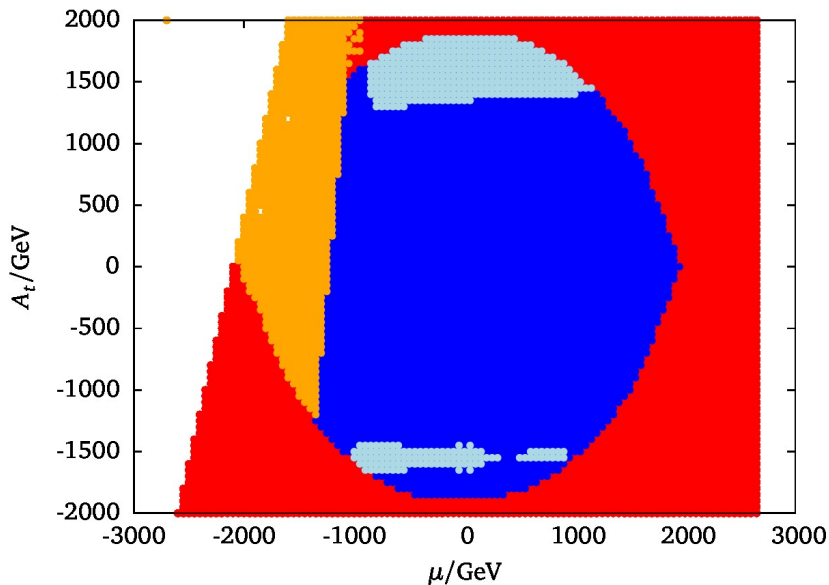
[WGH'15]

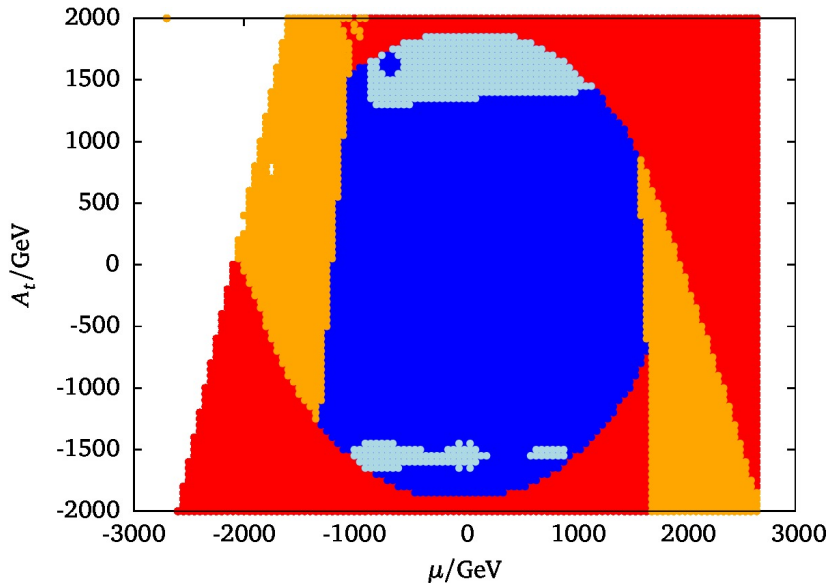
$$m_{H_u}^2 + \mu^2 + \tilde{m}_Q^2 + \tilde{m}_b^2 > \frac{(\mu y_b)^2}{y_b^2 + (g_1^2 + g_2^2)/2}$$

