
Strange, Charm and Bottom masses at NNLO and N^3LO

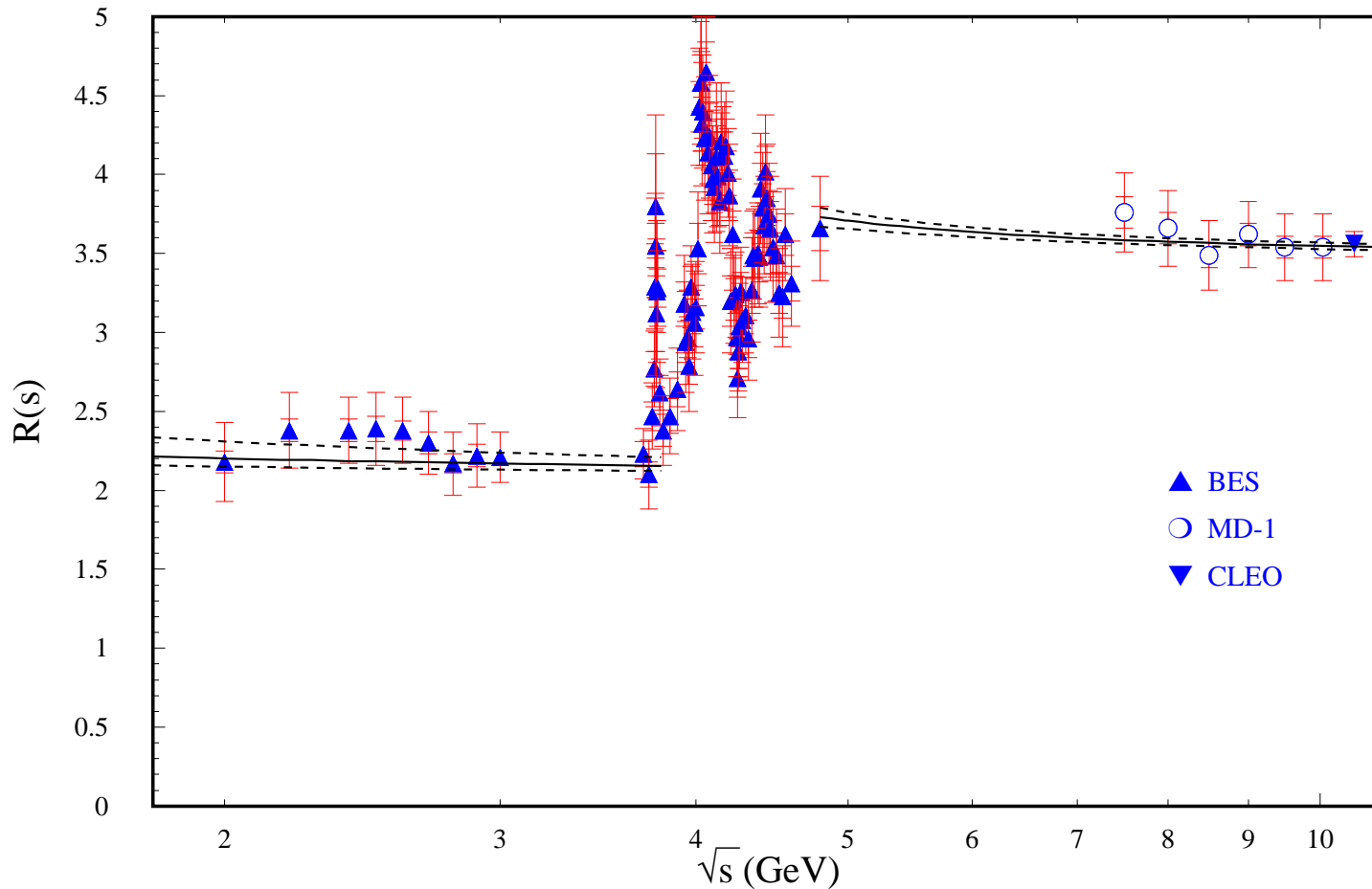
J.H. Kühn



SFB TR-9

- I. Sum Rules with Charm and Bottom Quarks:
recent data and N^3LO calculations
- II. m_c and m_b : overview
- III. m_s from τ -decays and sum rules

