

The Radiative Return at Φ - and B -Meson Factories

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KARLSRUHER INSTITUT FÜR TECHNOLOGIE

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- IV Nucleon Form Factors at B-Factories
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and $\tau \rightarrow \nu K^- K^0$
- VI Experimental Results
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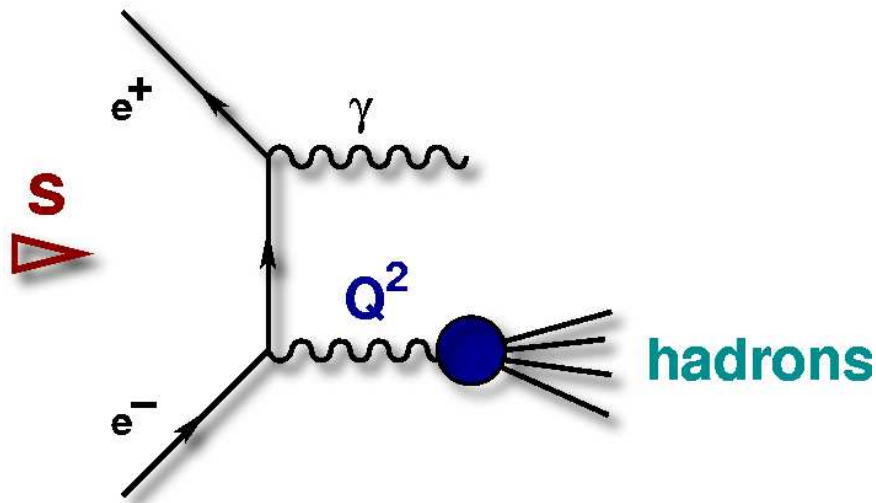
(with H. Czyż, G. Rodrigo, K. Melnikov;

and S. Binner, A. Grzelinska, E. Nowak, G. Rodrigo, A. Wapientnik)

I BASIC IDEA

photon radiated off the initial e^+e^- (ISR) reduces the effective energy of the collision

$$d\sigma(e^+e^- \rightarrow \text{hadrons} + \gamma) = H(Q^2, \theta_\gamma) d\sigma(e^+e^- \rightarrow \text{hadrons})$$



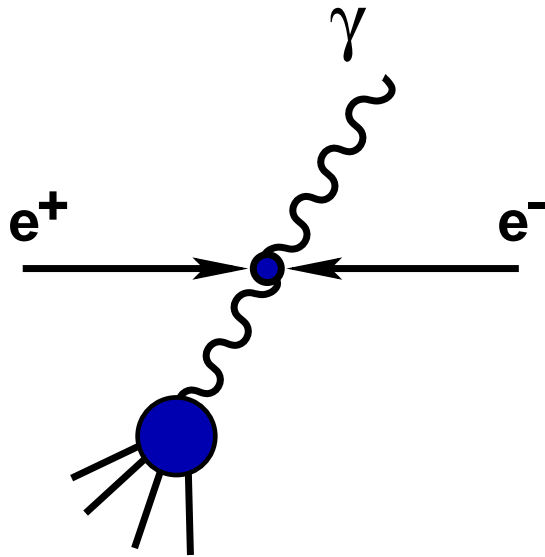
- ▶ measurement of $R(s)$ over the full range of energies, from threshold up to \sqrt{s}
- ▶ large luminosities of factories compensate α/π from photon radiation
- ▶ radiative corrections essential (NLO)
- ▶ advantage over energy scan (BES, CMD2, SND): systematics (e.g. normalization) only once

High precision measurement of the hadronic cross-section at DAΦNE, CLEO-C, B-factories

DAΦNE versus B-factories:

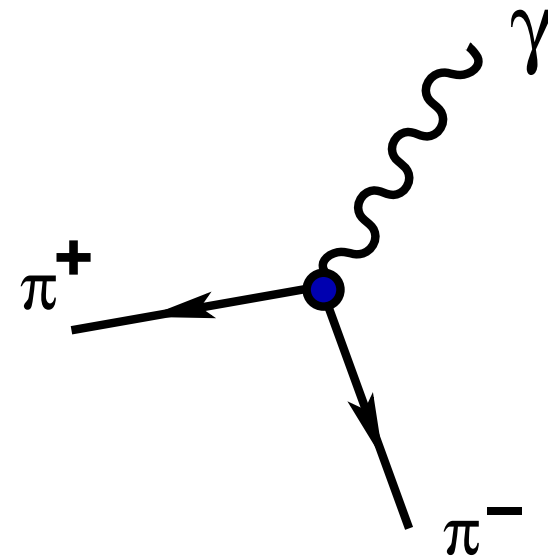
configurations in the cms - frame

10 GeV



very hard photon: clear kinematic separation between photon and hadrons

1 GeV



no natural kinematic separation
 \Rightarrow cuts to control FSR versus ISR

(two step process: $e^+e^- \rightarrow \gamma \rho(\rightarrow \gamma\pi\pi) \Rightarrow$ see below)

Rough estimates (1999) for rates:

$$\pi^+ \pi^- \gamma : E_\gamma > 100 \text{ MeV}$$

\sqrt{s} [GeV]	$\int \mathcal{L}$ [fb ⁻¹]	#events, $\theta_{min} = 7^\circ$
1.02	1.35	$16 \cdot 10^6$
10.6	100	$3.5 \cdot 10^6$

multi-hadron-events ($R \equiv 2$) $\sqrt{s} = 10.6 \text{ GeV}$

Q^2 -interval [GeV]	#events, $\theta_{min} = 7^\circ$
[1.5 , 2.0]	$9.9 \cdot 10^5$
[2.0 , 2.5]	$7.9 \cdot 10^5$
[2.5 , 3.0]	$6.6 \cdot 10^5$
[3.0 , 3.5]	$5.8 \cdot 10^5$

actually (2010): $\int \mathcal{L} \sim 3 \text{ fb}^{-1}$ (KLOE)

$\int \mathcal{L} \sim 500 \text{ fb}^{-1}$ (BABAR)

Lowest order

$$\frac{d\sigma}{dQ^2} (e^+e^- \rightarrow \gamma + \text{had}(Q^2)) = \sigma (e^+e^- \rightarrow \text{had}(Q^2))$$

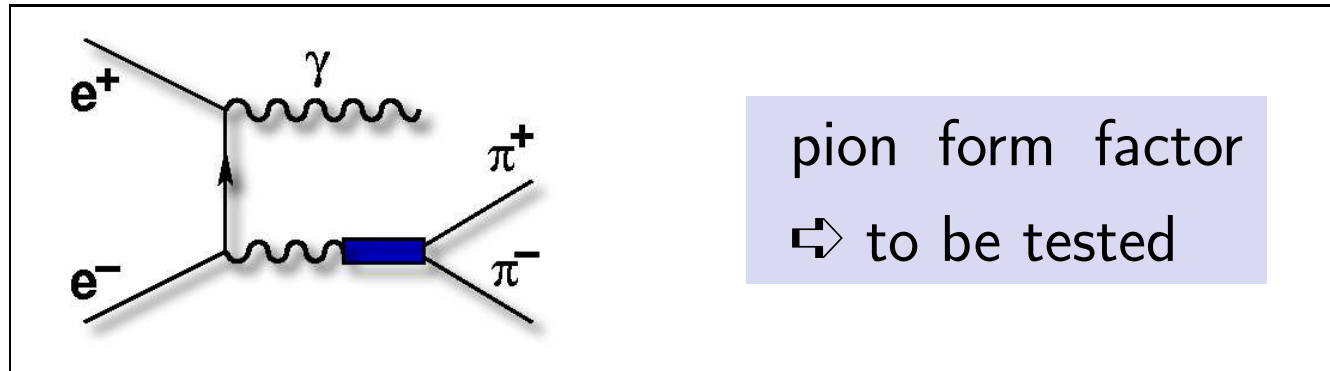
$$\times \frac{\alpha}{\pi s} \left\{ \begin{array}{l} \frac{s^2+Q^4}{s(s-Q^2)} (\log(s/m_e^2) - 1) , \text{ no angular cut} \\ \frac{s^2+Q^4}{s(s-Q^2)} \log \left(\frac{1+\cos \theta_{min}}{1-\cos \theta_{min}} \right) - \frac{s-Q^2}{s} \cos \theta_{min} \end{array} \right\}$$

$$\Rightarrow \text{differential luminosity: } \frac{dL}{dQ^2} (Q^2, s) = \frac{\alpha}{\pi s} \left\{ \dots \right\} L(\text{at } s)$$

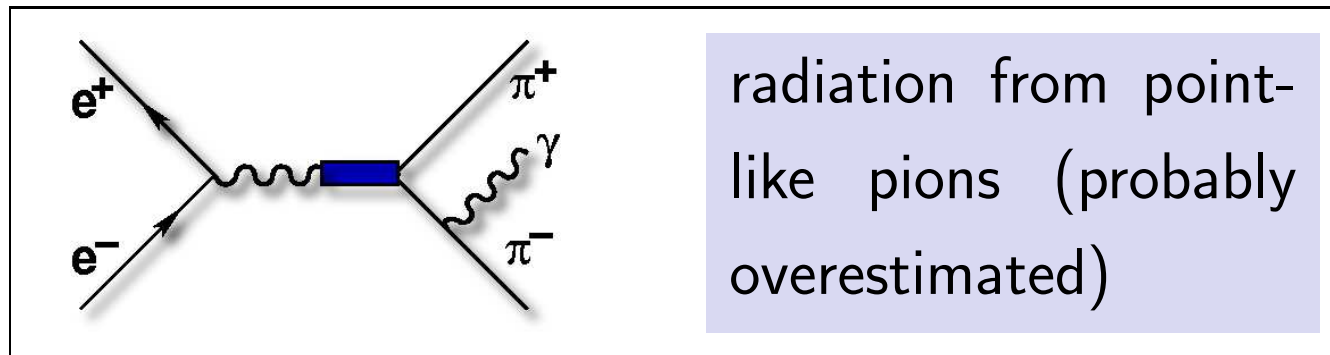
all errors dominated by systematics
⇒ Monte Carlo Generator

Basic Ingredients for Pion Formfactor

► ISR



► FSR



⇐ virtual Compton scattering

- **additional radiation:** collinear (EVA MC) (Binner, JK, Melnikov) or NLO calculation (PHOKHARA MC)

II MONTE CARLO GENERATORS



P
H
O
K
H
A
R
A

References etc. → <http://ific.uv.es/~rodrigo/phokhara>

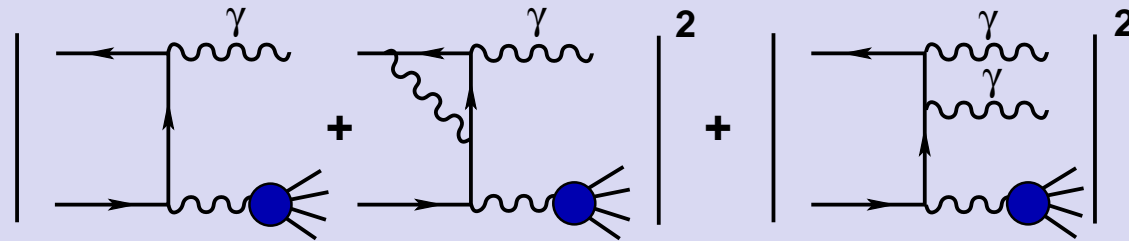
QED corrections at leptonic side

⇒ basic building block for all hadronic final states

PHOKHARA 2.0:

$$\pi^+\pi^-, \mu^+\mu^-, 4\pi$$

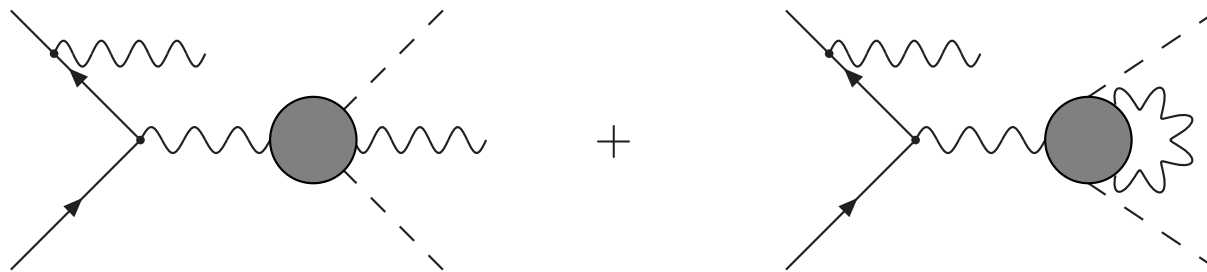
- **ISR at NLO:** virtual corrections to one photon events and two photon emission at tree level



- FSR at LO: $\pi^+\pi^-, \mu^+\mu^-$
- tagged or untagged photons
- modular structure