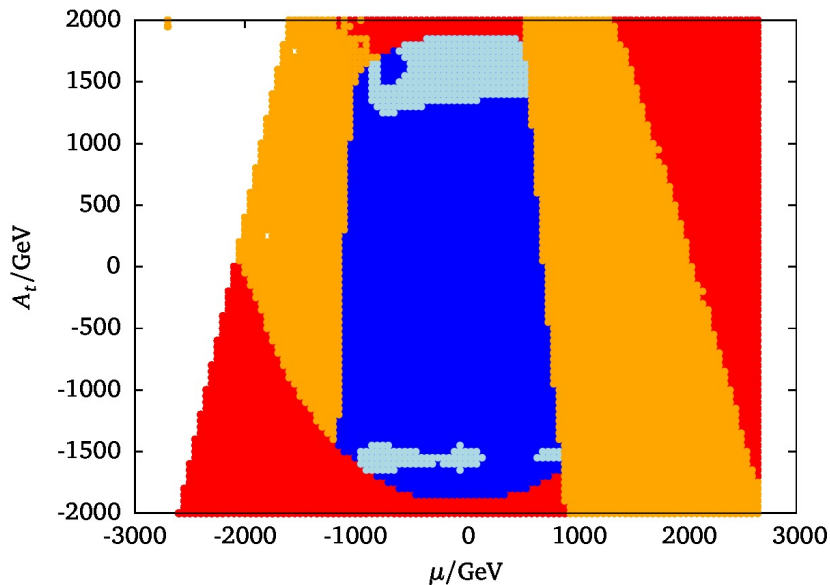
$$|h_d|^2 = |h_u|^2 + |\tilde{b}|^2, \tilde{b} = \alpha h_u$$

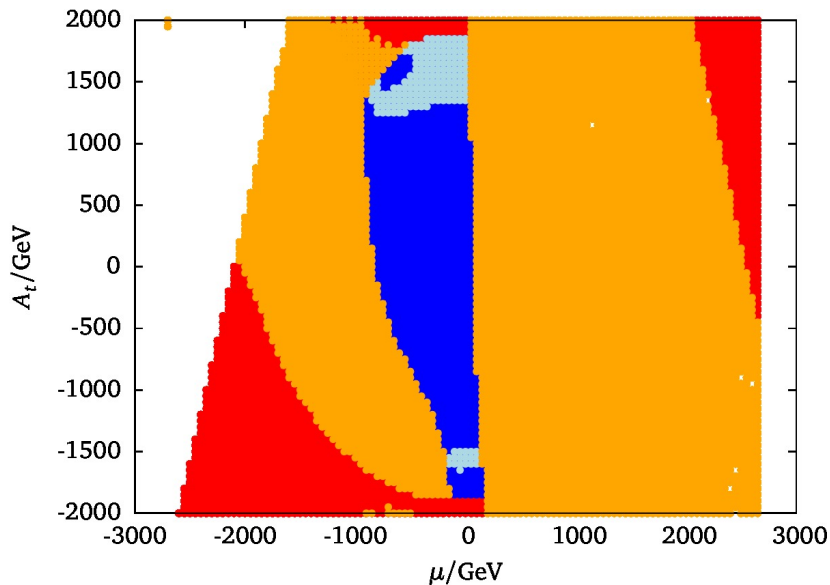
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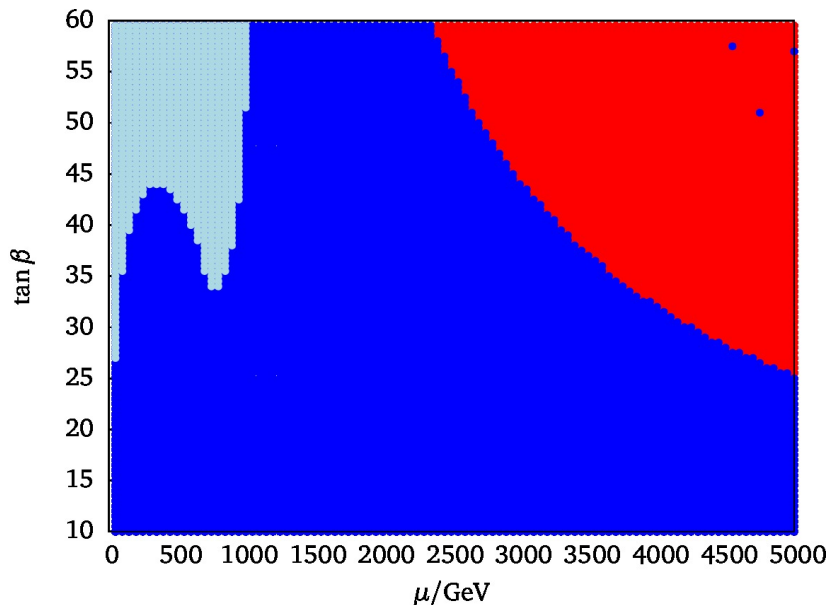
$$m_{11}^2(1 + \alpha^2) + m_{22}^2 \pm 2m_{12}^2\sqrt{1 + \alpha^2} + \alpha^2(\tilde{m}_Q^2 + \tilde{m}_b^2) > \frac{4\mu^2\alpha^2}{2 + 3\alpha^2}$$