I Sum Rules with Charm and Bottom Quarks



pQCD and data agree well in the regions
2 — 3.73 GeV and 5 — 10.52 GeV

experiment	energy [GeV]	date	systematic error
BES	2 — 5	2001	4%
MD-1	7.2 — 10.34	1996	4%
CLEO	10.52	1998	2%
PDG	J/ψ		(7%) 3%
PDG	ψ'		(9%) 5.7%
PDG	ψ''		15%

m_Q from SVZ Sum Rules, Moments and Tadpoles

Some definitions

$$R(s) = 12\pi \operatorname{Im} \left[\Pi(q^2 = s + i\epsilon) \right]$$

$$\left(-q^2 g_{\mu\nu} + q_\mu q_\nu \right) \Pi(q^2) \equiv i \int dx e^{iqx} \langle T j_\mu(x) j_\nu(0) \rangle$$

with the electromagnetic current j_μ

$$\text{Taylor expansion: } \Pi_Q(q^2) = Q_Q^2 \frac{3}{16\pi^2} \sum_{n \geq 0} \bar{C}_n z^n$$

with $z = q^2/(4m_Q^2)$ and $m_Q = m_Q(\mu)$ the $\overline{\text{MS}}$ mass.

Coefficients \bar{C}_n up to $n = 8$ known analytically in order α_s^2

(Chetyrkin, JK, Steinhauser)

recently up to $n = 30!$

(Boughezal, Czakon, Schutzmeier)

recently also \bar{C}_0 and \bar{C}_1 in order α_s^3 (four loops!)

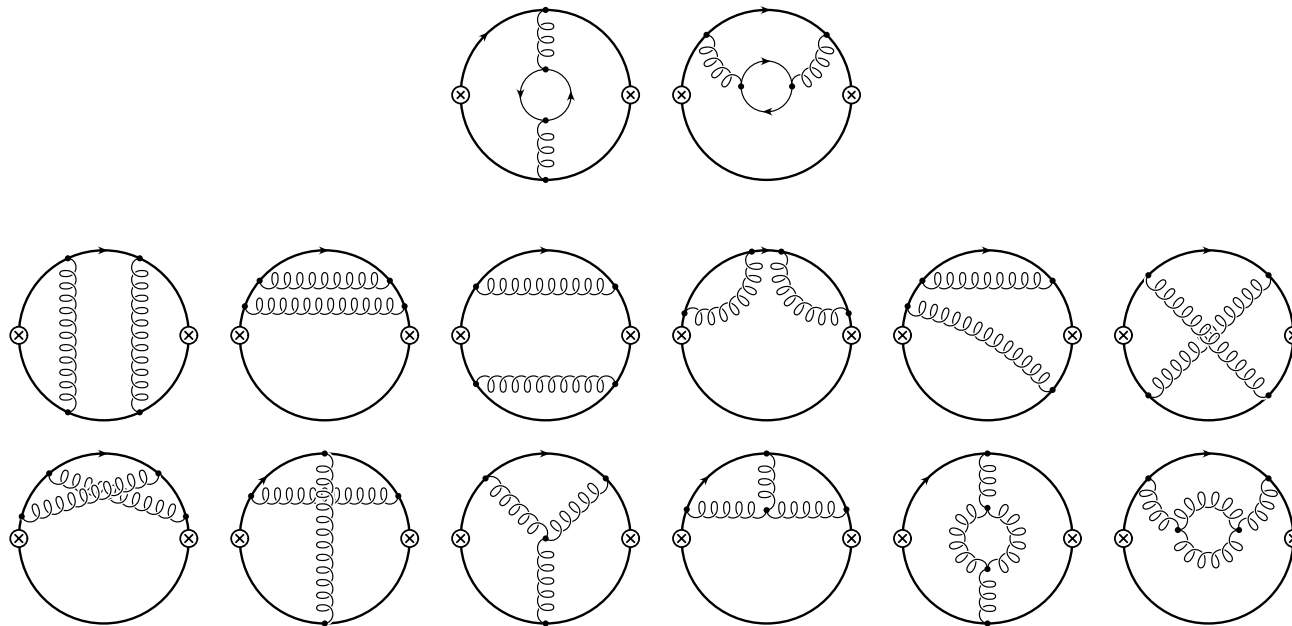
⇒ talk at session V (QCD hard interactions)

(Chetyrkin, JK, Sturm; confirmed by Boughezal, Czakon, Schutzmeier)