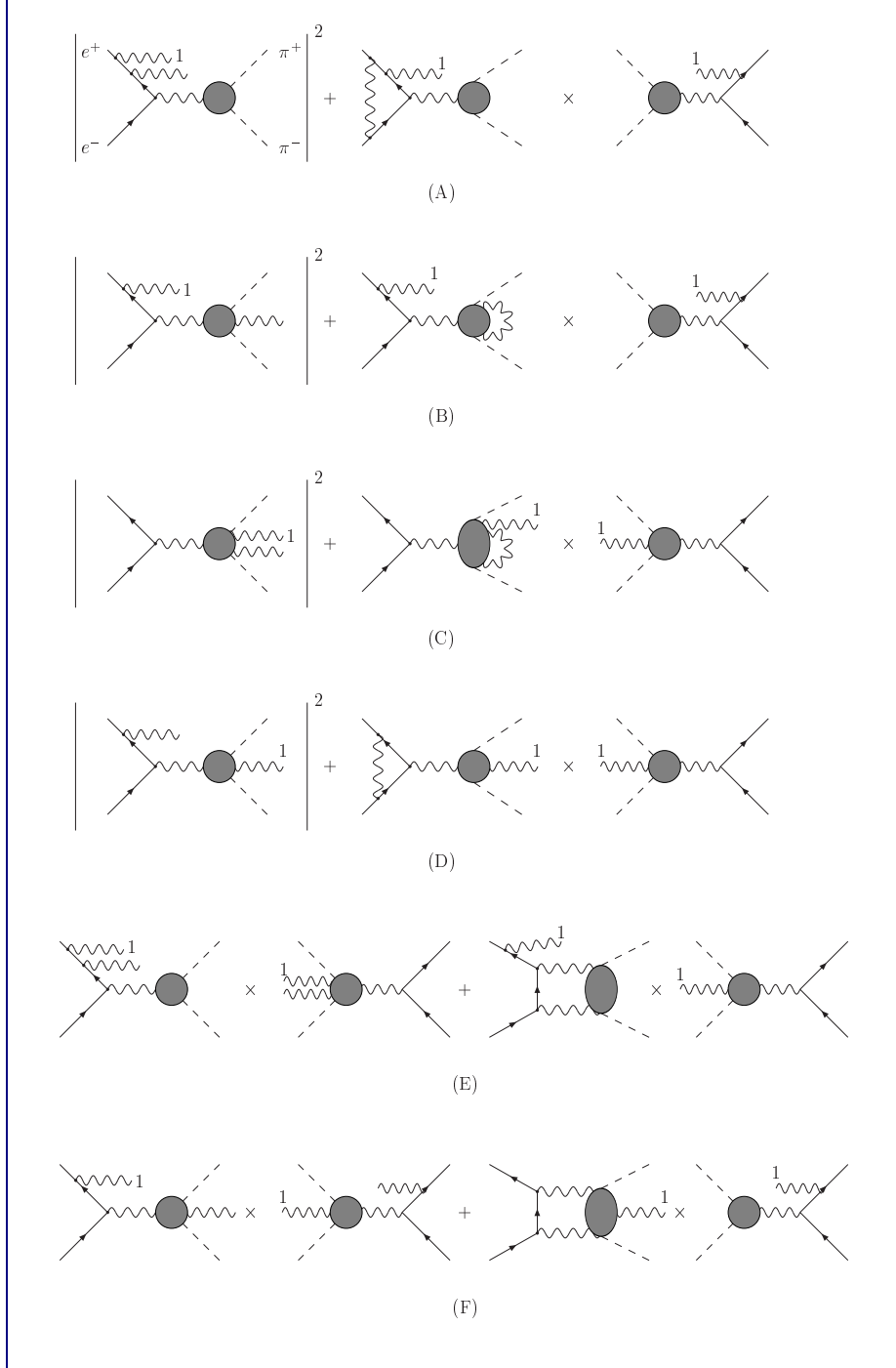
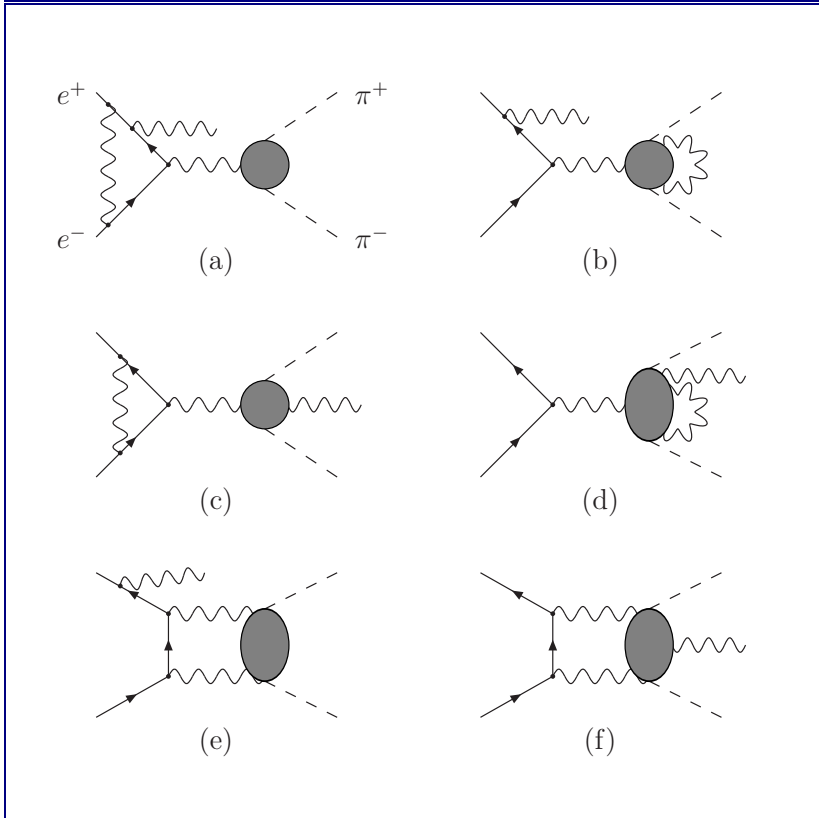
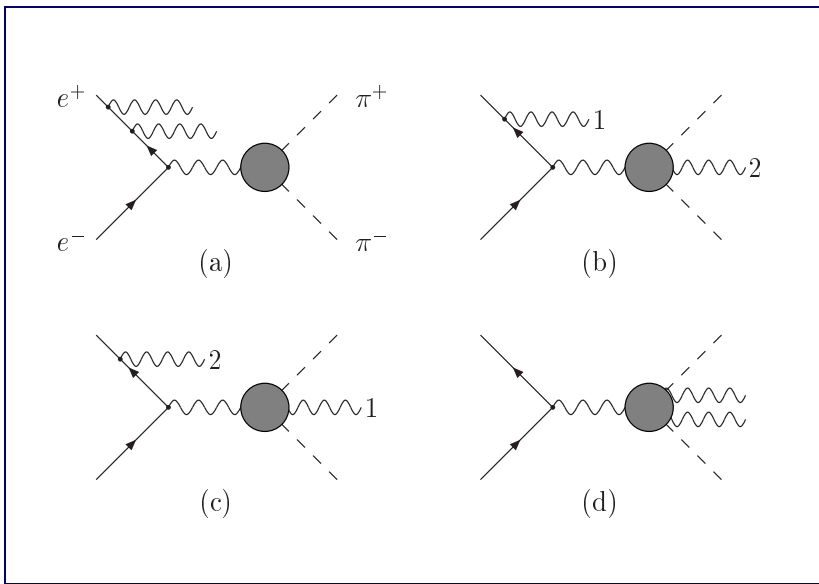
PHOKHARA 3.0

- ▶ specifically developed for $\pi^+\pi^-$ (plus photons)
- ▶ allows for **simultaneous** emission of photons from **initial and final state**, including virtual corrections (interference neglected).

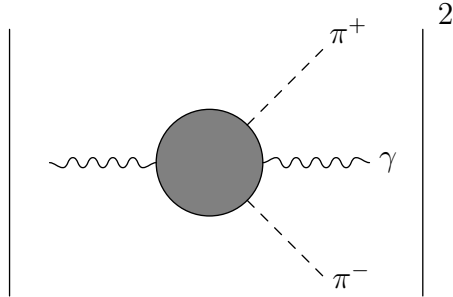


⇒ dominated by “two step process”: $e^+e^- \rightarrow \gamma \rho (\rightarrow \gamma \pi\pi)$

⇒ importance of $\pi\pi\gamma$ as input for a_μ



Contributions to a_μ from

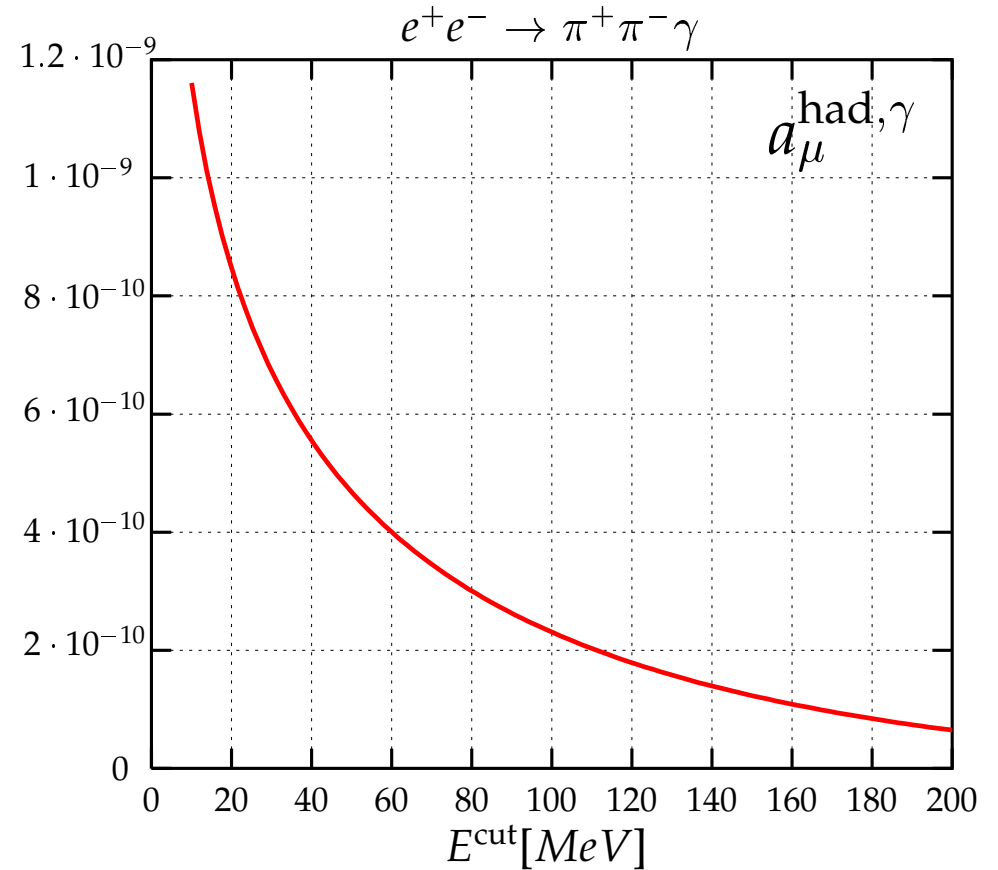
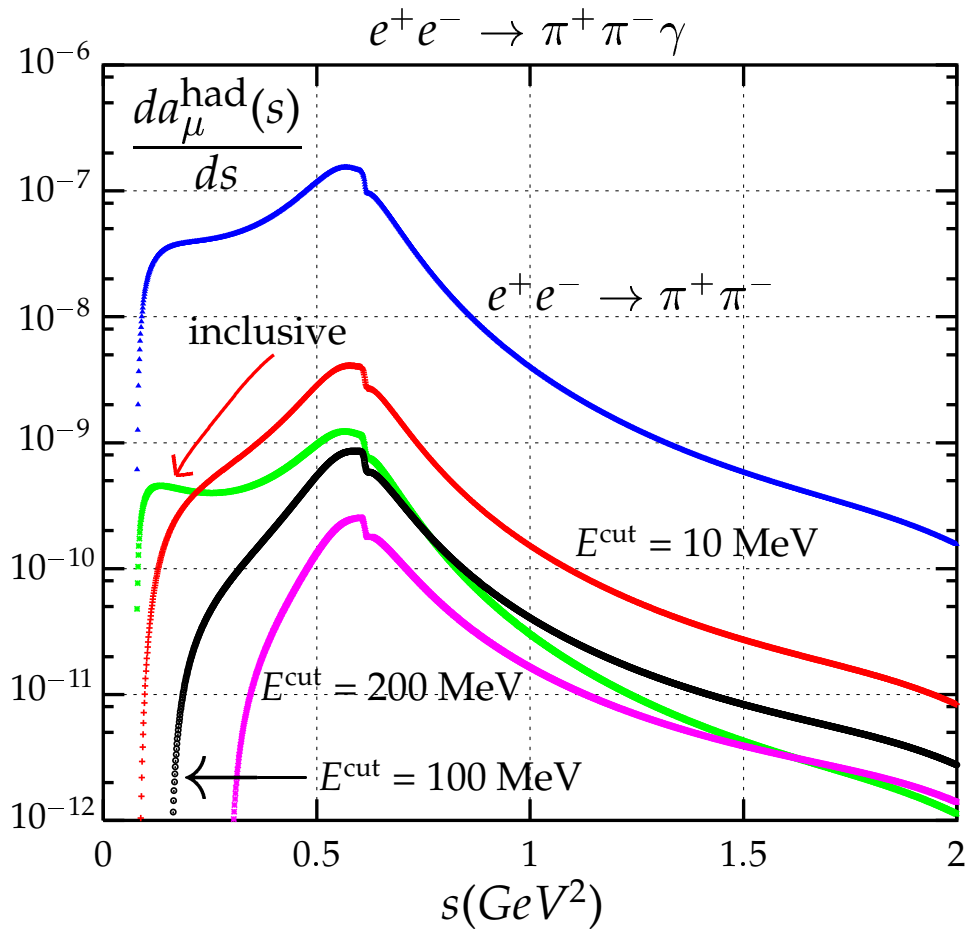


