

## Two-loop master-integrals

We define all integrals in the Euclidean space with  $D = 4 - 2\epsilon$  dimensions. In each integral we only compute a piece of the imaginary part, which includes only cuts crossing the line with mass  $x$  (corresponding to the normalized Higgs boson mass present in the denominator  $D_1$ ). Master integrals were identified with the help of Laporta algorithm. Integral U6 is a linear combination of U1, U1a, and U8, other integrals are linearly independent.

The forward scattering kinematics is assumed: the external momenta are on-shell,  $p_1^2 = p_2^2 = 0$ , normalized by  $(p_1 + p_2)^2 = -1$ . Each integral depends on  $x$ ,  $0 < x < 1$ . The topologies are defined as follows:

$$\text{TT}\#(a_1, \dots, a_7; x) = \frac{1}{\pi \mathcal{F}^2} \mathfrak{S}^* \int \frac{d^D v_1 d^D v_2}{(2\pi)^{2D}} \frac{1}{D_1^{a_1} \dots D_7^{a_7}}, \quad D_1 = v_1^2 + x, \quad D_2 = v_2^2, \quad D_i = k_i^2, \quad k = 3..7$$

with the loop factor  $\mathcal{F} = \frac{\Gamma(1+\epsilon)}{(4\pi)^{2-\epsilon}}$ , symbol  $\mathfrak{S}^*$  denoting the sum over appropriate cuts, and

| # | $k_3$                   | $k_4$                   | $k_5$                   | $k_6$             | $k_7$                   |
|---|-------------------------|-------------------------|-------------------------|-------------------|-------------------------|
| A | $v_1 + v_2 - p_1 - p_2$ | $v_1 + v_2 - p_2$       | $v_2 - p_1$             | $v_2 - p_2$       | $v_1 + v_2 - p_1$       |
| B | $v_1 + v_2 - p_1 - p_2$ | $v_2 - p_1$             | $v_1 - p_1$             | $v_1 + v_2 - p_1$ | $v_2 - p_2$             |
| C | $v_1 + v_2 - p_1 - p_2$ | $v_2 - p_1$             | $v_1 - p_1$             | $v_1 - p_2$       | $v_1 + v_2 - p_1$       |
| D | $v_1 + v_2 - p_1 - p_2$ | $v_2 - p_1$             | $v_1 - p_1$             | $v_1 - p_2$       | $v_2 - p_2$             |
| E | $v_2 + p_1$             | $v_2 - p_2$             | $v_1 + v_2 + p_1$       | $v_1 + v_2 - p_2$ | $v_1 + v_2$             |
| F | $v_2 + p_1$             | $v_2 - p_2$             | $v_1 + v_2 + p_1$       | $v_1 + p_1 + p_2$ | $v_1 + p_1$             |
| G | $v_2 + p_1$             | $v_2 - p_2$             | $v_1 + v_2 + p_1$       | $v_1 + p_1 + p_2$ | $v_1 + p_2$             |
| H | $v_2 + p_1$             | $v_2 - p_2$             | $v_1 + v_2 + p_1$       | $v_1 + v_2 - p_2$ | $v_1 + v_2 + p_1 - p_2$ |
| J | $v_2 + p_1$             | $v_1 + v_2 + p_1 + p_2$ | $v_1 + v_2 + p_1$       | $v_1 + p_1 + p_2$ | $v_1 + p_1$             |
| K | $v_2 + p_1$             | $v_1 - p_2$             | $v_1 + v_2 + p_1$       | $v_1 + v_2 - p_2$ | $v_1 + v_2$             |
| L | $v_2 + p_1$             | $v_2 - p_2$             | $v_1 + v_2 + p_1 - p_2$ | $v_1 + v_2 + p_1$ | $v_1 + p_1$             |

For example, using also the Cutkoski relation, integral U7a is

$$U7a(x, \epsilon) = \frac{1}{2\pi \mathcal{F}^2} \int \frac{d^D v_1 d^D v_2 v_2^2 (2\pi)^3 \delta(v_1^2 + x) \delta((v_2 - p_2)^2) \delta((v_1 + v_2 + p_1)^2)}{(2\pi)^{2D} (v_2 + p_1)^2 (v_1 + p_2)^2}. \quad (1)$$

We find the integrals by solving the system of differential equations  $d U_i(x, \epsilon)/d x = \sum c_j(x, \epsilon) U_j(x, \epsilon)$ , in terms of harmonic polylogarithms (HPLs), and fix the integration constants with the expansion of each integral in  $(1-x) \ll 1$  to a few orders: two in case of U1 (U1a) and U7 (U7a), and one for the remaining integrals. For most integrals, we have managed to over-constrain the derivation and produce three or more expansion terms.

The file `masters.m` contains Mathematica expressions according to the list below. For each integral, we present a few orders in the fixed-order  $\epsilon$ -expansion, e.g., `U1[x, ep]`, and the singular limit regulated by a factor  $(1-x)^{n\epsilon}$  with integer  $n$ , e.g. `U1[x -> 1, ep]`.