Analysis in NNLO

Coefficients \bar{C}_n from three-loop one-scale tadpole amplitudes with
“arbitrary” power of propagators;

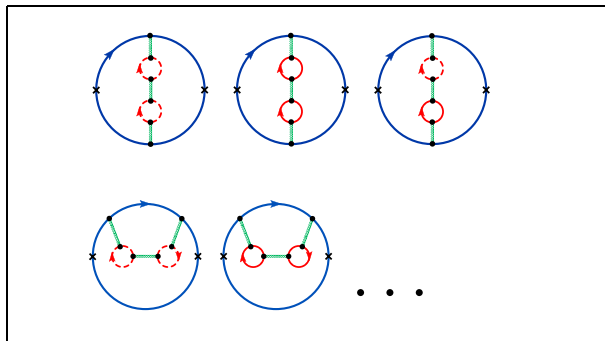
FORM-program MATAD



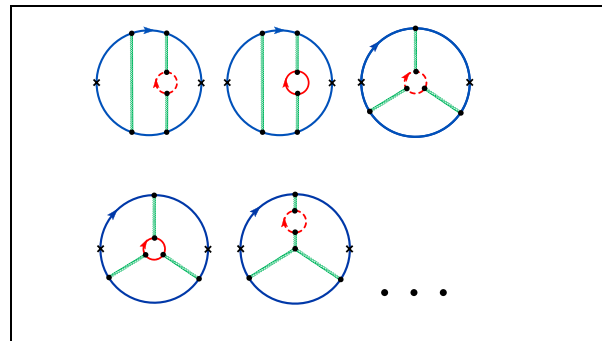
Analysis in N³LO

Algebraic reduction to 13 master integrals (Laporta algorithm);
numerical evaluation of master integrals

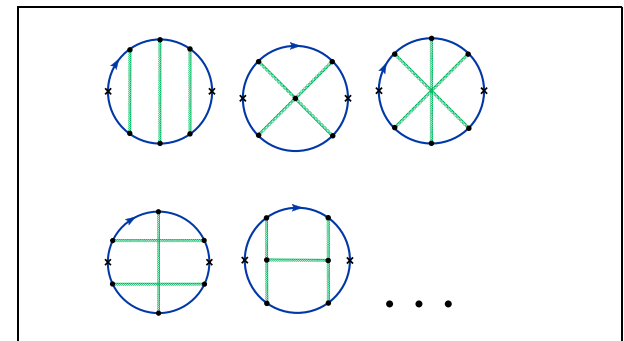
n_f^2 -contributions





n_f^1 -contributions



n_f^0 -contributions



 : heavy quarks,  : light quarks,

n_f : number of active quarks

\implies About 700 Feynman-diagrams

\bar{C}_n depend on the charm quark mass through $l_{m_c} \equiv \ln(m_c^2(\mu)/\mu^2)$

$$\begin{aligned}\bar{C}_n &= \bar{C}_n^{(0)} + \frac{\alpha_s(\mu)}{\pi} \left(\bar{C}_n^{(10)} + \bar{C}_n^{(11)} l_{m_c} \right) \\ &+ \left(\frac{\alpha_s(\mu)}{\pi} \right)^2 \left(\bar{C}_n^{(20)} + \bar{C}_n^{(21)} l_{m_c} + \bar{C}_n^{(22)} l_{m_c}^2 \right) \\ &+ \left(\frac{\alpha_s(\mu)}{\pi} \right)^3 \left(\dots \right)\end{aligned}$$

n	1	2	3	4
$\bar{C}_n^{(0)}$	1.0667	0.4571	0.2709	0.1847
$\bar{C}_n^{(10)}$	2.5547	1.1096	0.5194	0.2031
$\bar{C}_n^{(11)}$	2.1333	1.8286	1.6254	1.4776
$\bar{C}_n^{(20)}$	2.4967	2.7770	1.6388	0.7956
$\bar{C}_n^{(21)}$	3.3130	5.1489	4.7207	3.6440
$\bar{C}_n^{(22)}$	-0.0889	1.7524	3.1831	4.3713