Estimates

$$\delta a_\mu(\text{quark}, \gamma, m_q = 180 \text{ MeV}) = 1.880 \times 10^{-10}$$

$$\delta a_\mu(\text{quark}, \gamma, m_q = 66 \text{ MeV}) = 8.577 \times 10^{-10}$$

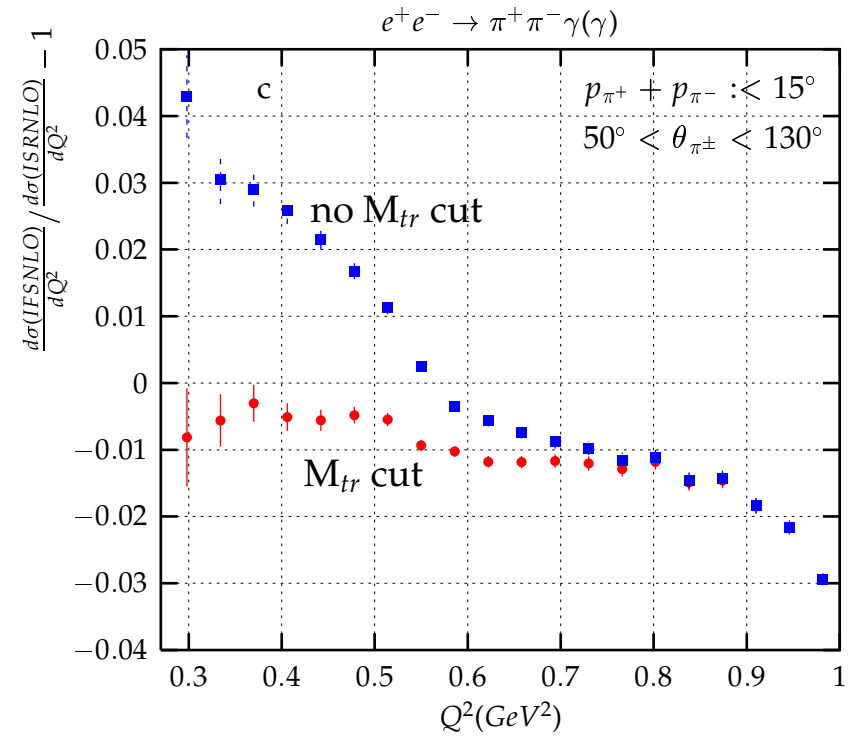
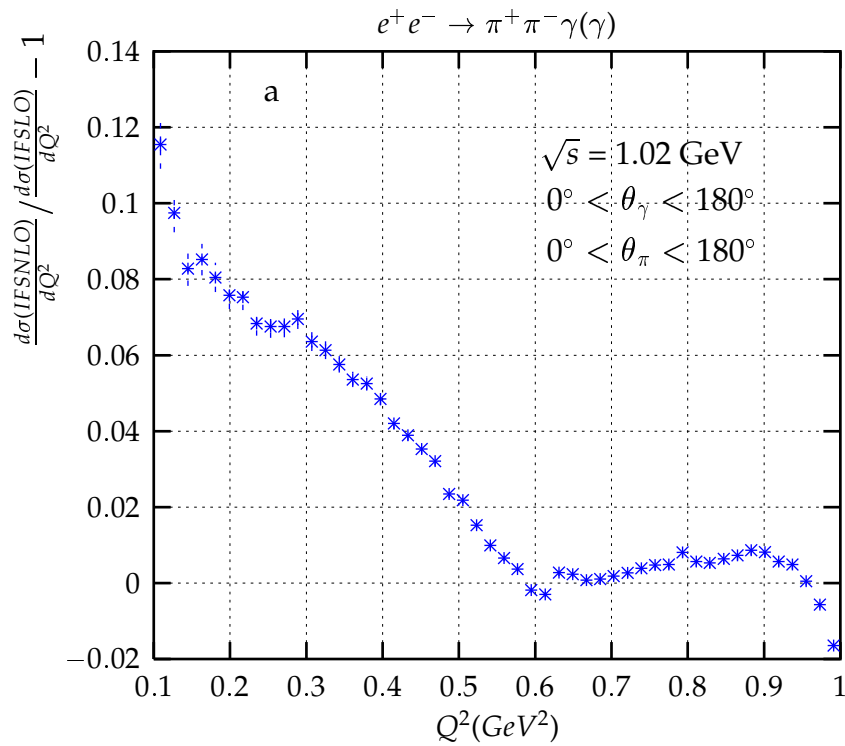
$$\delta a_\mu(\pi^+ \pi^-, \gamma) = 4.309 \times 10^{-10}$$



Differential contribution to $a_{\mu}^{\text{had},\gamma}$ from $\pi^+\pi^-\gamma$ intermediate states for different cutoff values compared with the complete contribution (virtual plus real corrections, labelled 'inclusive') evaluated in sQED (FSR), as well as with the contribution from the $\pi^+\pi^-$ intermediate state.

Integrated contribution to $a_{\mu}^{\text{had},\gamma}$ as a function of the cutoff E^{cut} .

Large effect for $Q^2 < m_\rho^2$ eliminated by suitable cuts on $\pi^+\pi^-$ configuration (suppress 2γ events)

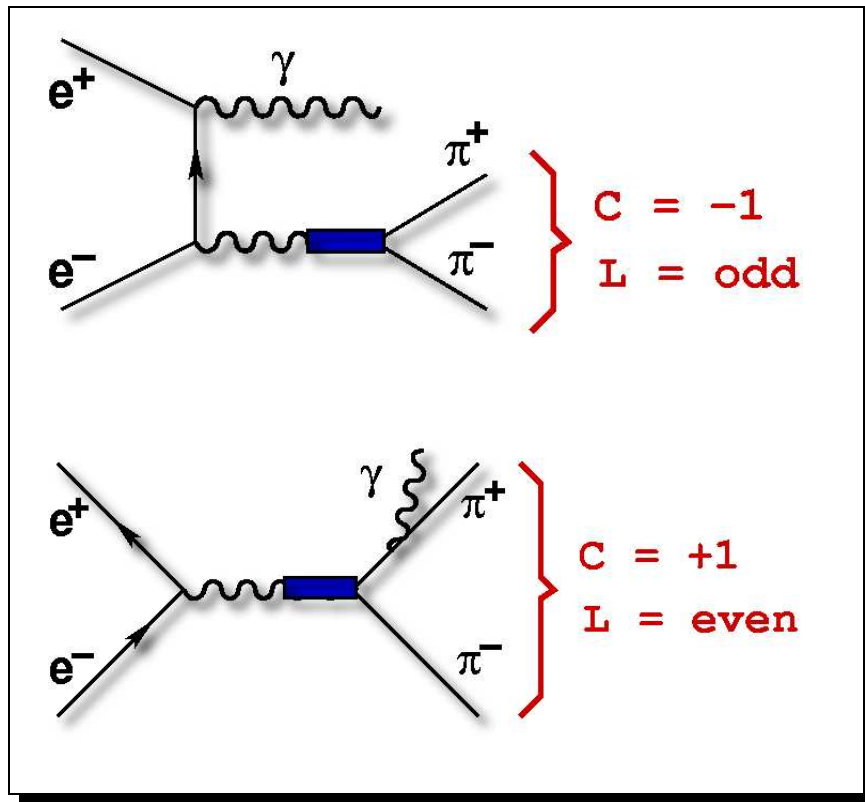


or measure photon

III Charge Asymmetries and Radiative Φ -Decays

(H. Czyż, A. Grzelinska, JK)

interference:



⇒ interference odd
under $\pi^+ \leftrightarrow \pi^-$

⇒ asymmetric differential
distribution: $\int \text{interf.} = 0$

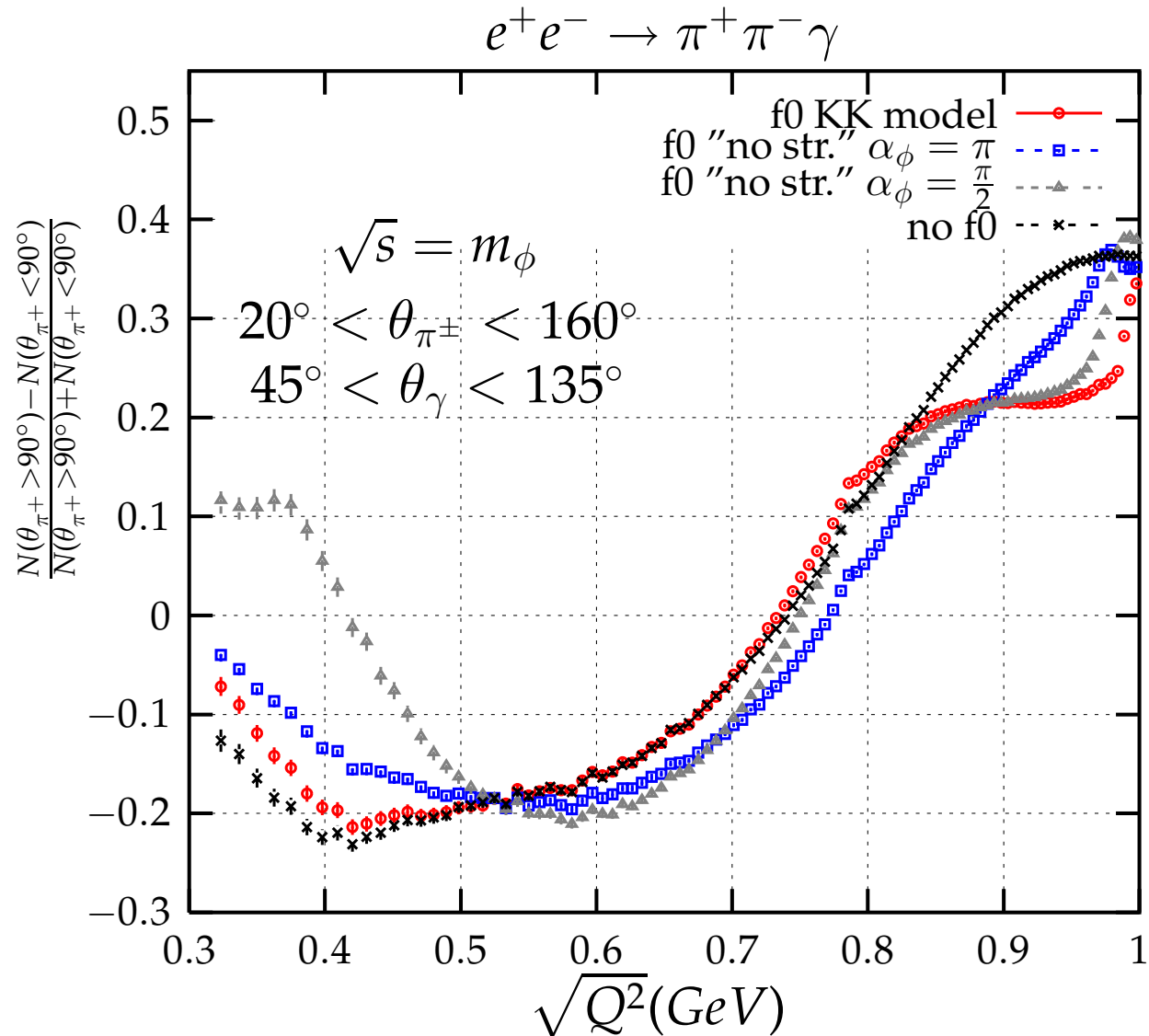
$$A(\theta) = \frac{N^{\pi^+}(\theta) - N^{\pi^-}(\theta)}{N^{\pi^+}(\theta) + N^{\pi^-}(\theta)}$$

Photon coupled to “pointlike” pion

additional contribution on top of Φ -resonance (KLOE !)

