

Introduction to Theoretical Particle Physics

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Exercise Sheet 8

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Exercise 19: Spontaneous symmetry breaking

7 points

Consider a theory with N real scalar fields governed by a Lagrangian

$$\mathcal{L} = \frac{1}{2}(\partial_\mu \Phi)^T (\partial^\mu \Phi) - V(\Phi^T \Phi), \quad (19.1)$$

where Φ is a vector, $\Phi = (\phi_1, \dots, \phi_N)$. Let the potential be

$$V(\Phi^T \Phi) = -\frac{\mu^2}{2}(\Phi^T \Phi) + \frac{\lambda}{4}(\Phi^T \Phi)^2, \quad (19.2)$$

where μ and λ are positive constants. This potential is manifestly symmetric under the $SO(N)$ group.

(a) Let $R = (\Phi^T \Phi)$. Find the minimum, R_{vac} , of the potential in Eq. (19.2).

Consider a possible pattern of symmetry breaking where one of the scalar fields obtains a non-zero expectation value. In this case one can write

$$\phi_1 = v + \chi_1, \quad \phi_i = \chi_i \text{ for } i \in \{2, \dots, N\}. \quad (19.3)$$

where v is a constant and χ_i describe small excitations around the minimum of the potential.

(b) Express v in terms of R_{vac} .

(c) Rewrite the Lagrangian in terms of the fields χ_i . What is the symmetry of the Lagrangian after the symmetry breaking? How many Goldstone bosons do you expect? How many massive scalar fields are present?

Now think of a different pattern of symmetry breaking, where two fields obtain a non-zero expectation value. In this case one can write

$$\phi_{1,2} = v_{1,2} + \chi_{1,2}, \quad \phi_i = \chi_i \text{ for } i \in \{3, \dots, N\} \quad (19.4)$$

where v_1 and v_2 are constants and, again, χ_i describe small excitations around the minimum of the potential.

(d) Express v_1 and v_2 in terms of R_{vac} and a mixing angle θ .

(e) Explain why, despite the fact that two fields obtain a non-zero expectation value, the above symmetry breaking pattern is equivalent to the previous one. Confirm this by rewriting the Lagrangian in terms of the fields χ_i and by diagonalising the mass matrix.

Exercise 20: Non-abelian gauge transformation**6 points**

In lecture 7 it was discussed why the gauge field \hat{A}_μ belongs to the Lie algebra for the case of an SU(2) gauge theory. This will be generalised in this exercise.

Consider a non-abelian gauge theory with a gauge group G . Under a non-abelian gauge transformation, the gauge field transforms as

$$\hat{A}_\mu(x) \rightarrow \hat{A}'_\mu(x) = U(x)\hat{A}_\mu(x)U^{-1}(x) + \frac{1}{ig}(\partial_\mu U^{-1}(x))U(x). \quad (20.1)$$

Recall that the gauge field is defined as $\hat{A}_\mu(x) = \sum_a A_\mu^a(x)t^a$, where the t^a are the generators of the Lie algebra \mathfrak{g} associated to the gauge group. In particular, this means that the gauge field belongs to the algebra, $\hat{A}_\mu \in \mathfrak{g}$. Now one has to show that also the gauge transformed field belongs to the algebra, $\hat{A}'_\mu \in \mathfrak{g}$.

The matrices t^a fulfil the commutation relations $[t^a, t^b] = F^{abc}t^c$ for some structure constants F^{abc} . Thus, the commutator of two generators is again a linear combination of generators, i.e. $[t^a, t^b] \in \mathfrak{g}$. The gauge transformations can be written as $U(x) = \exp(i \sum_a \theta^a(x)t^a)$. To make the notation more compact, we define the following shorthands for nested commutators:

$$\tilde{\Delta}_0(X, Y) = Y, \quad \Delta_0(Y, X) = Y, \quad (20.2)$$

$$\tilde{\Delta}_n(X, Y) = [X, \tilde{\Delta}_{n-1}(X, Y)], \quad \Delta_n(Y, X) = [\Delta_{n-1}(Y, X), X]. \quad (20.3)$$