Define the moments

$$\mathcal{M}_n^{\text{th}} \equiv \frac{12\pi^2}{n!} \left(\frac{d}{dq^2} \right)^n \Pi_c(q^2) \Big|_{q^2=0} = \frac{9}{4} Q_c^2 \left(\frac{1}{4m_c^2} \right)^n \bar{C}_n$$

dispersion relation:

$$\Pi_c(q^2) = \frac{q^2}{12\pi^2} \int ds \frac{R_c(s)}{s(s - q^2)} + \text{subtraction}$$

$$\Leftrightarrow \mathcal{M}_n^{\text{exp}} = \int \frac{ds}{s^{n+1}} R_c(s)$$

constraint: $\mathcal{M}_n^{\text{exp}} = \mathcal{M}_n^{\text{th}}$

$\Leftrightarrow m_c$

SVZ:

$\mathcal{M}_n^{\text{th}}$ can be reliably calculated in pQCD: low n :

- fixed order in α_s is sufficient, in particular no resummation of $1/v$ - terms from higher orders required
- condensates are unimportant
- pQCD in terms of short distance mass : $m_c(3 \text{ GeV}) \Leftrightarrow m_c(m_c)$
stable expansion : no pole mass or closely related definition
(1S-mass, potential-subtracted mass) involved
- moments available in NNLO
- and \bar{C}_0, \bar{C}_1 in N³LO

anatomy of errors and update: m_c

old results (NPB 619(2001) and

new results (update on $\Gamma_e(J/\psi, \psi')$ and $\alpha_s = 0.1187 \pm 0.0020$)

	$J/\psi, \psi'$	charm threshold region	continuum	sum
n	$\mathcal{M}_n^{\text{exp, res}}$ $\times 10^{(n-1)}$	$\mathcal{M}_n^{\text{exp, cc}}$ $\times 10^{(n-1)}$	$\mathcal{M}_n^{\text{cont}}$ $\times 10^{(n-1)}$	$\mathcal{M}_n^{\text{exp}}$ $\times 10^{(n-1)}$
1	0.1114(82)	0.0313(15)	0.0638(10)	0.2065(84)
1	0.1138(40)	0.0313(15)	0.0639(10)	0.2090(44)
2	0.1096(79)	0.0174(8)	0.0142(3)	0.1412(80)
2	0.1121(38)	0.0174(8)	0.0142(3)	0.1437(39)

new analysis

$$\boxed{n = 1} \quad N^2LO$$

$$m_c(3 \text{ GeV}) = 1017 \pm 14(\text{exp}) \pm 7(\alpha_s) \pm 1.3(\text{scale})$$

$$\Leftrightarrow m_c(m_c) = 1301 \pm 15 \quad (\text{3-loop-running})$$

$$\boxed{n = 1} \quad N^3LO$$

$$m_c(3 \text{ GeV}) = 1012 \pm 14(\text{exp}) \pm 8(\alpha_s) \pm 0.6(\text{scale})$$

$$\Leftrightarrow m_c(m_c) = 1297 \pm 15 \quad (\text{4-loop-running})$$

$$\boxed{n = 2} \quad N^2LO$$

$$m_c(3 \text{ GeV}) = 982 \pm 10(\text{exp}) \pm 17(\alpha_s) \pm 19(\text{scale})$$

$$\Leftrightarrow m_c(m_c) = 1269 \pm 25 \quad (\text{3-loop-running})$$

$$\text{combine} \quad \Leftrightarrow m_c(3 \text{ GeV}) = 1007 \pm 16$$

$$m_c(m_c) = 1290 \pm 15, \quad \text{old result } 1304 \pm 27$$

anatomy of errors and update: m_b

old results (NPP 619 (2001)) and update (Γ_e from CLEO);

$$\alpha_s = 0.1187 \pm 0.0020$$

n	$\mathcal{M}_n^{\text{exp,res}}$ $\times 10^{(2n+1)}$	$\mathcal{M}_n^{\text{exp,thr}}$ $\times 10^{(2n+1)}$	$\mathcal{M}_n^{\text{cont}}$ $\times 10^{(2n+1)}$	$\mathcal{M}_n^{\text{exp}}$ $\times 10^{(2n+1)}$
1	1.237(63)	0.306(62)	2.913(21)	4.456(121)
1	1.271(24)	0.306(62)	2.918(16)	4.494(84)
2	1.312(65)	0.261(54)	1.182(12)	2.756(113)
2	1.348(25)	0.261(52)	1.185(9)	2.795(75)
3	1.399(68)	0.223(44)	0.634(8)	2.256(108)
3	1.437(26)	0.223(44)	0.636(6)	2.296(68)

