

# Introduction to Theoretical Particle Physics

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## Exercise Sheet 3

Issue: 25.10. – Submission: 31.10. @ 18:00 Uhr – Discussion: 05.11/06.11

### Exercise 6: Two-point correlation function

**3 points**

In the lecture you have seen how the matrix element  $\langle x_f, t_f | q(t_1) | x_i, t_i \rangle$ , where  $q(t_1)$  is the position operator in the Heisenberg representation, can be written as a path integral with an additional factor  $x(t_1)$  in the integrand:

$$\langle x_f, t_f | q(t_1) | x_i, t_i \rangle = \int [\mathcal{D}x(t)] x(t_1) e^{\frac{i}{\hbar} S} \Big|_{x(t_f)=x_f, x(t_i)=x_i}. \quad (6.1)$$

Show that a similar identity holds for the matrix element of two position operators at times  $t_1$  and  $t_2$ ,

$$\langle x_f, t_f | T q(t_1) q(t_2) | x_i, t_i \rangle = \int [\mathcal{D}x(t)] x(t_1) x(t_2) e^{\frac{i}{\hbar} S} \Big|_{x(t_f)=x_f, x(t_i)=x_i}, \quad (6.2)$$

Where  $T$  is the time-ordering operator,

$$T q(t_1) q(t_2) = \theta(t_1 - t_2) q(t_1) q(t_2) + \theta(t_2 - t_1) q(t_2) q(t_1). \quad (6.3)$$

Explain why the time-ordering operator has to be introduced.

### Exercise 7: Harmonic oscillator with external force

**9 points**

Consider a free harmonic oscillator with the Lagrangian

$$L_{\text{H.O.}}(\dot{q}, q) = \frac{1}{2} m \dot{q}(t)^2 - \frac{1}{2} m \omega^2 q(t)^2, \quad (7.1)$$

where  $m$  is the mass parameter and  $\omega$  is the frequency. The action of an external force  $m f(t)$  on the oscillator is modelled by an additional term in the Lagrangian

$$\delta L_{\text{ext}} = m f(t) q(t). \quad (7.2)$$

In the lectures it was shown that taking  $H \rightarrow (1 - i\epsilon)H$ , where  $\epsilon$  is infinitesimal and positive, the transition amplitude from the ground state in the remote past ( $t = -\infty$ ) to the ground state in the remote future ( $t = +\infty$ ) is given by

$$\langle 0|0 \rangle = \int [\mathcal{D}q] \exp \{iS\}, \quad (7.3)$$

where  $[\mathcal{D}q]$  denotes the usual path integral. In this exercise we study the path integral in the presence of an external force where the full Lagrangian is given by  $L_{\text{full}} = (L_{\text{H.O.}} + \delta L_{\text{ext}})$ .

- (a) What is the Hamiltonian  $H$  of the free harmonic oscillator? Show that the limit,  $H \rightarrow (1 - i\epsilon)H$ , is equivalent to the replacements

$$m \rightarrow m(1 + i\epsilon), \quad m\omega^2 \rightarrow m\omega^2(1 - i\epsilon). \quad (7.4)$$

Since  $\epsilon$  is infinitesimally small, keep only linear terms and neglect all terms of  $\mathcal{O}(\epsilon^2)$ .

- (b) Perform the shift of Eq. (7.4) on the H.O. Lagrangian (7.1) and write it in terms of the Fourier transformed variables

$$\tilde{q}(E) = \int_{-\infty}^{+\infty} dt e^{iEt} q(t), \quad q(t) = \int_{-\infty}^{+\infty} \frac{dE}{2\pi} e^{-iEt} \tilde{q}(E). \quad (7.5)$$

Show that the action,  $S = \int_{-\infty}^{+\infty} dt L_{\text{full}}$ , can be written in terms of an integral over one Fourier parameter,  $\tilde{E}$ ,

$$S = \frac{m}{2} \int_{-\infty}^{\infty} \frac{dE}{2\pi} \left[ (E^2 - \omega^2 + i\epsilon) \tilde{q}(E) \tilde{q}(-E) + \tilde{f}(E) \tilde{q}(-E) + \tilde{f}(-E) \tilde{q}(E) \right]. \quad (7.6)$$

*Hints:*

- ★ Use a separate Fourier parameter for each power of  $q(t)$  and perform the integral over time

$$\int_{-\infty}^{+\infty} dt e^{i(E+E')t} = (2\pi) \delta(E + E'). \quad (7.7)$$

- ★ Since the  $i\epsilon$  prescription is considered in the  $\epsilon \rightarrow 0$  limit, you can rescale  $i(E^2 + \omega^2)\epsilon \rightarrow i\epsilon$ .

- (c) Perform a shift of the variable  $\tilde{q}$ ,

$$\tilde{q}(E) \longrightarrow \tilde{x}(E) = \tilde{q}(E) + \frac{\tilde{f}(E)}{E^2 - \omega^2 + i\epsilon}. \quad (7.8)$$

Argue that the path integral measure does not change under this transformation, i.e.  $[\mathcal{D}q] = [\mathcal{D}x]$ . Verify that the action splits into two pieces, one that is quadratic in  $\tilde{x}$  and another which is independent of  $\tilde{x}$ . Show that

$$\langle 0|0 \rangle_f = \exp \left\{ -\frac{i}{2} m \int_{-\infty}^{+\infty} \frac{dE}{2\pi} \frac{\tilde{f}(E) \tilde{f}(-E)}{E^2 - \omega^2 + i\epsilon} \right\} \cdot \langle 0|0 \rangle_{f=0}. \quad (7.9)$$

- (d) The second factor on the right-hand-side of Eq. (7.9) is the transition amplitude in the absence of an external force and hence  $\langle 0|0 \rangle_{f=0} = 1$ . Perform an inverse Fourier transform and show that the transition amplitude can be written as

$$\langle 0|0 \rangle_f = \exp \left\{ \frac{i}{2} m \int_{-\infty}^{+\infty} dt dt' f(t) G(t-t') f(t') \right\}, \quad (7.10)$$

where

$$G(t-t') = - \int_{-\infty}^{+\infty} \frac{dE}{2\pi} \frac{e^{iE(t-t')}}{E^2 - \omega^2 + i\epsilon}. \quad (7.11)$$

- (e) Verify that the function  $G(t - t')$  of Eq. (7.11) is a Green's function for the harmonic oscillator, i.e. check that after taking  $\epsilon \rightarrow 0$  limit the following equation holds

$$(\partial_t^2 + \omega^2) G(t - t') = \delta(t - t'). \quad (7.12)$$

- (f) Using the residue theorem, show that the Green's function of Eq. (7.11) can be written explicitly as

$$G(t - t') = \frac{i}{2\omega} e^{-i\omega|t-t'|}. \quad (7.13)$$

*Hint:* See the discussion in Sec. 2.4 of Peskin & Schroeder.