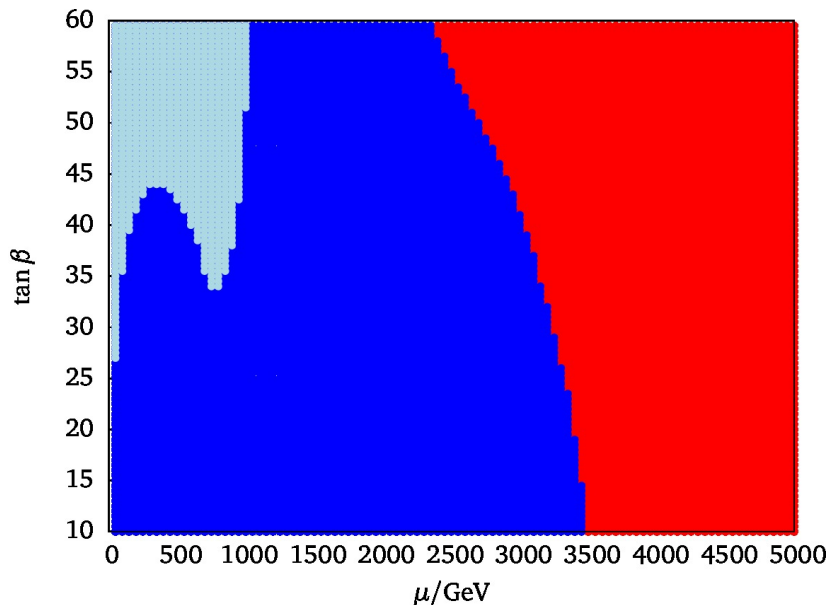


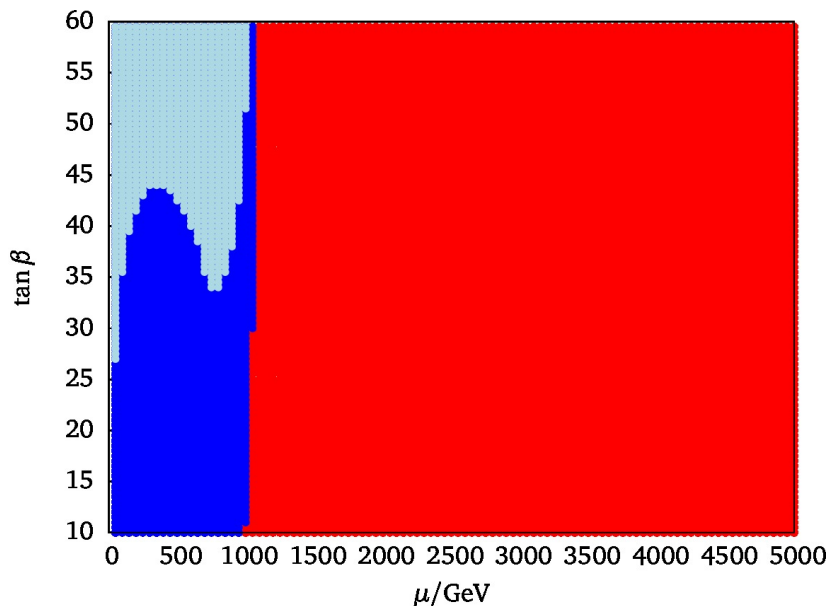


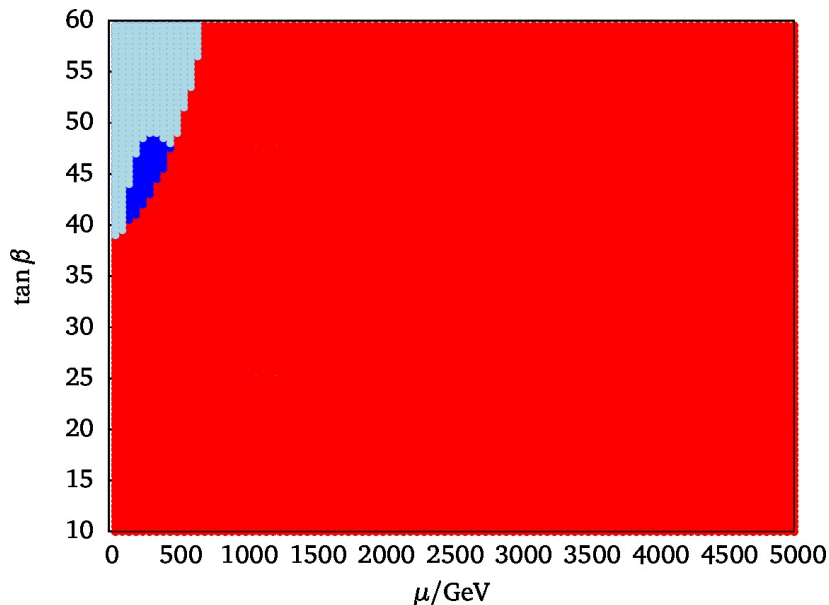




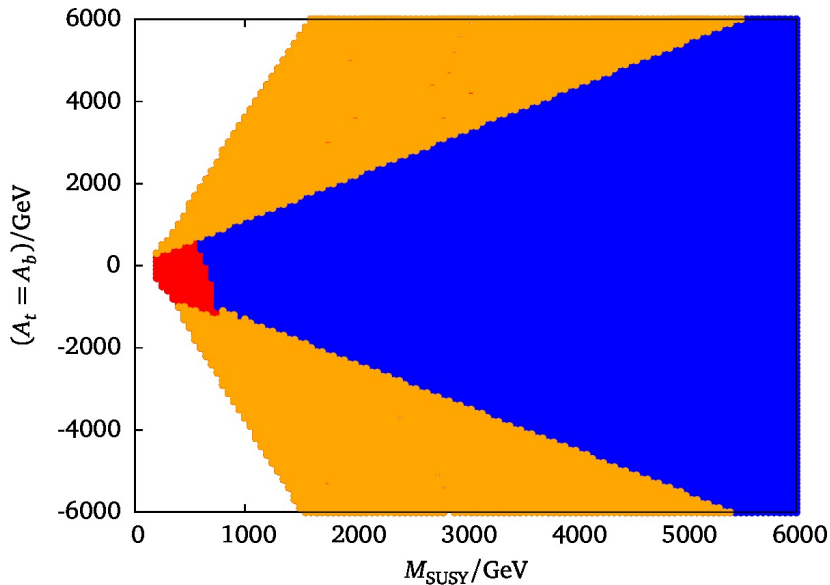


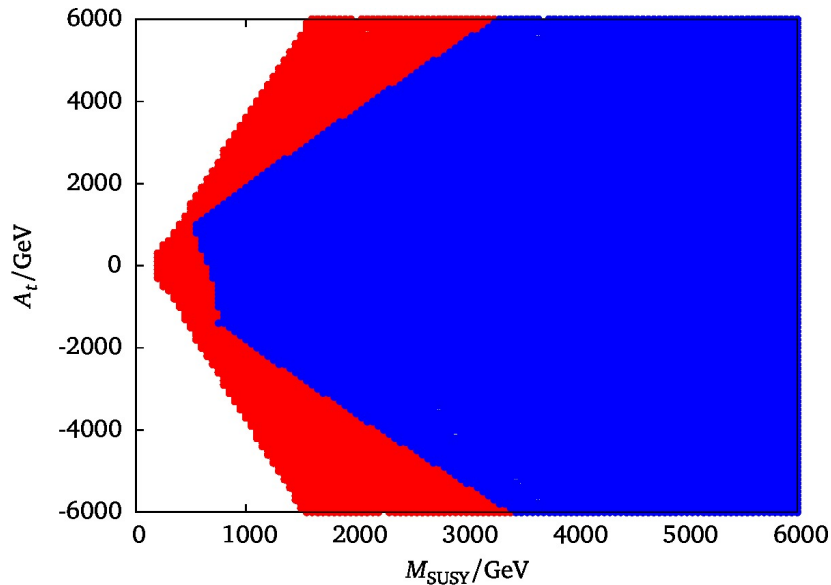


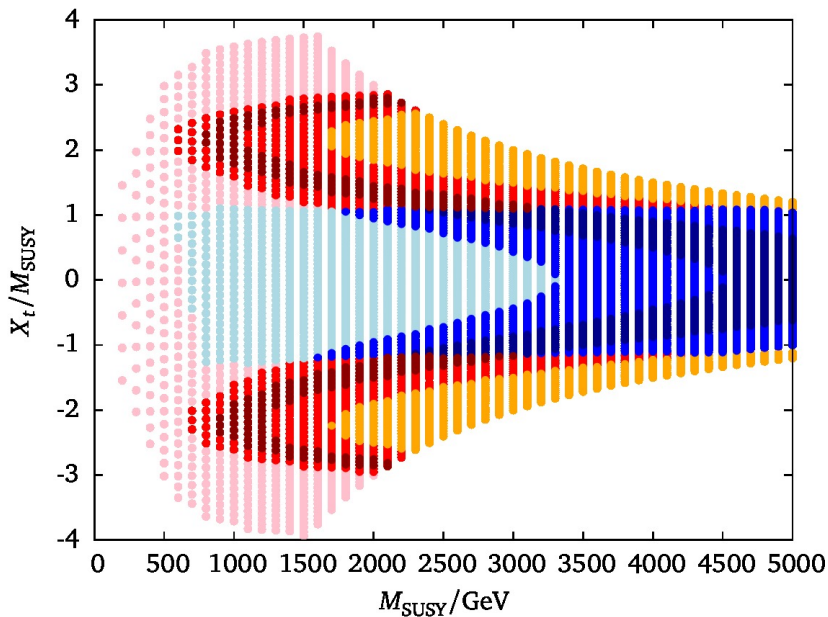


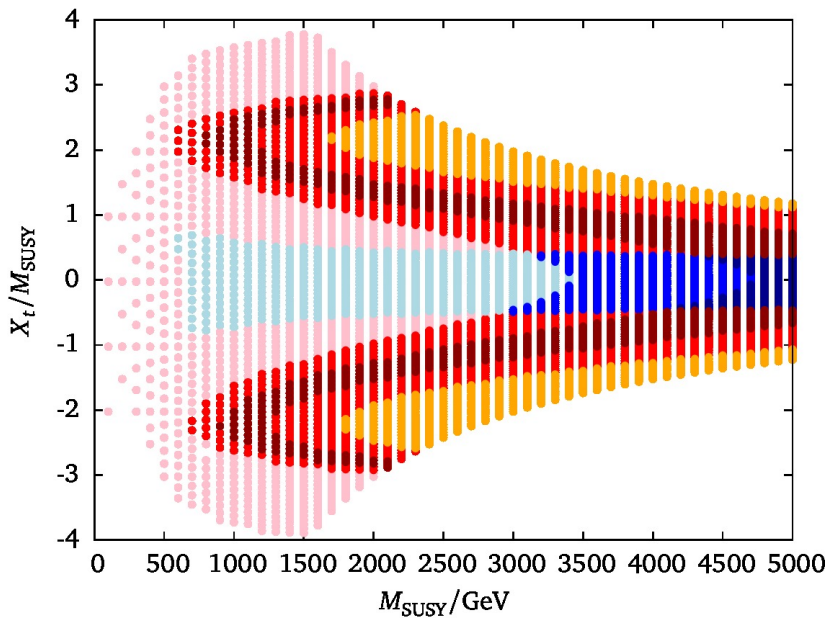