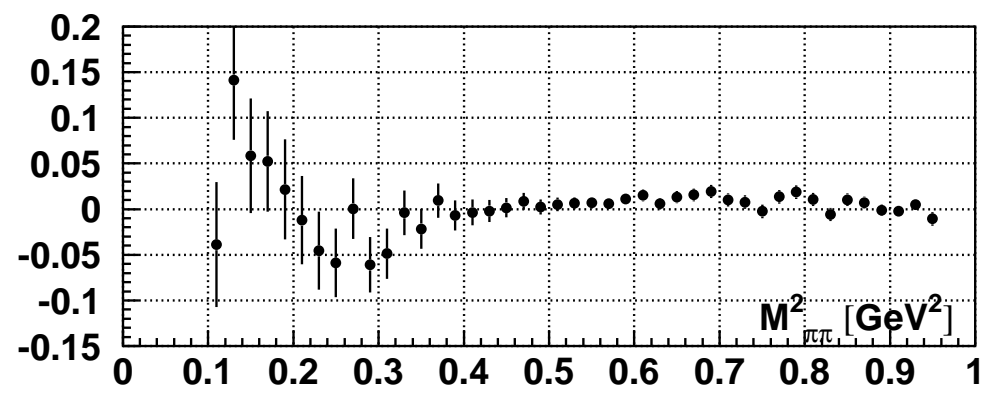
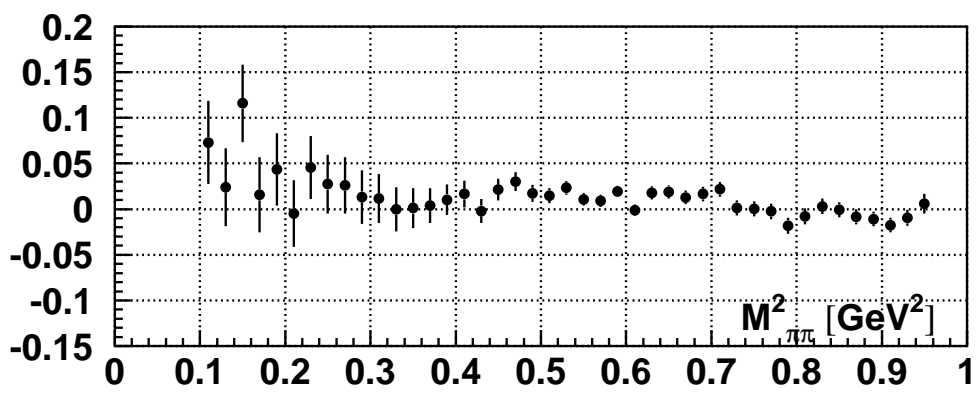
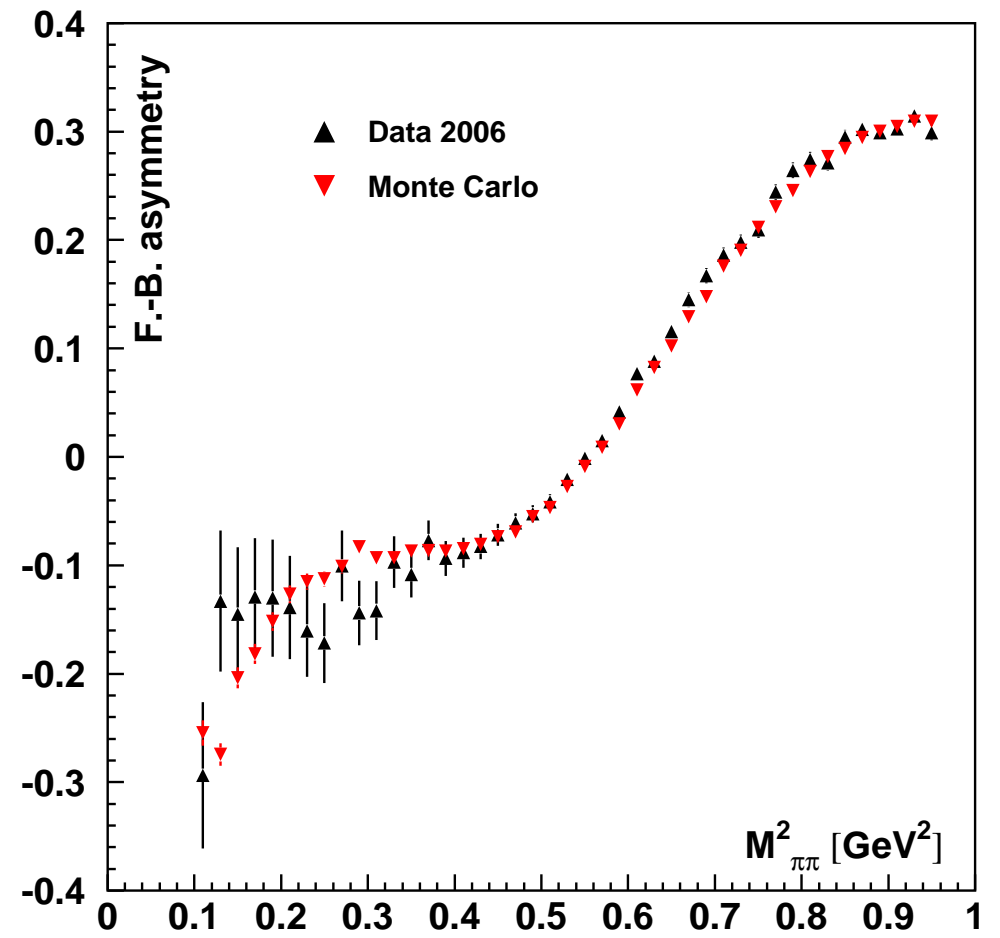
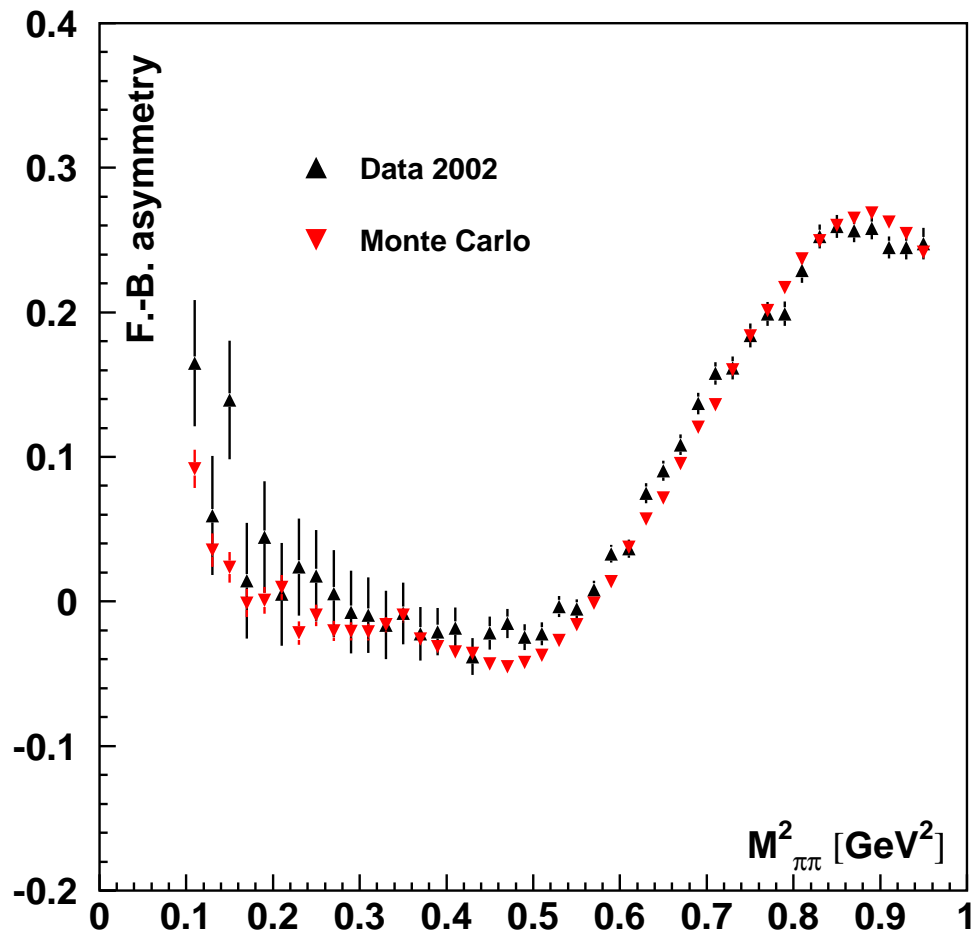
$e^+e^- \rightarrow \Phi \rightarrow \gamma f_{0,2} (\rightarrow \pi^+\pi^-)$ interference !

Significant influence of scalar resonances on charge asymmetry



Interference between ISR and FSR

\Rightarrow amplitude for $\Phi \rightarrow \gamma\pi\pi$



IV NUCLEON FORM FACTORS

(PHOKHARA 4.0)

$Q^2 \gtrsim 4m_N^2$ accessible at B-factories
 \Rightarrow study $e^+e^- \rightarrow \gamma N\bar{N}$ (with $N = p$ or n)

hadronic current:

$$J_\mu = -ie \cdot \bar{u}(q_2) \left(F_1^N(Q^2) \gamma_\mu - \frac{F_2^N(Q^2)}{4m_N} [\gamma_\mu, \not{Q}] \right) v(q_1),$$

$$Q = q_1 + q_2, \quad q = (q_1 - q_2)/2$$

or

$$G_M = F_1 + F_2, \quad G_E = F_1 + \frac{Q^2}{4m^2} F_2$$

Result:

$$d\sigma = \frac{1}{2s} L_{\mu\nu} H^{\mu\nu} d\Phi_2(p_1 + p_2; Q, k) d\Phi_2(Q; q_1, q_2) \frac{dQ^2}{2\pi},$$

$$L_{\mu\nu} H^{\mu\nu} = \frac{(4\pi\alpha)^3}{Q^2} \left\{ \left(|G_M^N|^2 - \frac{1}{\tau} |G_E^N|^2 \right) \right. \\ \times \frac{32s}{\beta_N^2 (s - Q^2)} \left(\frac{1}{y_1} + \frac{1}{y_2} \right) \left(\frac{(p_1 \cdot q)^2 + (p_2 \cdot q)^2}{s^2} \right) \\ \left. + 2 \left(|G_M^N|^2 + \frac{1}{\tau} |G_E^N|^2 \right) \left[\left(\frac{1}{y_1} + \frac{1}{y_2} \right) \frac{(s^2 + Q^4)}{s(s - Q^2)} - 2 \right] \right\},$$

where

$$y_{1,2} = \frac{s - Q^2}{2s} (1 \mp \cos \theta_\gamma), \quad \tau = \frac{Q^2}{4m_N^2}, \quad \beta_N^2 = 1 - \frac{4m_N^2}{Q^2}$$

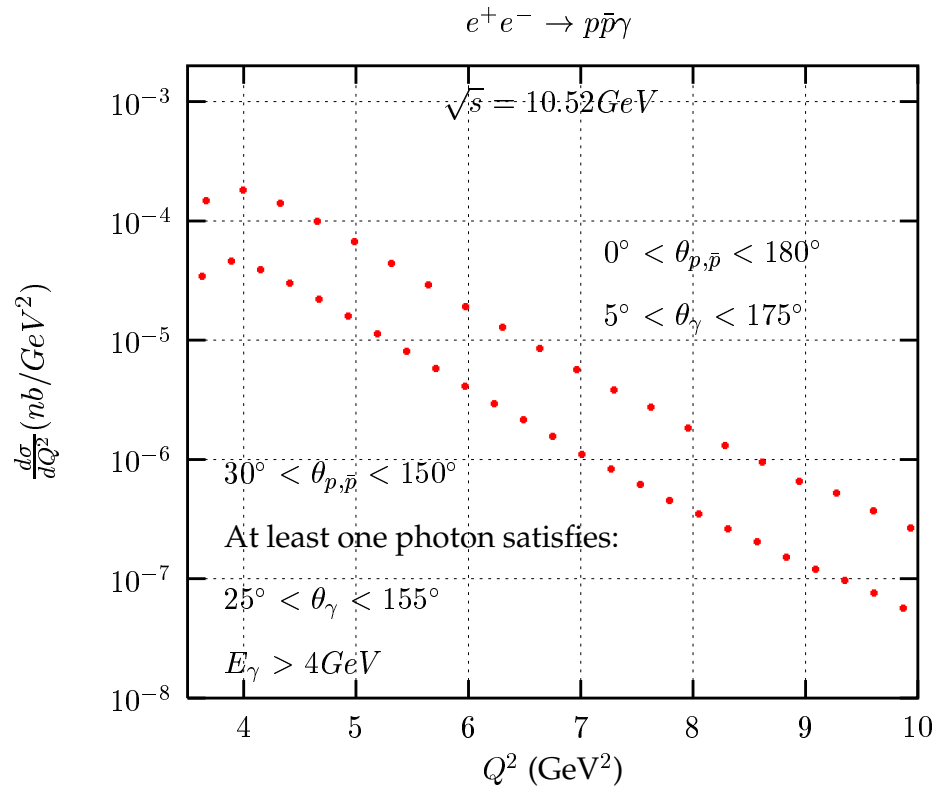
Separation of $|G_M|^2$ and $|G_E|^2$ through angular distribution:

$$L_{\mu\nu}H^{\mu\nu} = \frac{(4\pi\alpha)^3 (1 + \cos^2 \theta_\gamma)}{Q^2 (1 - \cos^2 \theta_\gamma)} \times 4 \left(|G_M^N|^2 (1 + \cos^2 \hat{\theta}) + \frac{1}{\tau} |G_E^N|^2 \sin^2 \hat{\theta} \right)$$

$\hat{\theta}$ = angle of nucleon with respect to γ -direction in hadronic rest frame
 (valid for $s/Q^2 \ll 1$, corrections and “optimal frame” → EPJ C 35 (2004) 527)

Similarity to $e^+e^- \rightarrow N\bar{N}$:

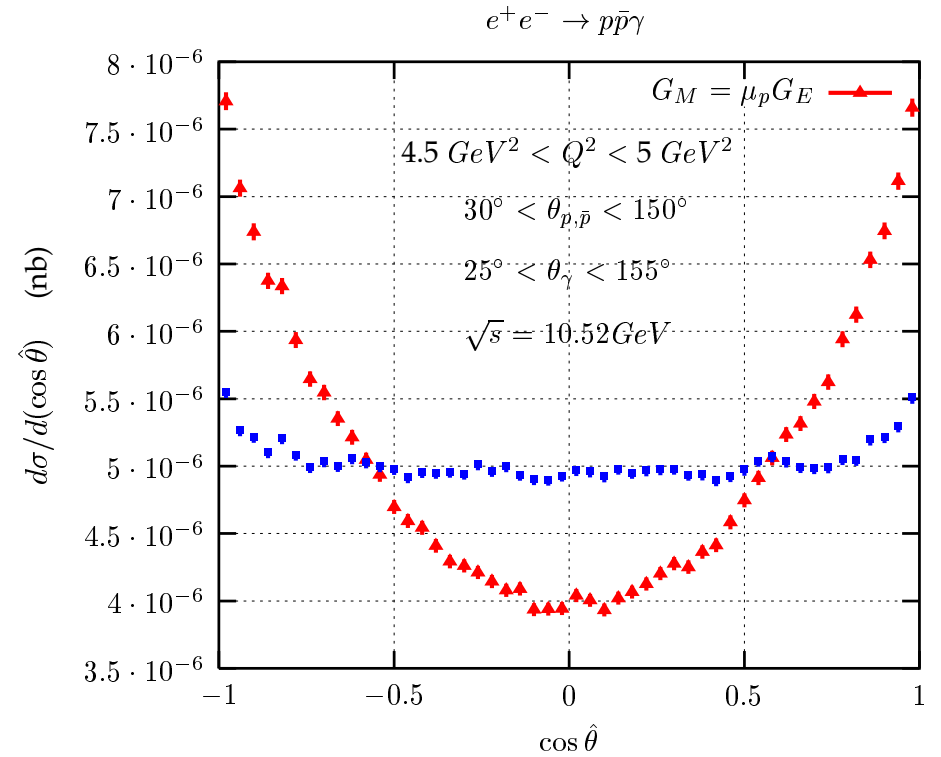
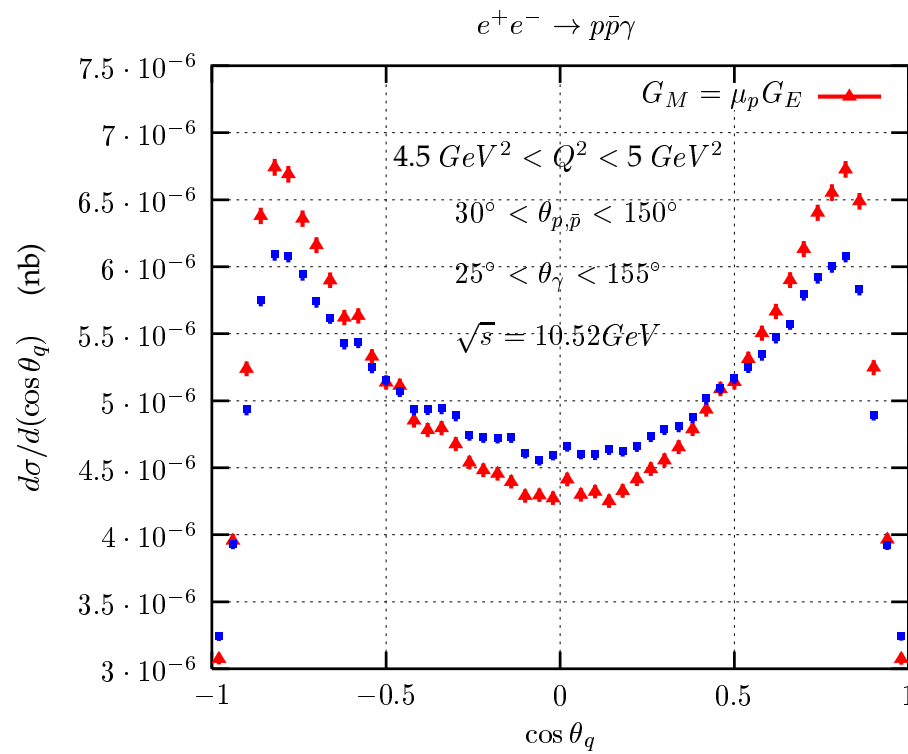
$$\frac{d\sigma}{d\Omega} = \frac{\alpha^2\beta_N}{4Q^2} \left(|G_M^N|^2 (1 + \cos^2 \theta) + \frac{1}{\tau} |G_E^N|^2 \sin^2 \theta \right)$$



$$e^+e^- \rightarrow p\bar{p}\gamma$$

implementation in PHOKHARA

Angular distributions of nucleon



(two choices for G_M/G_E)

Comments

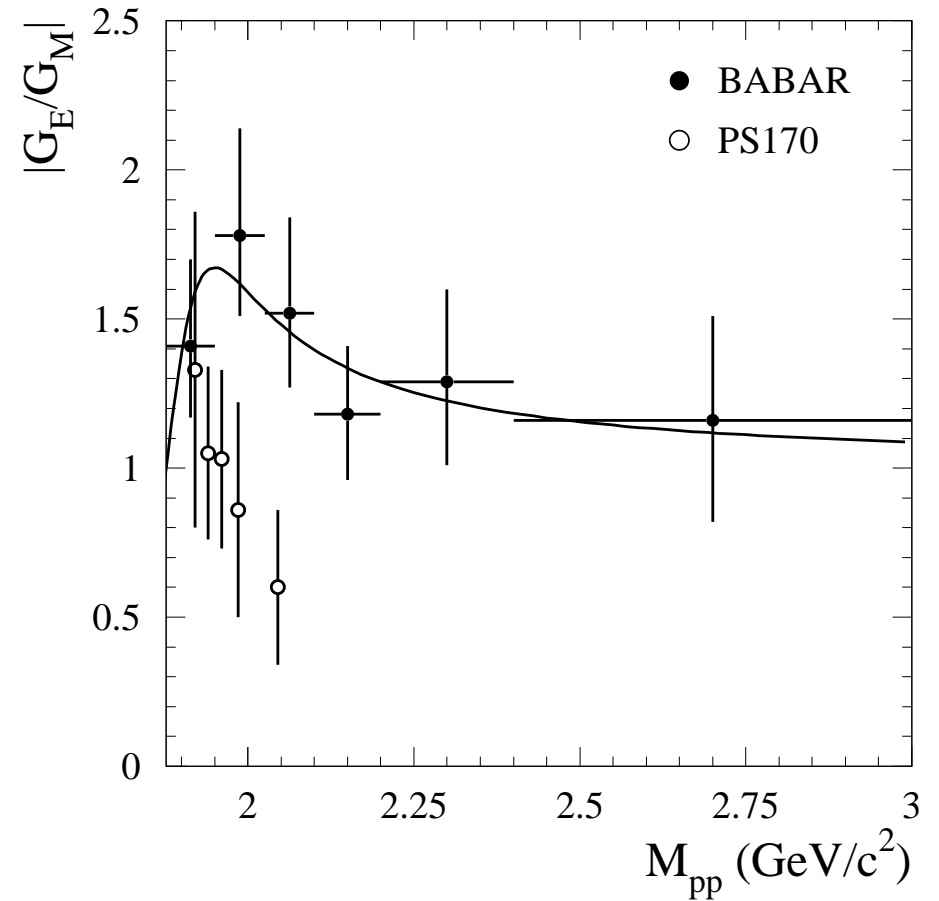
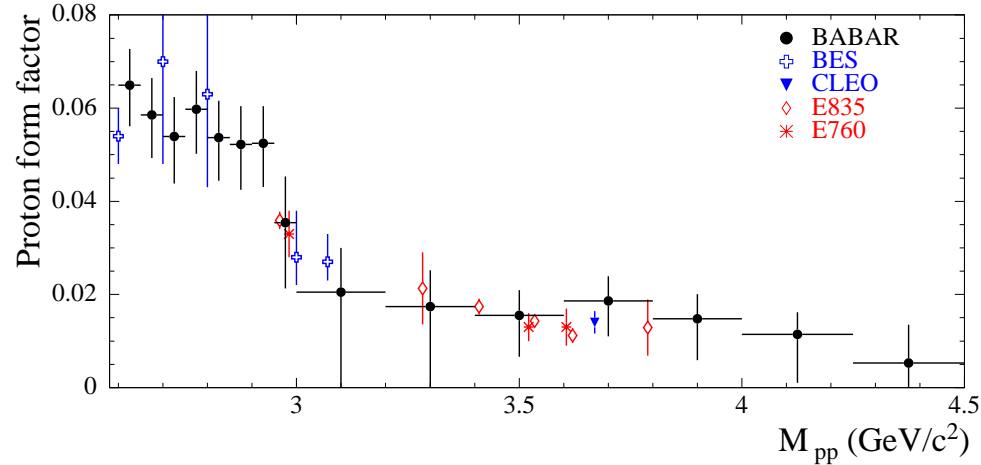
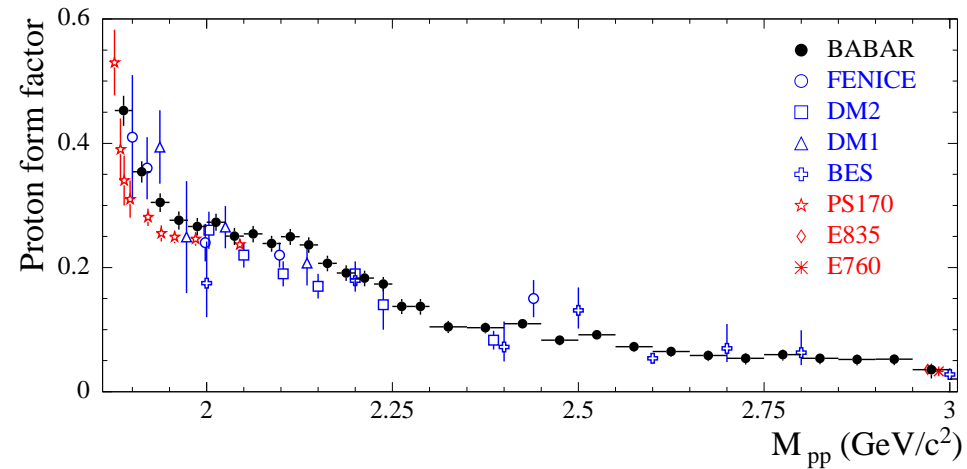
- similar results for **neutron** pair or $\Lambda\bar{\Lambda}$ production
- NLO corrections from ISR included (corrections $\sim 1-2\%$)
- no FSR
- $\Lambda\bar{\Lambda}$ can be studied with analysis of Λ -polarization

thousands of events around $4-5 \text{ GeV}^2$

several events up to $7-8 \text{ GeV}^2$

Results (from BABAR)

PRD 73, 012005 (2006)



Recent improvements

- model amplitudes

- Improved formfactors for

$$\pi^+ \pi^- 2\pi^0, \quad 2\pi^+ 2\pi^-,$$

$$\pi^+ \pi^- \quad (\text{high } Q^2 \text{ region})$$

$$K^+ K^-, K^0 \bar{K}^0$$

- narrow resonances ($J/\psi, \psi'$)
(direct vs. indirect amplitudes)

V MESON FORM FACTORS at LARGE Q^2

Khodjamirian, JK, EPJ C 39 (2004) 41

Czyz, Grzelinska, JK, PRD 81, 094014 (2010)

radiative return will explore large Q^2

convenient representation for F_π :

generalized VDM with ρ, ρ', \dots

combined with Veneziano-type tower of resonances (Dominguez)

$$F_\pi(s) = \sum_{n=0}^{\infty} c_n \frac{m_n^2}{m_n^2 - s},$$
$$c_n = \frac{(-1)^n \Gamma(\beta - 1/2)}{\sqrt{\pi} (\frac{1}{2} + n) \Gamma(n + 1) \Gamma(\beta - 1 - n)},$$
$$m_n^2 = m_\rho^2 (1 + 2n),$$