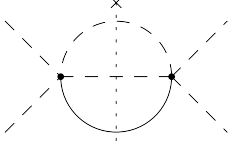
In order to use those expression, one has to add and subtract the singularity. E.g. if the limit is  $U(x \rightarrow 1, \epsilon) = C(\epsilon)(1-x)^{k-m\epsilon}$ , the regularized solution is

$$U_r(x) = U(x) + C(\epsilon)(1-x)^{k+1} \left[ \left\{ \frac{\delta(1-x)}{m\epsilon} + \left[ \frac{1}{1-x} \right]_+ + \dots \right\} - \left\{ \frac{1}{1-x} - \frac{m \ln 1-x}{1-x} + \dots \right\} \right]. \quad (2)$$

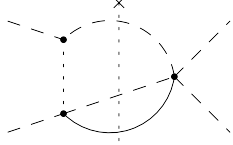
The expression in the square brackets represents the difference  $[(1-x)^{-1-m\epsilon} - (1-x)^{-1-m\epsilon}]$ , where the first term is expanded in plus-distributions and the second term is expanded ‘naively’, canceling the non-integrable divergences in  $U(x)$ .

Integrals U17 and U12 have two possible cuts, presented separately as U12d (two-particle) and U12t (three-particle cut), and U17d and U17t.

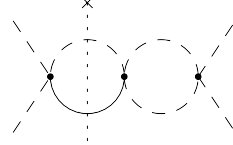
U1 = TTA(1,1,1,0,0,0,0)



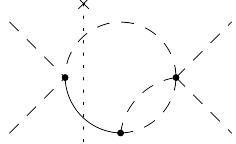
U1a = TTA(1,1,1,-1,0,0,0)



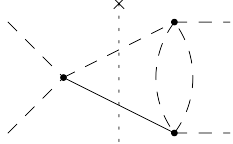
U2 = TTF(1,0,1,1,0,1,0)



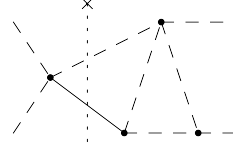
U3 = TTF(1,0,1,0,1,1,0)



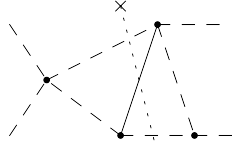
U4 = TTF(1,1,0,0,1,1,0)



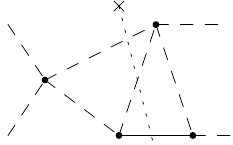
U5 = TTF(1,1,1,0,1,1,0)



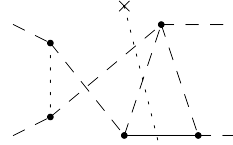
U6 = TTE(1,0,1,1,1,0,1)



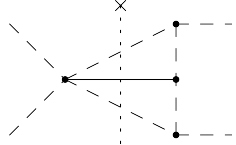
U7 = TTG(1,0,1,1,1,0,1)



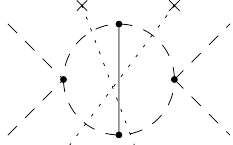
U7a = TTG(1,-1,1,1,1,0,1)



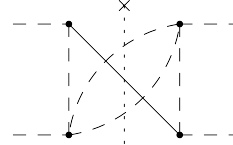
U8 = TTA(1,1,1,1,0,1,0)



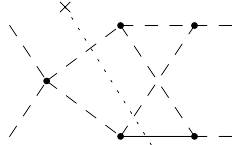
U9 = TTE(1,0,1,1,1,1,0)



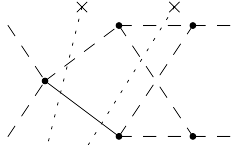
U10 = TTC(1,1,1,0,1,1,0)



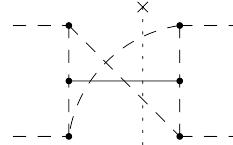
U11 = TTK(1,1,1,1,1,1,0)



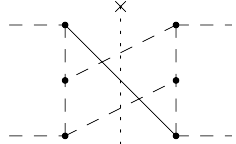
U12 = TTJ(1,1,1,1,1,1,0)



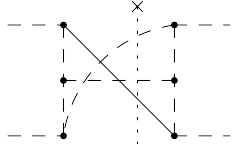
U13 = TTA(1,1,1,1,1,1,1)



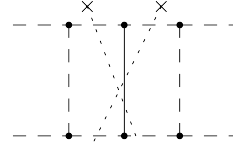
U14 = TTC(1,1,1,1,1,1,1)



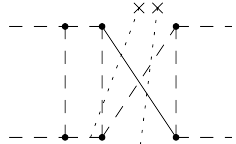
U15 = TTD(1,1,1,1,1,1,1)



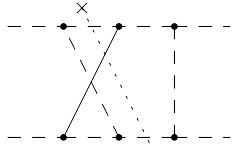
U16 = TTE(1,1,1,1,1,1,1)



U17 = TTG(1,1,1,1,1,1,1)



U18 = TTK(1,1,1,1,1,1,1)



U19 = TTH(1,1,1,1,1,1,1)