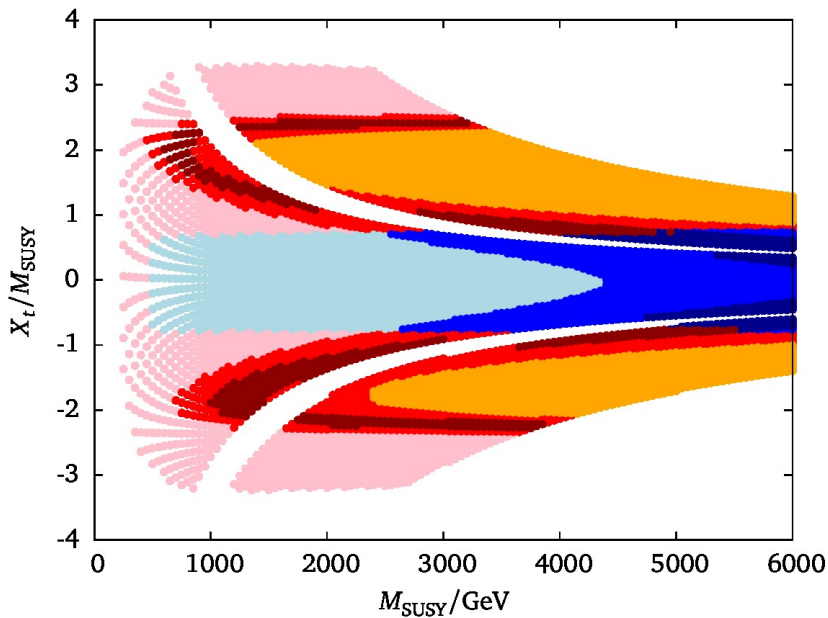
Resort comes close (increasing M_{SUSY})

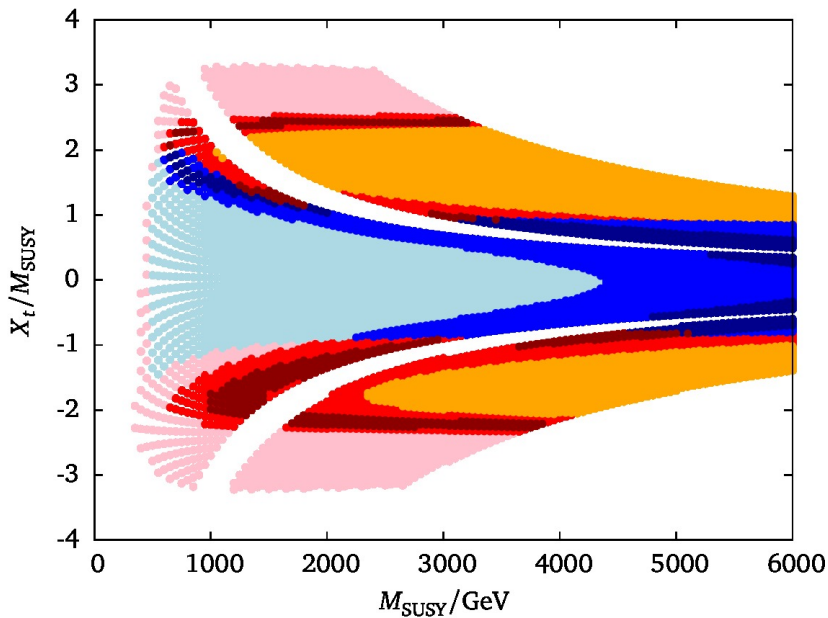












Short summary

- Vacuum stability and the origin of mass:
 $v = 246 \text{ GeV}$ vs. $v > M_{\text{Planck}}$ or $\tilde{v} \approx M_{\text{SUSY}}$
- constraints on model parameters from theoretical consistency:
global minimum has to be electroweak minimum
- “heavy” light Higgs @ 125 GeV: large SUSY corrections
 - e.g. MSSM: large $A_{t,b}$ and μ induce squark vevs



Backup

Slides

Cosmological stability

bounce action

$$B \gtrsim 400$$

↔ life-time longer than age of the universe

Decay probability (per unit volume)

$$\frac{\Gamma}{V} = Ae^{-B/\hbar}$$

[Coleman '77]

Death and doom

- value of B crucially depends on field space path
- very different conclusions for different η, α, β
- *independent* of SUSY parameter choice