β = free parameter

$$\sum c_n = 1$$

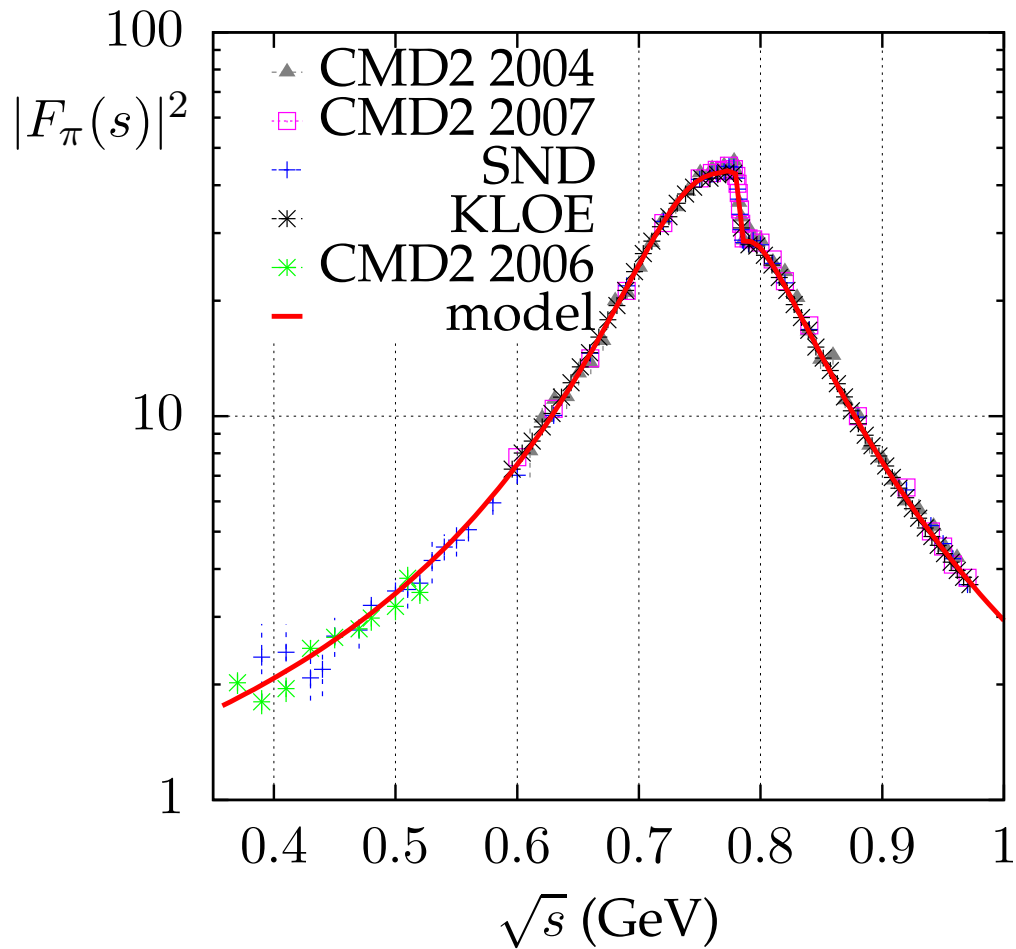
Modifications:

- finite widths
- parameters of ρ , ρ' , ρ'' fitted to data
- Breit-Wigner for ρ , ρ' , ρ'' with Q^2 -dependent widths
 \Rightarrow reasonable agreement between model and fit

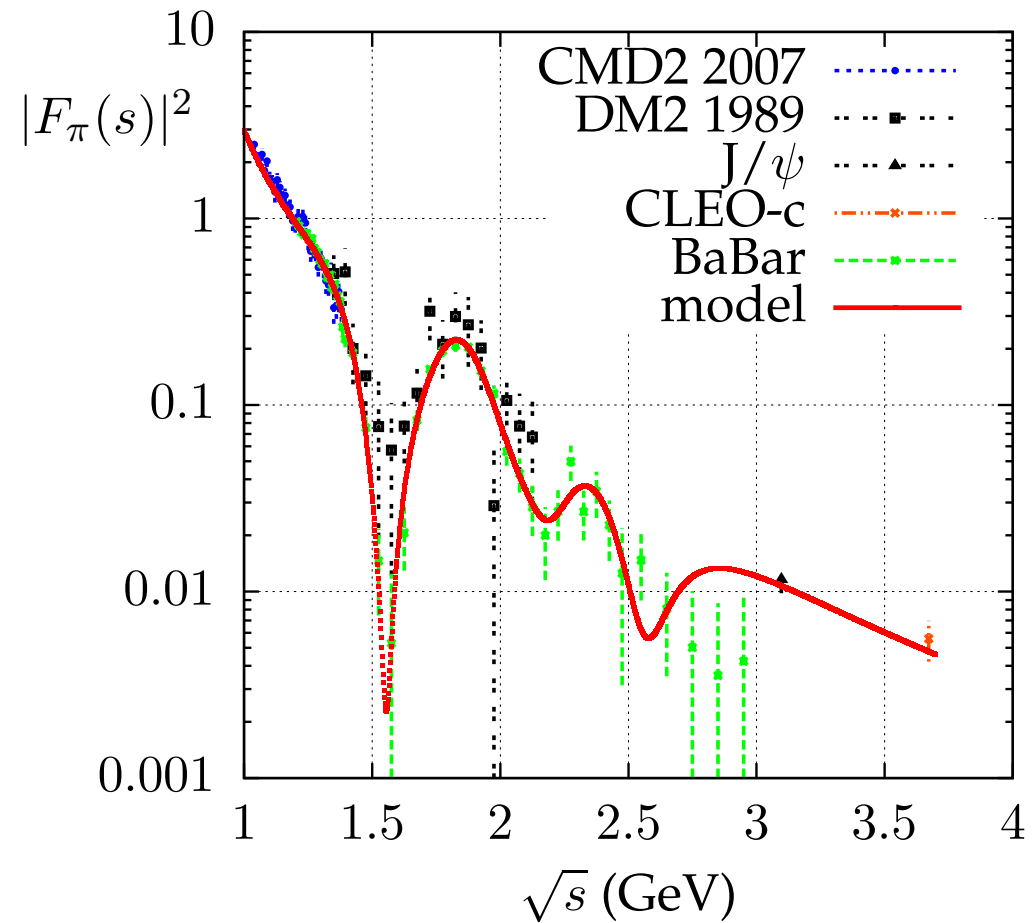
Parameter	model(fit)	PDG value	model
m_{ρ_0}	773.37 ± 0.19	775.49 ± 0.34	input
Γ_{ρ_0}	147.1 ± 1.0	149.4 ± 1.0	input
m_{ω}	782.4 ± 0.5	782.41 ± 0.12	-
Γ_{ω}	8.33 ± 0.27	8.49 ± 0.08	-
m_{ρ_1}	1490 ± 11	1465 ± 25	1340
Γ_{ρ_1}	429 ± 27	400 ± 60	256
m_{ρ_2}	1870 ± 25	1720 ± 20	1730
Γ_{ρ_2}	357 ± 46	250 ± 100	330
m_{ρ_3}	2120	-	2047
Γ_{ρ_3}	300	-	391
m_{ρ_4}	model	-	2321
Γ_{ρ_4}	model	-	444
m_{ρ_5}	model	-	2567
Γ_{ρ_5}	model	-	491

Parameter	model(fit)	PDG value	model
β	2.148 ± 0.003	-	input
$ c_{\omega}^{\pi} $	$(18.7 \pm 0.5) \cdot 10^{-4}$	-	-
$Arg(c_{\omega}^{\pi})$	0.106 ± 0.020	-	-
$ F_2 $	0.59 ± 0.10	-	-
$Arg(F_2)$	-2.20 ± 0.16	-	-
$ F_3 $	0.048 ± 0.056	-	-
$Arg(F_3)$	$-2. \pm 1.4$	-	-
$ F_4 $	0.40 ± 0.07	-	-
$Arg(F_4)$	-2.9 ± 0.3	-	-
$ F_5 $	0.43 ± 0.05	-	-
$Arg(F_5)$	1.19 ± 0.18	-	-
$\chi^2/d.o.f.$	271/270	-	-

$$e^+e^- \rightarrow \pi^+\pi^-$$



$$e^+e^- \rightarrow \pi^+\pi^-$$



data point at 3.1 GeV ($J/\Psi \rightarrow \pi\pi$) leads to strong constraints

$$e^+e^- \rightarrow K^+K^-, K^0\bar{K}^0$$

isospin symmetry:

$$F_{K^+} = +F^{(I=1)} + F^{(I=0)}$$

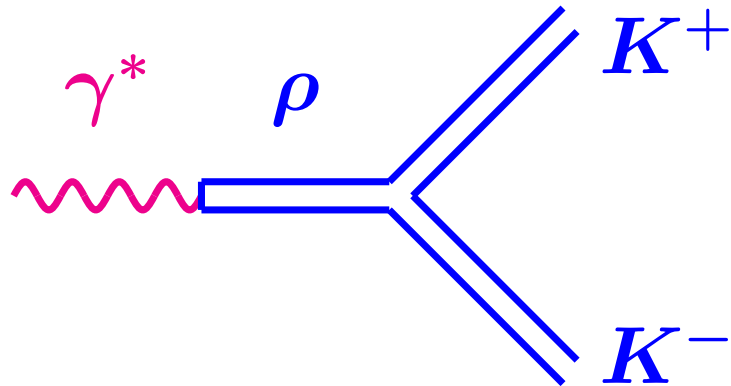
$$F_{K^0} = -F^{(I=1)} + F^{(I=0)}$$

resonances:

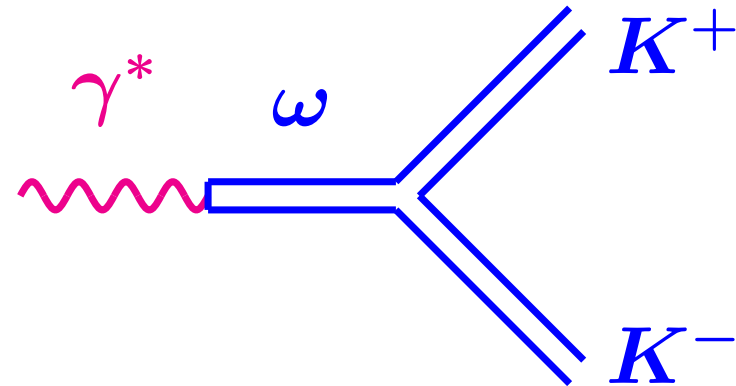
$$\begin{aligned} F_{K^+}(s) = & +\frac{1}{2} \left(c_\rho^K BW_\rho(s) + c_{\rho'}^K BW_{\rho'}(s) + c_{\rho''}^K BW_{\rho''}(s) \right) \\ & + \frac{1}{6} \left(c_\omega^K BW_\omega(s) + c_{\omega'}^K BW_{\omega'}(s) + c_{\omega''}^K BW_{\omega''}(s) \right) \\ & + \frac{1}{3} \left(c_\phi^K BW_\phi(s) + c_{\phi'}^K BW_{\phi'}(s) \right), \end{aligned}$$

$$\begin{aligned} F_{K^0}(s) = & -\frac{1}{2} \left(c_\rho^K BW_\rho(s) + c_{\rho'}^K BW_{\rho'}(s) + c_{\rho''}^K BW_{\rho''}(s) \right) \\ & + \frac{1}{6} \left(c_\omega^K BW_\omega(s) + c_{\omega'}^K BW_{\omega'}(s) + c_{\omega''}^K BW_{\omega''}(s) \right) \\ & + \frac{1}{3} \left(\eta_\phi c_\phi^K BW_\phi(s) + c_{\phi'}^K BW_{\phi'}(s) \right) \end{aligned}$$

quark model:



$$\frac{1}{\sqrt{2}} f_\rho \quad g_{\rho KK}$$

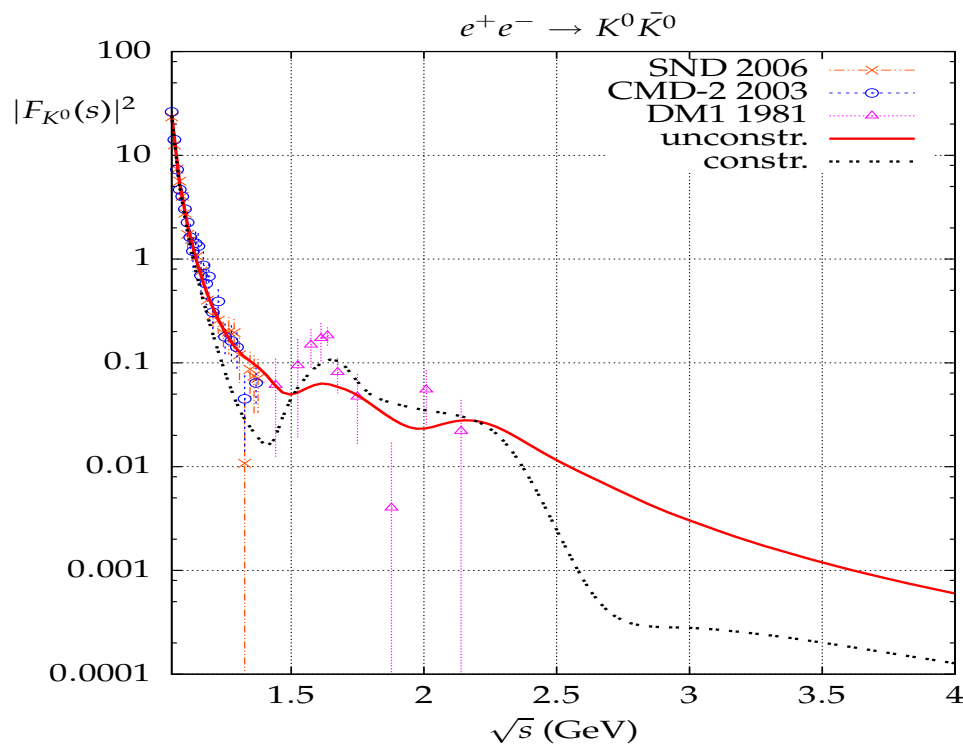
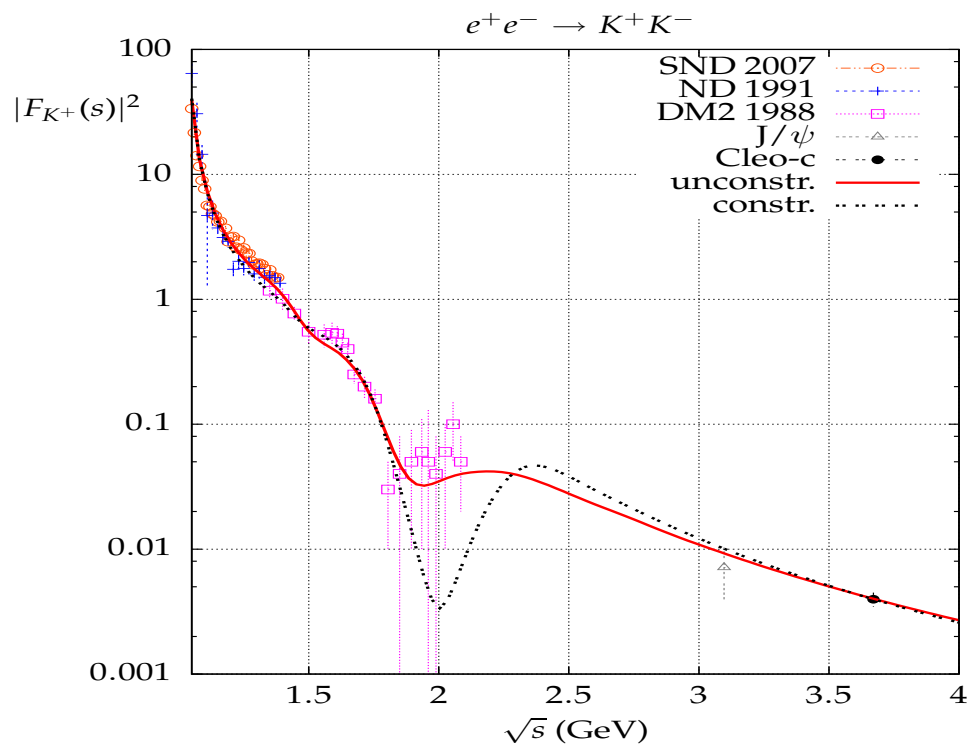
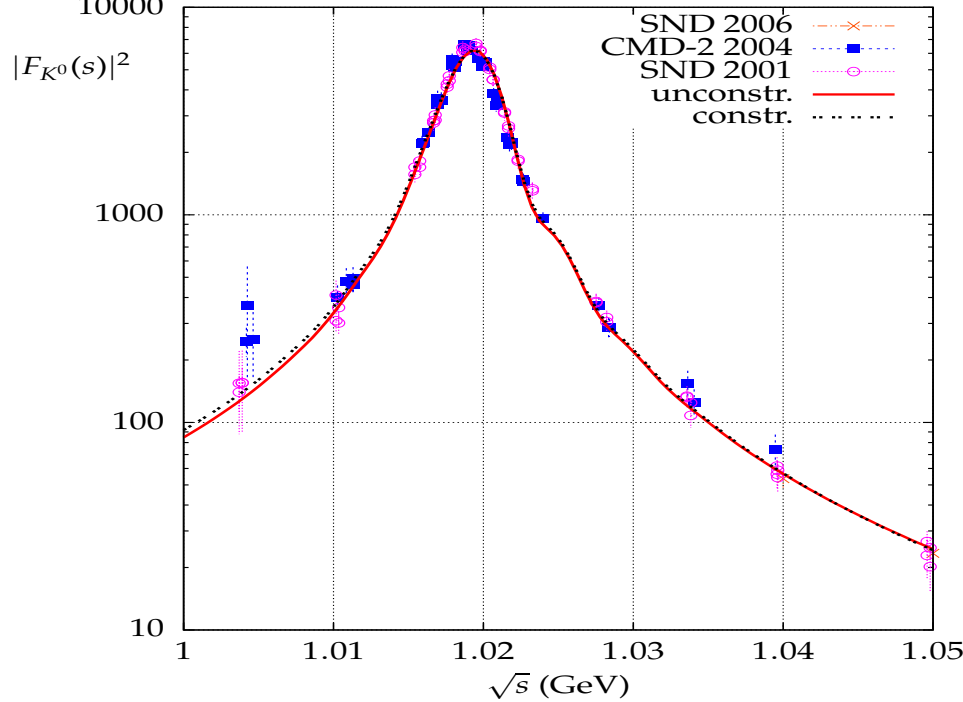
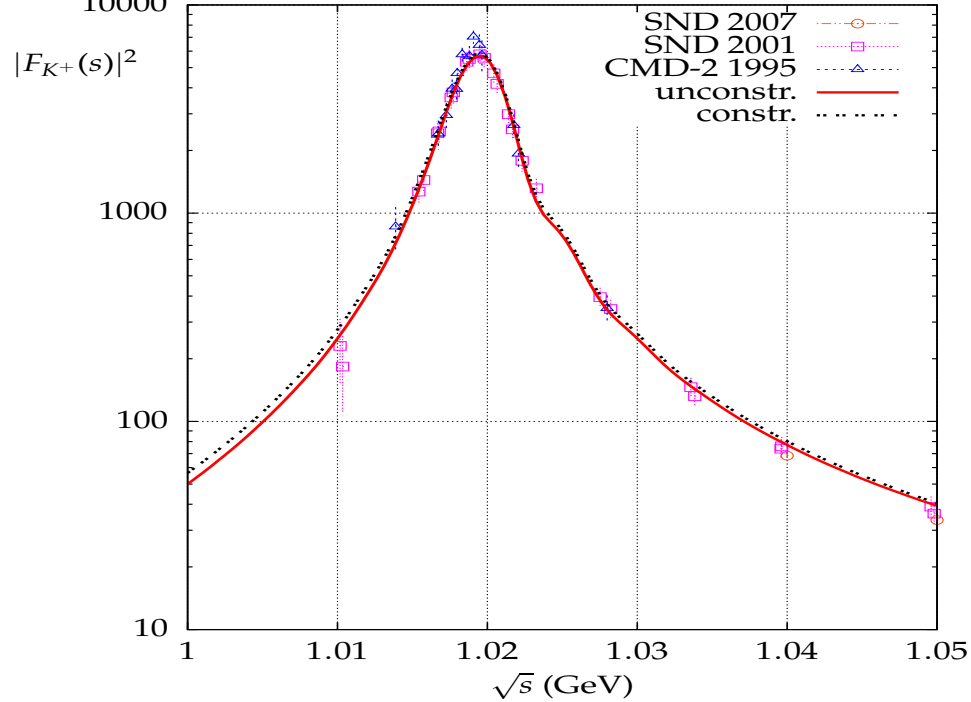


$$\frac{1}{3\sqrt{2}} f_\omega \quad g_{\omega KK}$$

constraint: $f_\rho = f_\omega, \quad g_{\rho KK} = g_{\omega KK}$

$$\Rightarrow c_\rho = c_\omega$$

fit performed with **(solid curves)**
 or **without (dashed curves)** this constraint



$$\tau \rightarrow K^- K^0 \nu$$

Predictions based on isospin symmetry and $I = 1$ part of form factor:

$$\left(\frac{1}{BR(\tau \rightarrow \mu^- \bar{\nu}_\mu \nu_\tau)} \right) \frac{dBR(\tau \rightarrow K^- K^0 \nu_\tau)}{d\sqrt{Q^2}} =$$

$$\frac{|V_{ud}|^2}{2m_\tau^2} \left(1 + \frac{2Q^2}{m_\tau^2} \right) \left(1 - \frac{Q^2}{m_\tau^2} \right)^2 \left(1 - \frac{4m_K^2}{Q^2} \right)^{3/2}$$

$$\times \sqrt{Q^2} |F_{K^- K^0}(Q^2)|^2$$

and $F_{K^- K^0} = -F_{K^+} + F_{K^0}$

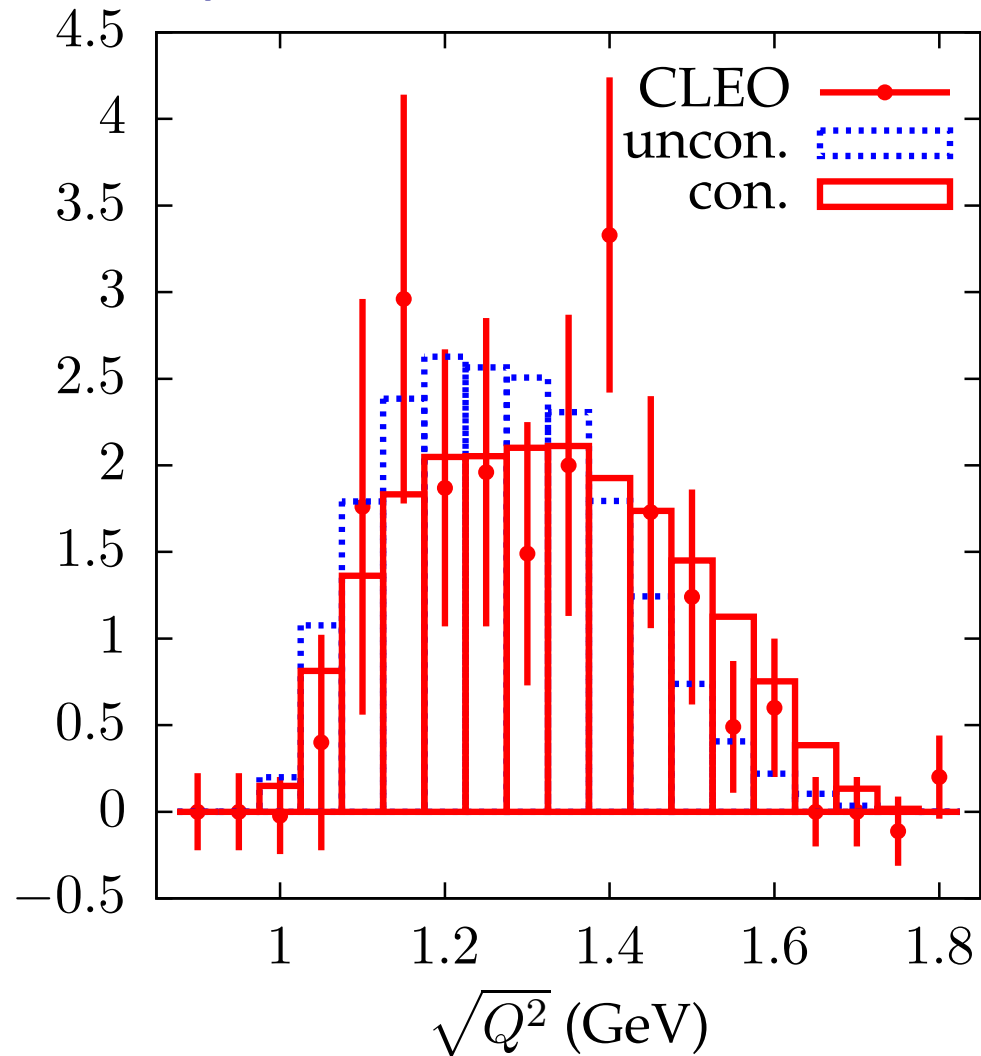
$$\Rightarrow BR(\tau \rightarrow K^- K^0 \nu_\tau) = 0.135 - 0.190\%$$

to be compared with

$$BR(\tau \rightarrow K^- K^0 \nu_\tau) = 0.158 \pm 0.016\%.$$

Q^2 distribution:

will provide further constraints!