new analysis

$$\boxed{n = 1} \quad N^2LO$$

$$m_b(10 \text{ GeV}) = 3644 \pm 40(\text{exp}) \pm 7(\alpha_s) \pm 4(\text{scale})$$

$$\Leftrightarrow m_b(m_b) = 4191 \pm 40 \quad (\text{3-loop-running})$$

$$\boxed{n = 1} \quad N^3LO$$

$$m_b(10 \text{ GeV}) = 3644 \pm 40(\text{exp}) \pm 7(\alpha_s) \pm 1(\text{scale})$$

$$\Leftrightarrow m_b(m_b) = 4191 \pm 40 \quad (\text{4-loop-running})$$

$$\boxed{n = 2} \quad N^2LO$$

$$m_b(10 \text{ GeV}) = 3631 \pm 29(\text{exp}) \pm 13(\alpha_s) \pm 17(\text{scale})$$

$$\Leftrightarrow m_b(m_b) = 4179 \pm 35 \quad (\text{3-loop-running})$$

$$\boxed{n = 3} \quad N^2LO$$

$$m_b(10 \text{ GeV}) = 3622 \pm 22(\text{exp}) \pm 17(\alpha_s) \pm 20(\text{scale})$$

$$\Leftrightarrow m_b(m_b) = 4170 \pm 33 \quad (\text{3-loop-running})$$

$$\text{combine} \quad \Leftrightarrow m_b(10 \text{ GeV}) = 3630 \pm 35$$

$$m_b(m_b) = 4180 \pm 35, \quad \text{old result } 4210 \pm 50$$

Summary on m_c and m_b

- ⇒ drastic improvement in δm_c , δm_b from moments with low n in N^2LO
- ⇒ direct determination of short-distance mass

improved measurements of $\Gamma_e(J/\psi, \psi')$ and $\Gamma_e(\Upsilon, \Upsilon', \Upsilon'')$
and new N^3LO results lead to significant improvements

preliminary results:

$$m_c(3 \text{ GeV}) = 1.007(16) \text{ GeV}$$

$$m_c(m_c) = 1.300(15) \text{ GeV}$$

$$m_b(10 \text{ GeV}) = 3.630(35) \text{ GeV}$$

$$m_b(m_b) = 4.179(35) \text{ GeV}$$

II m_c and m_b : other characteristic results

no review

charm

- moments of B -decay distributions (hadron mass, lepton energy)

HQE up to $\mathcal{O}(1/m_b^3)$, pQCD up to $\mathcal{O}(\alpha_s^2\beta_0)$

1240 ± 70

Buchmüller, Flächer

$1224 \pm 17 \pm 54$

Hoang, Manohar

- Lattice, from D_s (quenched $\Leftrightarrow \pm(40 - 60)$)

$1260 \pm 40 \pm 120$

Becirevic, Lubicz, Martinelli

1301 ± 34

Rolf, Sint

bottom

- moments of B -decay distributions (hadron mass, lepton energy)

HQE up to $\mathcal{O}(1/m_b^3)$, pQCD up to $\mathcal{O}(\alpha_s^2\beta_0)$

4200 ± 40

Buchmüller, Flächer

4170 ± 30

Bauer et al.

- Υ -spectroscopy (1S-state), pNRQCD + nonperturbative effects

4346 ± 70

Penin, Steinhauser (N^3LO)

$4210 \pm 90 \pm 25$

Pineda (N^2LO)

- Lattice (HQET + $1/m_b$ terms, quenched)

4301 ± 70

ALPHA-Coll. Delta Morte et al.