(a) It will be shown in the next subquestions that

$$\begin{aligned} e^{X}Y e^{-X} &= \sum_{n=0}^{\infty} \frac{\tilde{\Delta}_n(X, Y)}{n!} \\ &= Y + [X, Y] + \frac{1}{2!}[X, [X, Y]] + \frac{1}{3!}[X, [X, [X, Y]]] + \dots \end{aligned} \quad (20.4)$$

and that

$$\begin{aligned} \frac{d}{dt}e^{X(t)} &= e^{X(t)} \left(\sum_{n=0}^{\infty} \frac{\Delta_n(\frac{dX}{dt}, X)}{(n+1)!} \right) \\ &= e^{X(t)} \left(\frac{dX}{dt} + \frac{1}{2!} [\frac{dX}{dt}, X] + \frac{1}{3!} [[\frac{dX}{dt}, X], X] + \dots \right). \end{aligned} \quad (20.5)$$

For now, use Eqs. (20.4) and (20.5) to show that $\hat{A}'_\mu \in \mathfrak{g}$.

(b) Prove Eq. (20.4). To that end, it is useful to define an auxiliary function $F(z) = e^{zX}Y e^{-zX}$. For $z = 1$ this becomes the left-hand side of Eq. (20.4). Make a power series ansatz for $F(z)$,

$$F(z) = \sum_{n=0}^{\infty} \frac{1}{n!} F_n z^n, \quad (20.6)$$

and use the derivative $\frac{dF(z)}{dz}$ to derive a recurrence relation for the coefficients F_n . Use this to show Eq. (20.4).

Now, prove Eq. (20.5). Use the following steps:

- (c) Show that the left-hand side of Eq. (20.5) can be written as

$$\frac{d}{dt}e^{X(t)} = \sum_{n=0}^{\infty} \frac{1}{(n+1)!} \sum_{k=0}^n X^{n-k}(t) \left(\frac{dX(t)}{dt} \right) X^k(t). \quad (20.7)$$

- (d) **Optional** Show that

$$[Y, X^n] = \sum_{k=0}^{n-1} \binom{n}{k} X^k \Delta_{n-k}(Y, X), \quad (20.8)$$

and conclude from it that

$$Y X^n = \sum_{k=0}^n \binom{n}{k} X^k \Delta_{n-k}(Y, X). \quad (20.9)$$

Proceed by induction over n . You may find the identity $\binom{n+1}{k+1} = \binom{n}{k} + \binom{n}{k+1}$ to be useful.

- (e) Use the previous result to further reexpress Eq. (20.7) as

$$\frac{d}{dt}e^{X(t)} = \sum_{n=0}^{\infty} \frac{1}{(n+1)!} \sum_{k=0}^n \sum_{j=0}^k \binom{k}{j} X^{n-j}(t) \Delta_j\left(\frac{dX(t)}{dt}, X(t)\right). \quad (20.10)$$

- (f) Show that the right-hand side of Eq. (20.5) can be transformed into

$$e^{X(t)} \left(\sum_{n=0}^{\infty} \frac{\Delta_n\left(\frac{dX}{dt}, X\right)}{(n+1)!} \right) = \sum_{n=0}^{\infty} \frac{1}{(n+1)!} \sum_{j=0}^n \binom{n+1}{j+1} X^{n-j}(t) \Delta_j\left(\frac{dX}{dt}, X\right). \quad (20.11)$$

Use for this the Cauchy product of two series,

$$\left(\sum_{n=0}^{\infty} a_n \right) \cdot \left(\sum_{m=0}^{\infty} b_m \right) = \sum_{n=0}^{\infty} \sum_{k=0}^n a_k b_{n-k}. \quad (20.12)$$

- (g) **Optional** Prove that for arbitrary coefficients $c_{n,j}$

$$\sum_{k=0}^n \sum_{j=0}^k \binom{k}{j} c_{n,j} = \sum_{j=0}^n \binom{n+1}{j+1} c_{n,j}. \quad (20.13)$$

Change the order of summation on the left-hand side and use the ‘‘hockey-stick identity’’ $\sum_{k=j}^n \binom{k}{j} = \binom{n+1}{j+1}$.

- (h) Use this result to show that the right-hand sides of Eqs. (20.10) and (20.11) are equal and, therefore, that Eq. (20.5) holds.