(data from CLEO)

VI Experimental Results

KLOE

pion form factor, asymmetry

BABAR, BELLE

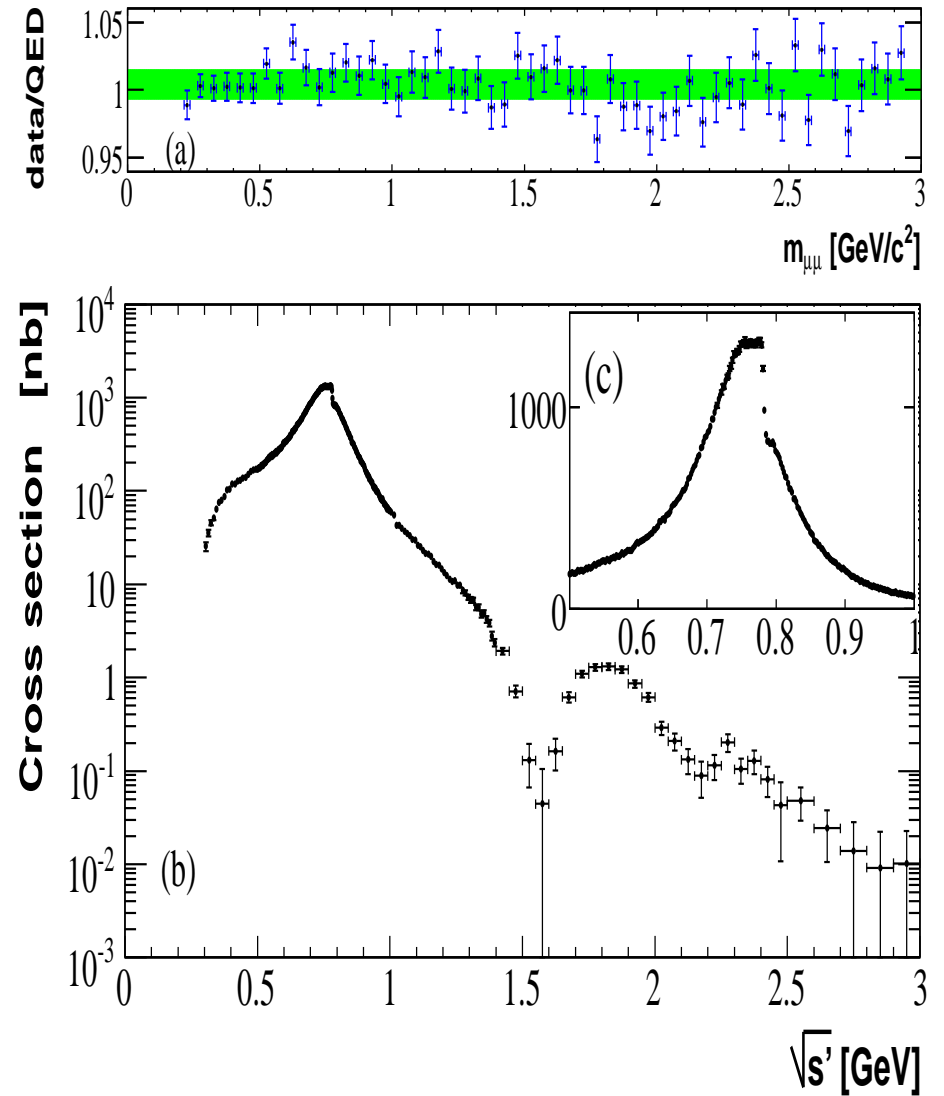
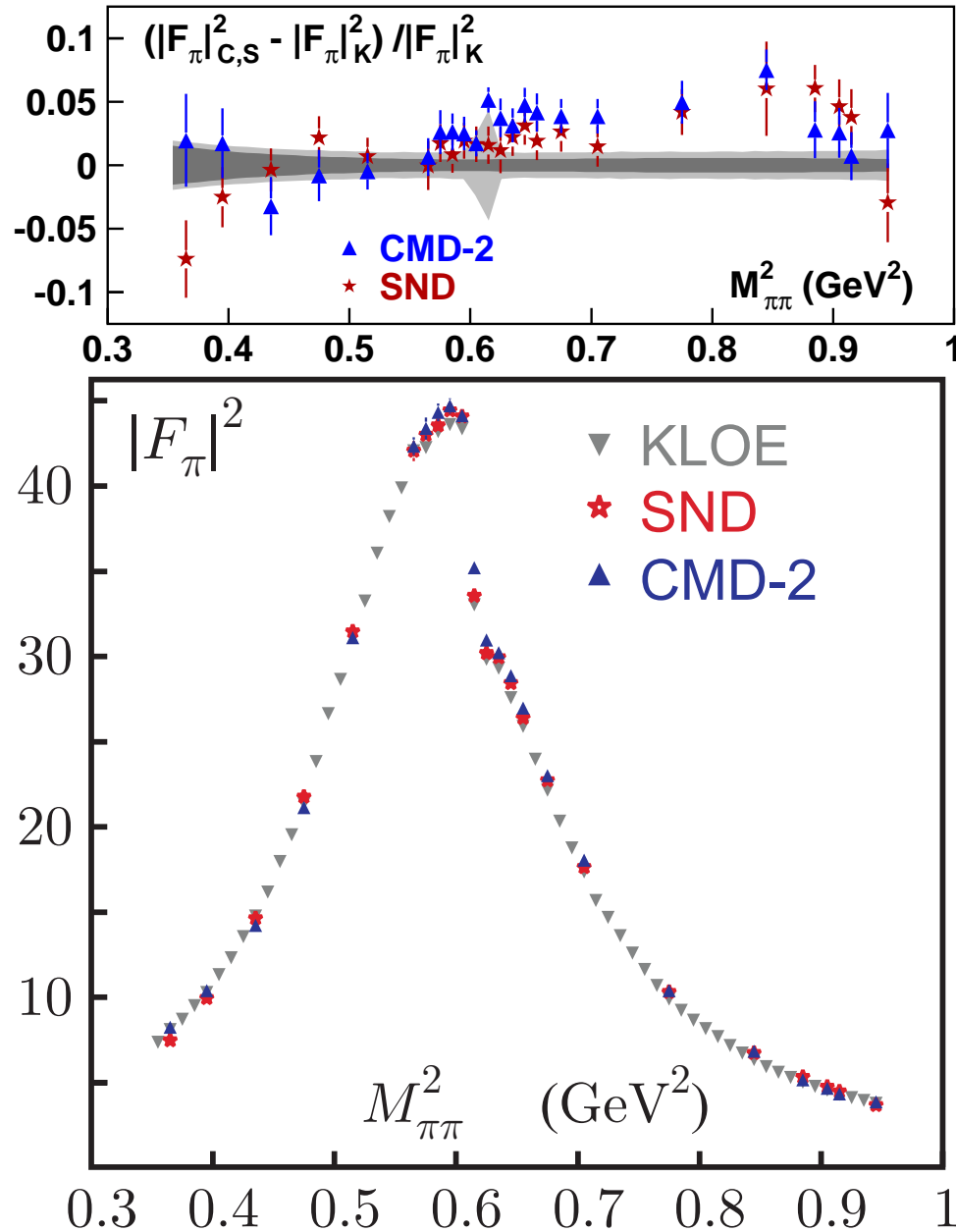
higher Q^2 available

⇒ measurement of $R(Q^2)$ from threshold up to at least 5 GeV.

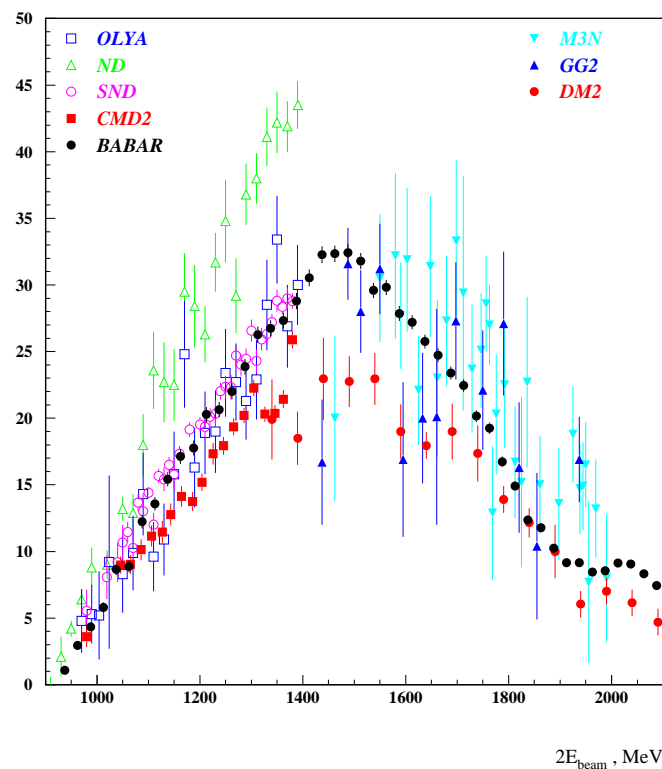
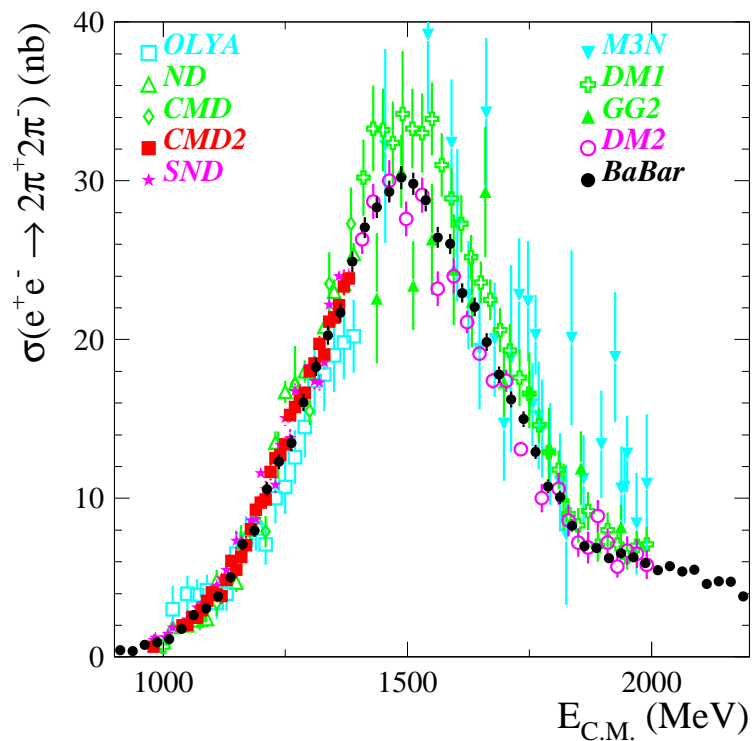
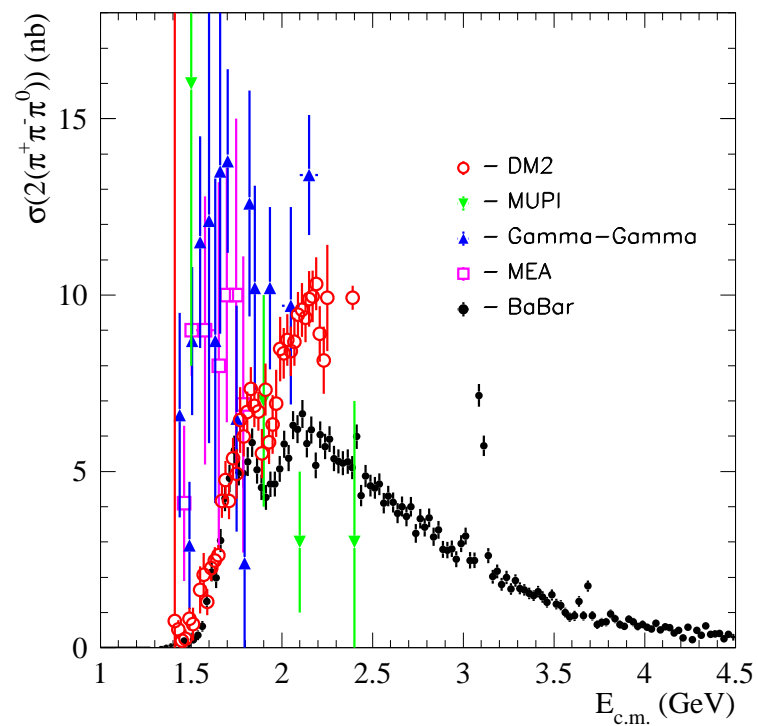
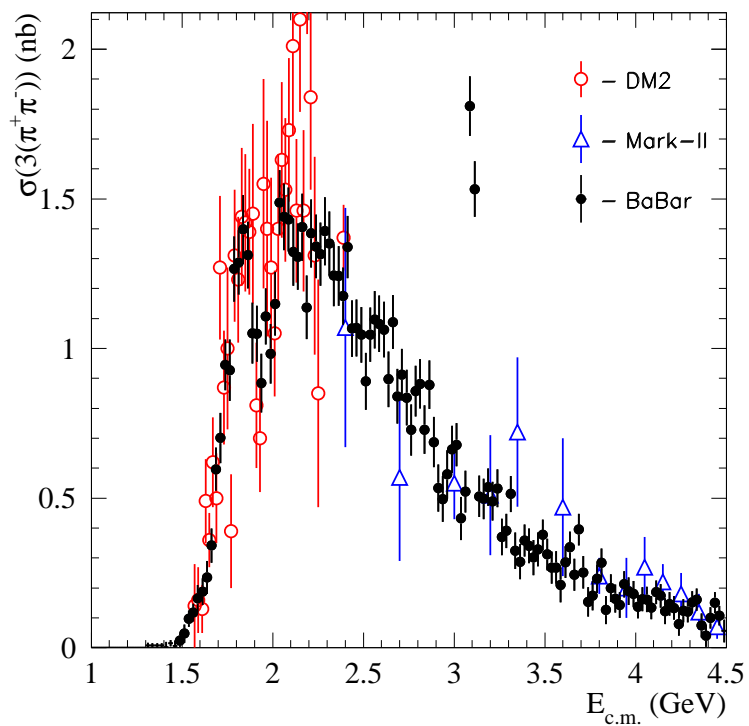
Examples:

- ▶ $\pi\pi$
- ▶ $4\pi^\pm, \pi^+\pi^-2\pi^0$
- ▶ 6π
- ▶ $K K \pi\pi$
- ▶ $K K K K$
- ▶ $p\bar{p}, \Lambda\bar{\Lambda}, \dots$

Pion Form Factor



BABAR

4π  6π 

VII Conclusions

- continuous development of **PHOKHARA**
 - ⇒ radiative corrections
 - ⇒ more channels
 - ⇒ cooperation between theory and experiment crucial
 - charge asymmetry as analysis tool
 - nucleon form factors:
 G_E and G_M can be measured for a wide range of Q^2
 - pion form factor: structures at large Q^2
kaon form factors: K^+K^- & $K^0\bar{K}^0 \Rightarrow K^-K^0$
⇒ prediction for $\tau \rightarrow \nu K^- K^0$
- ⇒ numerous experimental results during past 5 years