III m_s from τ -decays and sum rules

$$\tau \rightarrow \nu s \bar{d}$$

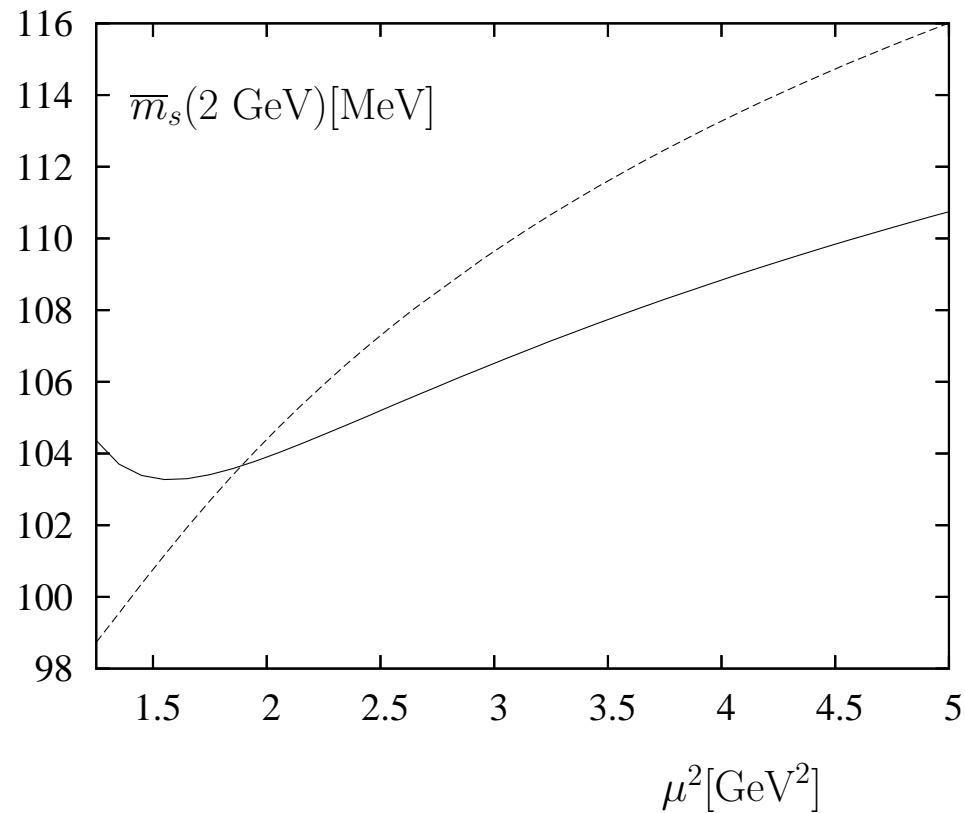
input: moments of $m(s\bar{d})$ (ALEPH, OPAL)
 V_{us} (Czarnecki, Marciano, Sirlin)
phenomenology (Gamiz et al)
pQCD in $\mathcal{O}(\alpha^3)$ (Baikov, Chetyrkin, JK)
(finite part of massless four-loop correlator)

$$\Leftrightarrow m_s(M_\tau) = 100 \pm \begin{pmatrix} +5 \\ -3 \end{pmatrix}_{\text{theo}} \pm \begin{pmatrix} +17 \\ -19 \end{pmatrix}_{\text{rest}}$$

pseudoscalar sum rules in $\mathcal{O}(\alpha_s^4)$

$$\bar{m}_s(2 \text{ GeV}) = 105 \pm 6(\text{param}) \pm 7(\text{had})$$

Chetyrkin, Khodjamirian



Method	$\overline{m}_s(2 \text{ GeV})$ [MeV]	Ref.
Pseudoscalar Borel sum rule	$105 \pm 6 \pm 7$	Chetyrkin
Pseudoscalar FESR	100 ± 12	Maltman
Scalar Borel sum rule	99 ± 16	Jamin
Vector FESR	139 ± 31	Eidemüller
Spectral function	> 77	Baikov
Hadronic τ decays	81 ± 22	Gamiz
	96_{-3}^{+5+16}	Baikov
	104 ± 28	Narison
τ decays \oplus sum rules	99 ± 28	Narison
Lattice QCD ($n_f = 2$)	97 ± 22	Della Morte
	100 -130	Gockeler
	$101 \pm 8_{-0}^{+25}$	Becirevic
Lattice QCD ($n_f = 3$)	$76 \pm 3 \pm 7$	Aubin
	86.7 ± 5.9	Ishikawa
	$87 \pm 4 \pm 4$	Mason
PDG04 average	80 -130	Eidelman

Summary

new multiloop results from pQCD + improved data

$$m_c(3 \text{ GeV}) = 1007 \pm 16 \text{ MeV} \quad m_c(m_c) = 1290 \pm 15 \text{ MeV}$$

$$m_b(10 \text{ GeV}) = 3630 \pm 35 \text{ MeV} \quad m_b(m_b) = 4180 \pm 35 \text{ MeV}$$

significantly reduced errors, consistent with other determinations,
but more precise

$$m_s(2 \text{ GeV}) = 105 \pm 10 \text{ MeV}$$

on the basis of N^3LO pseudoscalar sumrules